

COMPUTATIONAL SEISMOLOGY



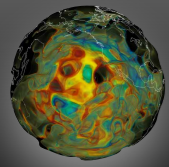
Imaging Earth's interior with
3D wave simulations

Ebru Bozdağ

Colorado School of Mines

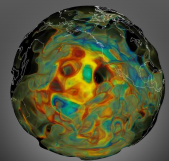
with contributions from:

Heiner Igel, Carene Larmat, Daniel Peter, Carl Tape, Jeroen Tromp

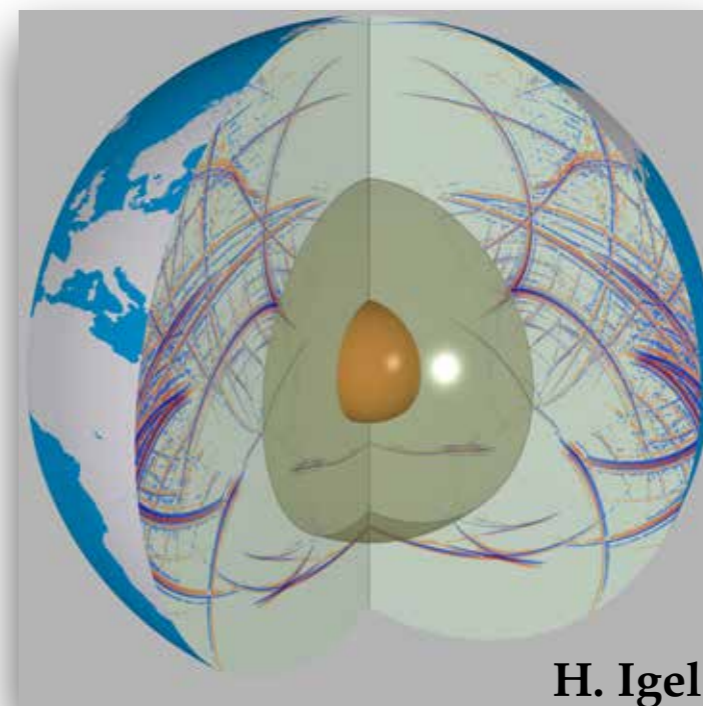
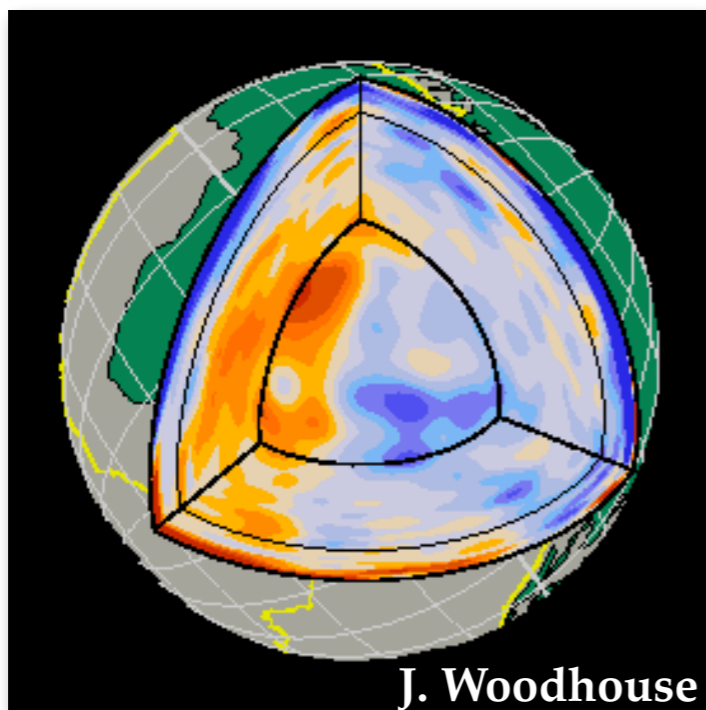


Seismology in a 3D planet

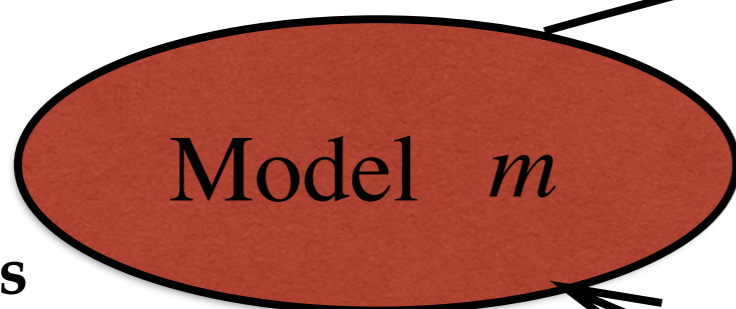
- **Forward problem**
 - ray theoretical methods (Sanne)
 - quasi-analytic methods - normal modes (Jessica)
 - **numerical methods**
- **Inverse problem (Ved)**
 - **gradient-based optimizations**
 - **seismic tomography with 3D wave simulations**
- **Current status in global imaging**
- **Challenges, future directions..**



Seismic modeling



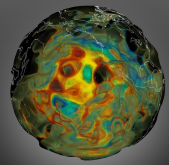
Forward problem



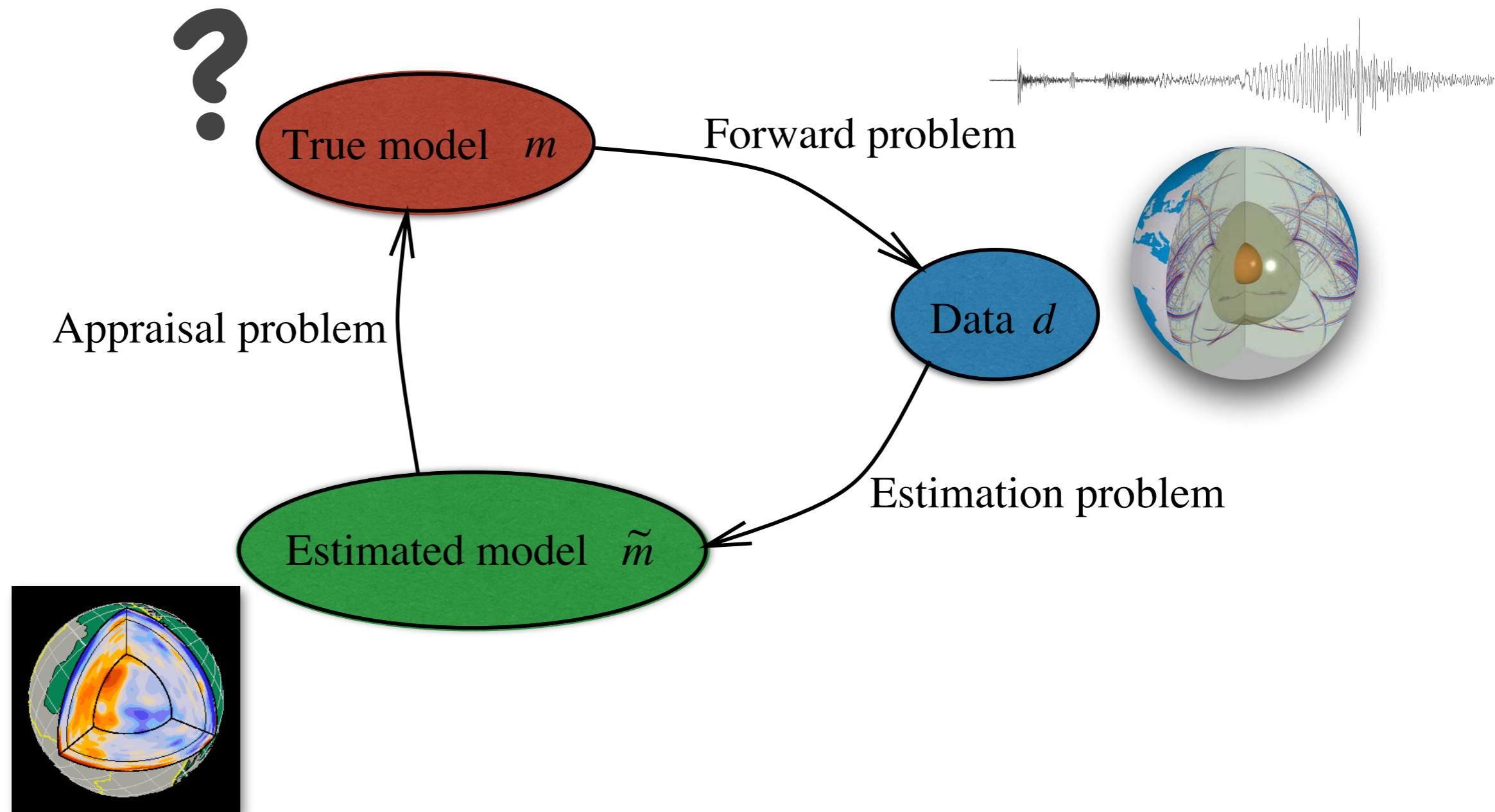
wavespeeds
density
anisotropy
anelasticity

seismograms
or
secondary observables

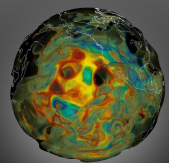
Inverse problem



Seismic modeling



inversion = estimation problem + appraisal problem



Seismic period (frequency) band

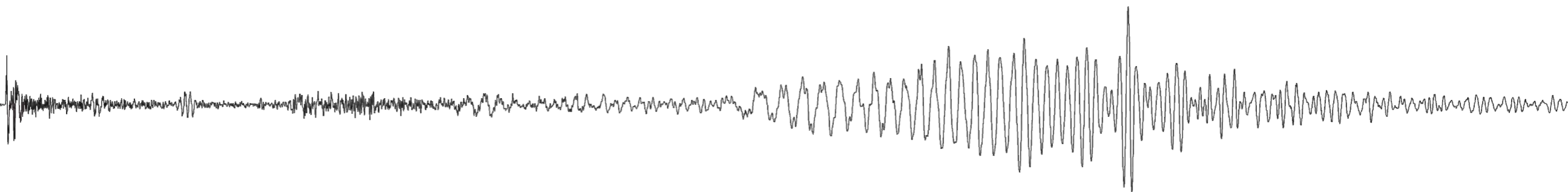
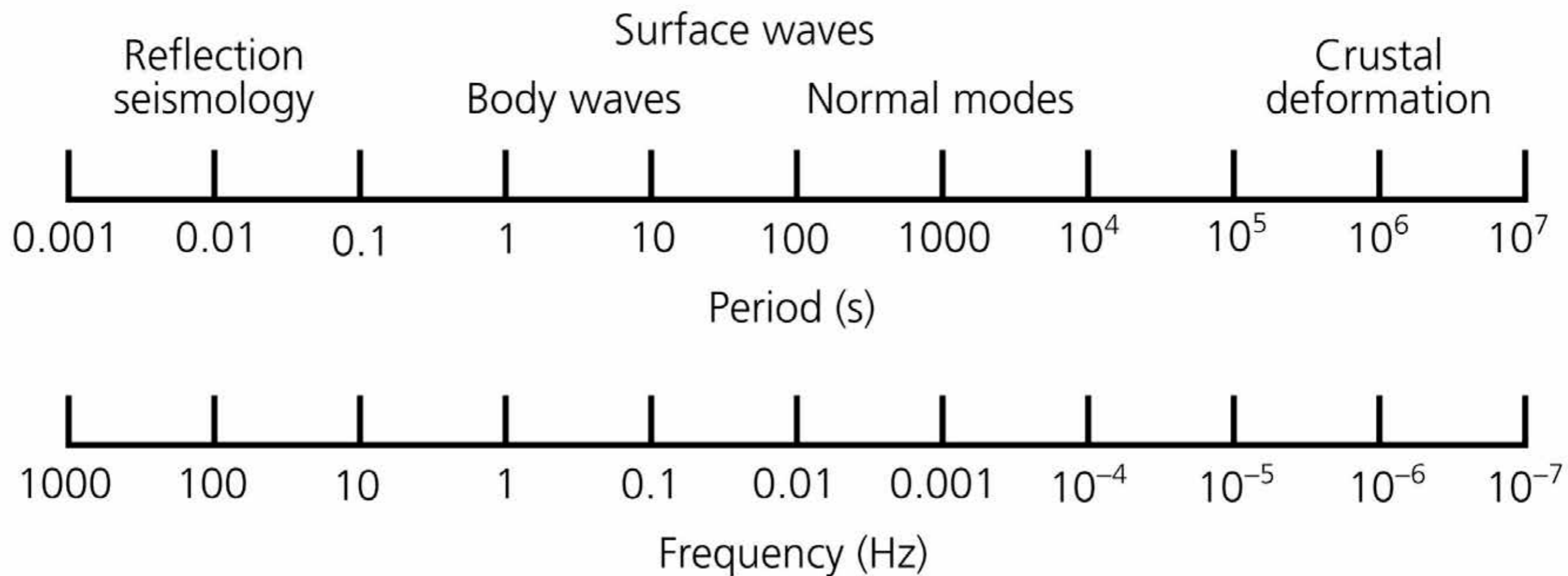
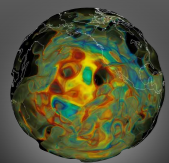


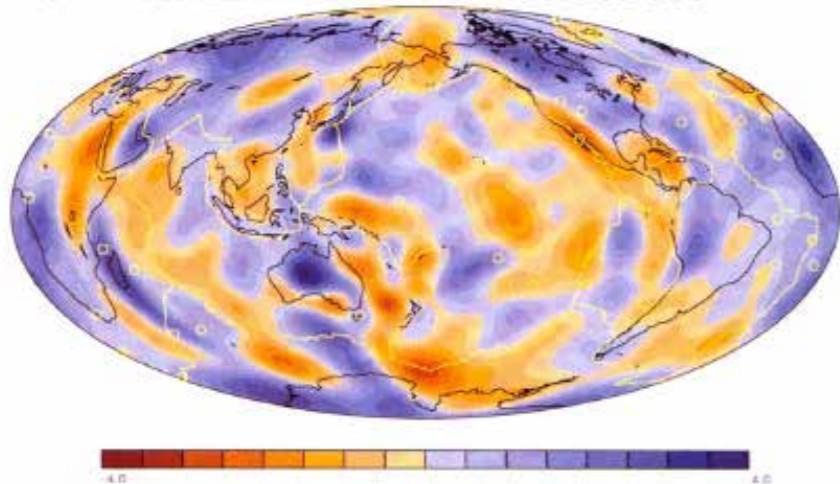
Figure 2.4-7: Seismic spectrum for various studies.





Ray theory vs. finite-frequency theory

(c) RAYLEIGH Period = 149.124 sec



Trampert & Woodhouse, 1995, GJI

Ray theory is a high-frequency approximation

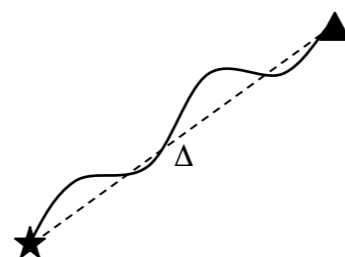
Its validity depends on:

- 1) frequency of waves
- 2) propagation distance of waves
- 3) smoothness of heterogeneities

see e.g., Wang & Dahlen, 1995, GJI

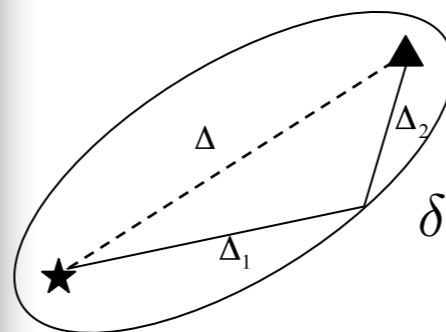
great circle approximation

$$\delta\phi = -\frac{\omega}{c_0} \int_0^\Delta \frac{\delta c}{c_0}(\theta, \varphi) d\Delta$$



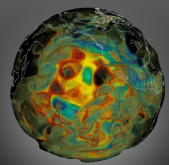
exact ray theory

$$\delta\phi = -\frac{\omega}{c_0} \int_{ray} \frac{\delta c}{c_0}(\theta, \varphi) ds$$



finite frequency theory

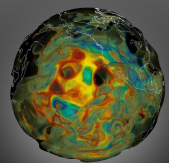
$$\delta\phi = -\frac{\omega}{c_0} \int \int_{sphere} K(\theta, \varphi) \frac{\delta c}{c_0}(\theta, \varphi) d\theta d\varphi$$



Ray theoretical phase and amplitude anomalies

$$\begin{aligned}\delta\psi &= \int_{\text{ray}} \frac{\omega}{c(\theta, \phi)} ds - \frac{\omega a}{c_0} \Delta \\ &= \frac{\omega a}{c_0} \int_0^\Delta \left[\frac{c_0}{c(\theta, \phi)} \left\{ \frac{v(\phi)^2}{[1 + \gamma^2(\phi)]^2} + \frac{1}{1 + \gamma^2(\phi)} \right\}^{1/2} - 1 \right] d\phi\end{aligned}$$

$$\ln A = \frac{1}{2} \operatorname{cosec} \Delta \int_0^\Delta \sin(\Delta - \phi) c_0^{-1} [\sin \phi \partial_\theta^2 - \cos \phi \partial_\phi] \delta c d\phi$$



Phase space of scattering

wavenumber

$$k = \frac{2\pi}{\lambda}$$

wavelength

$$c = \lambda f$$

a : dominant wavelength of the heterogeneous medium

L : propagation distance

$$ka = 2\pi \frac{a}{\lambda}$$

Correlation length

1,000

Homogeneous media

$$ka \gg 1$$

$$a \gg \lambda$$

Ray theory

100

Computational wave propagation

Random media theory

10

$$a \approx \lambda$$

1

Effective medium theory

$$a \ll \lambda$$

1

10

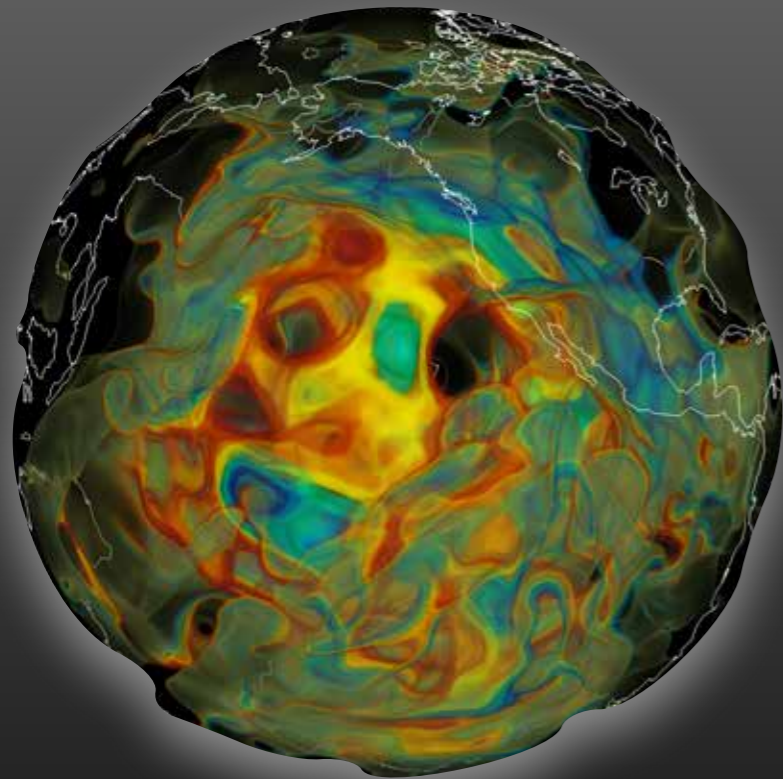
100

1,000

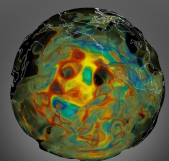
10,000

Propagation distance

$$kL = 2\pi \frac{L}{\lambda}$$



Numerical seismic wave simulations



Computational Infrastructure for Geodynamics

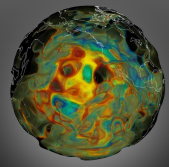
Short-Term Crustal Dynamics
 Long-Term Tectonics
 Mantle Convection
 Seismology
 Geodynamo
 Computational Science

SEISMOLOGY

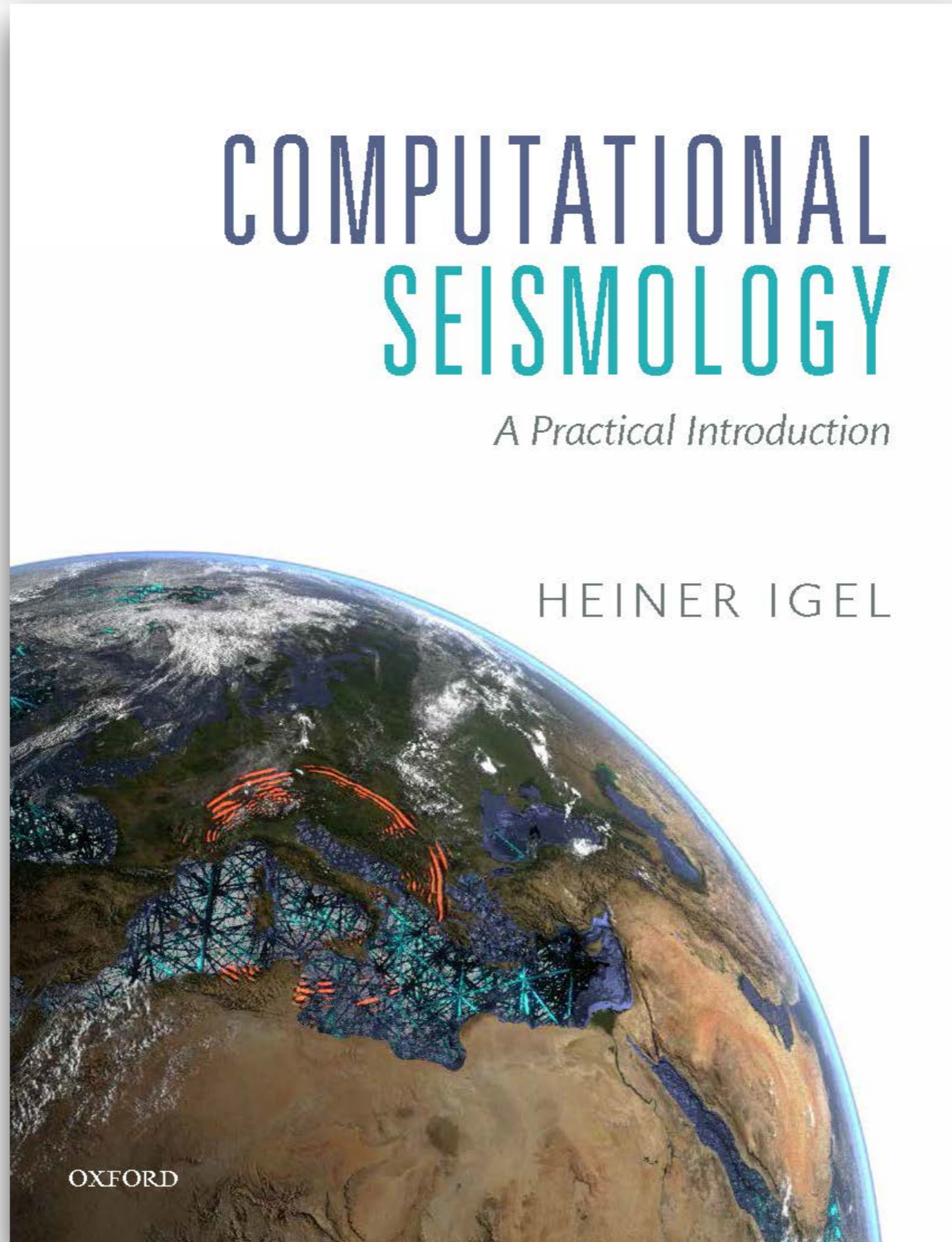
SPECFEM3D GLOBE	SPECFEM3D_GLOBE simulates global and regional (continental-scale) seismic wave propagation.	D_CIG
SPECFEM3D Cartesian	SPECFEM3D Cartesian simulates acoustic (fluid), elastic (solid), coupled acoustic/elastic, poroelastic or seismic wave propagation in any type of conforming mesh of hexahedra (structured or not.) It can, for instance, model seismic waves propagating in sedimentary basins or any other regional geological model following earthquakes. It can also be used for non-destructive testing or for ocean acoustics.	D_CIG
BurnMan	BurnMan is an open source mineral physics toolbox written in Python to determine seismic velocities for the lower mantle. BurnMan calculates the isotropic thermoelastic moduli by solving the equations-of-state for a mixture of minerals defined by the user. The user may select from a list of minerals applicable to the lower mantle included or easily define one of their own.	D_CONTRIB
AxiSEM	AxiSEM is a parallel spectral-element method for 3D (an-)elastic, anisotropic and acoustic wave propagation in spherical domains. It requires axisymmetric background models and runs within a 2D computational domain, thereby reaching all desired highest observable frequencies (up to 2Hz) in global seismology. The Fortran2003 (OpenMP, MPI) code scales very well on supercomputers and is available under the GPL license. For more details see the website at www.axisem.info .	D_CONTRIB
SPECFEM3D GEOTECH	SPECFEM3D_GEOTECH is an open-source command-driven software for 3D slope stability analysis and simulation of 3D multistage excavation based on the spectral-element method.	D_CONTRIB
SW4	SW4 implements substantial capabilities for 3-D seismic modeling, with a free surface condition on the top boundary, absorbing super-grid conditions on the far-field boundaries, and an arbitrary number of point force and/or point moment tensor source terms. Each source time function can have one of many predefined analytical time dependencies, or interpolate a user defined discrete time series.	D_CONTRIB
SPECFEM2D	SPECFEM2D simulates forward and adjoint seismic wave propagation in two-dimensional acoustic, (an)elastic, poroelastic or coupled acoustic-(an)elastic-poroelastic media, with Convolution PML absorbing conditions.	D_CONTRIB
SEISMIC_CPML	SEISMIC_CPML is a set of eleven open-source Fortran90 programs to solve the two-dimensional or three-dimensional isotropic or anisotropic elastic, viscoelastic or poroelastic wave equation using a finite-difference method with Convolutional or Auxiliary Perfectly Matched Layer (C-PML or ADE-PML) conditions, developed by Dimitri Komatitsch and Roland Martin from CNRS, France. Contributions by other authors have recently been added.	S_CONTRIB
SPECFEM1D	SPECFEM1D simulates seismic wave propagation in a one-dimensional heterogeneous medium. It is a small code that allows users to learn how a spectral-element program is written.	S_CONTRIB
Flexwin	The FLEXWIN software package automates the time-window selection problem for seismologists. It operates on pairs of observed and synthetic single component seismograms, defining windows that cover as much of a given seismogram as possible, while avoiding portions of the waveform that are dominated by noise.	A
Mineos	Mineos computes synthetic seismograms in a spherically symmetric non-rotating Earth by summing normal modes.	A

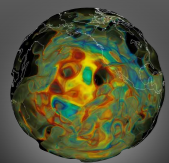


<https://geodynamics.org/cig/software/#seismo>



Reference book suggestion





Jupyter notebooks

http://krischer.github.io/seismo_live/

Seismo-Live

Live Jupyter Notebooks
for Seismology

25 of 25

Seismo-Live containers are
currently available

*Not persistent. Will be deleted after being idle for 30
minutes.*

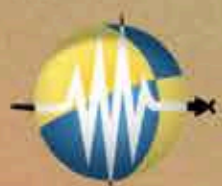
Launch Seismo-Live



Help - What is Seismo-Live?



More Information



ObsPy

A Python Framework for Seismology

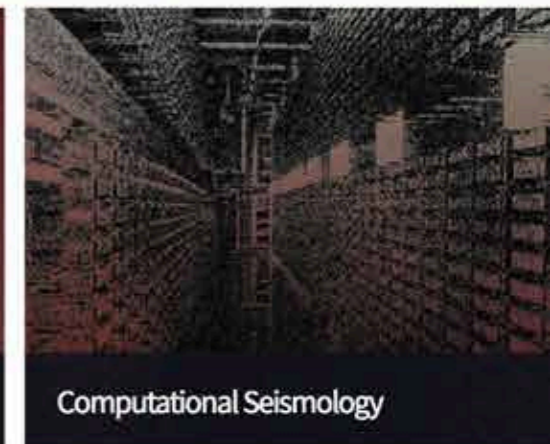
ObsPy



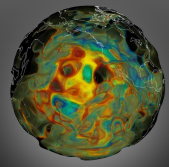
ASDF



Instaseis

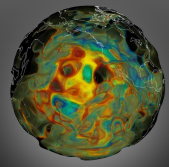


Computational Seismology



Numerical wave simulations

- Forward problem in seismology
- Numerical methods
 - **Finite difference**
 - Pseudo-spectral
 - Finite elements
 - **Spectral elements**
- How to use a black box?



Forward problem in seismology

Equation of motion:

$$\rho \partial_t^2 \mathbf{s} - \nabla \cdot \mathbf{T} = \mathbf{f}$$

Earthquake source:

$$\mathbf{f} = \mathbf{M} \cdot \nabla \delta(\mathbf{x} - \mathbf{x}_s) S(t)$$

Hook's law:

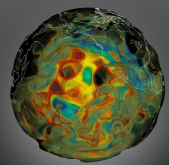
$$\mathbf{T} = \mathbf{c} : \epsilon$$

Boundary conditions:

$$\hat{\mathbf{n}} \cdot \mathbf{T} = 0$$

Initial conditions:

$$\mathbf{s}(\mathbf{x}, 0) = 0 \quad \partial_t \mathbf{s}(\mathbf{x}, 0) = 0$$



Finite-difference method

Equation of motion:

$$\rho \partial_t^2 \mathbf{s} - \nabla \cdot \mathbf{T} = \mathbf{f}$$

density
displacement
stress tensor
source term

Strong form:

$$\partial_t \mathbf{v} = \rho^{-1} (\nabla \cdot \mathbf{T} + \mathbf{f}) \quad \partial_t \mathbf{T} = \mathbf{c} : \nabla \mathbf{v}$$

$$\mathbf{v} = \partial_t \mathbf{s}$$

1D wave equation in a homogeneous medium

$$\partial_t^2 s(x, t) = \alpha^2 \partial_x^2 s(x, t)$$

P-wave speed S-wave speed

$$\alpha = \sqrt{\frac{\lambda + 2\mu}{\rho}} \quad \left(\beta = \sqrt{\frac{\mu}{\rho}} \right)$$

$$c = \lambda + 2\mu \quad (c = \mu)$$

Velocity-stress formulation

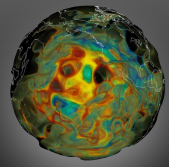
$$\partial_t v(x, t) = \rho^{-1} \partial_x T(x, t)$$

$$v(x, t) = \partial_t s(x, t)$$

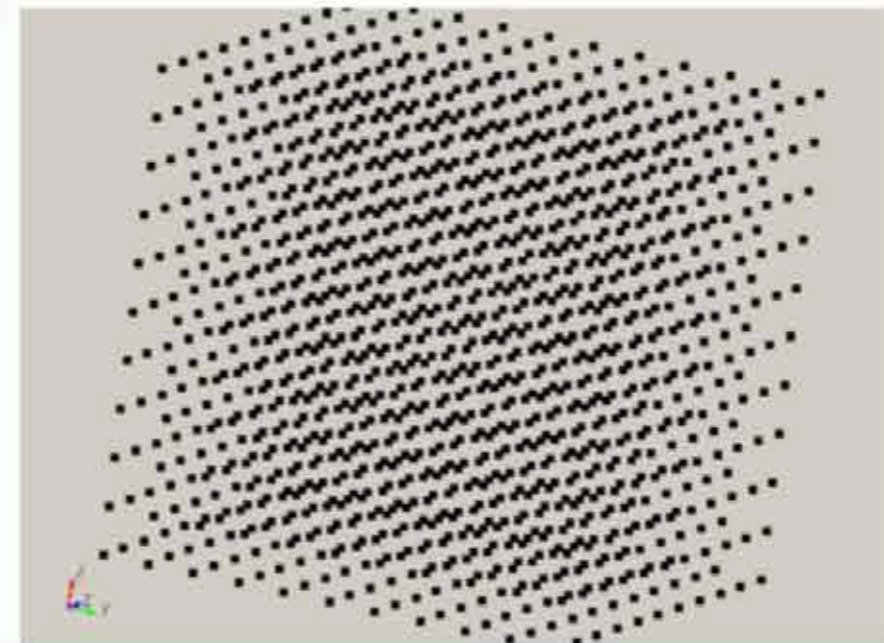
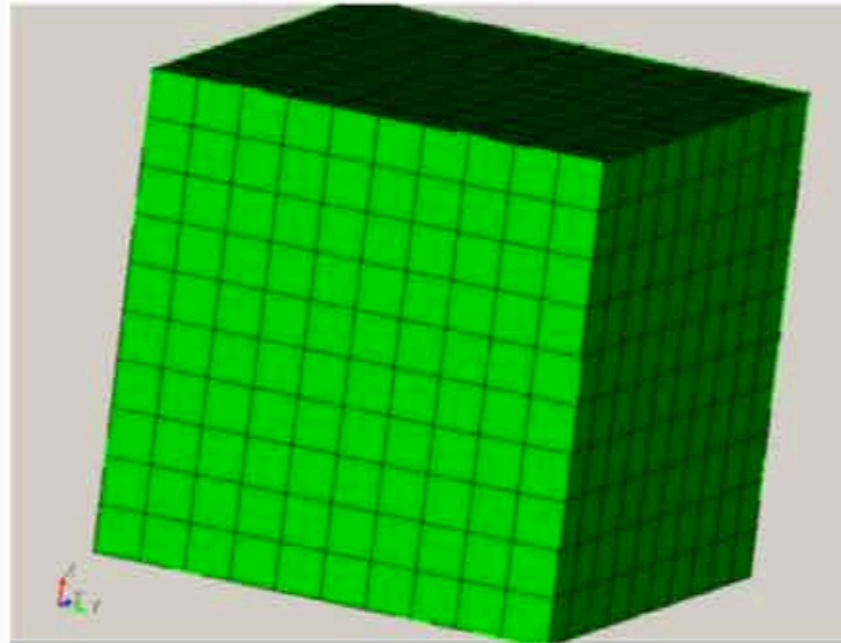
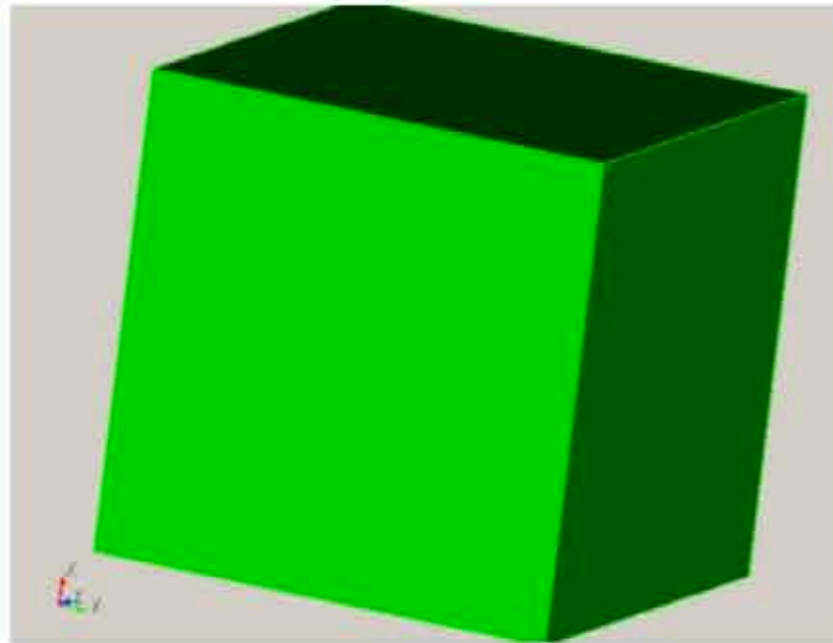
where

$$\partial_t T(x, t) = c \partial_x v(x, t)$$

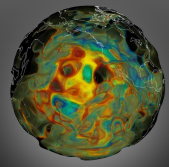
$$T(x, t) = c \partial_x s(x, t)$$



Discretized Earth



Discretise in both space & time!



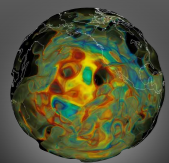
Finite difference & Taylor series

Accuracy of numerical simulations

Taylor series

$$f(x + dx) = f(x) + f'(x)dx + \frac{1}{2!}f''(x)dx^2 + \frac{1}{3!}f'''(x)dx^3 + \dots$$

$$f(x - dx) = f(x) - f'(x)dx + \frac{1}{2!}f''(x)dx^2 - \frac{1}{3!}f'''(x)dx^3 + \dots$$



Finite difference & Taylor series

$$\partial_t v(x, t) = \rho^{-1} \partial_x T(x, t) \quad \partial_t T(x, t) = c \partial_x v(x, t)$$

$$\frac{f(x + dx) - f(x)}{dx} = f'(x) + O(dx)$$

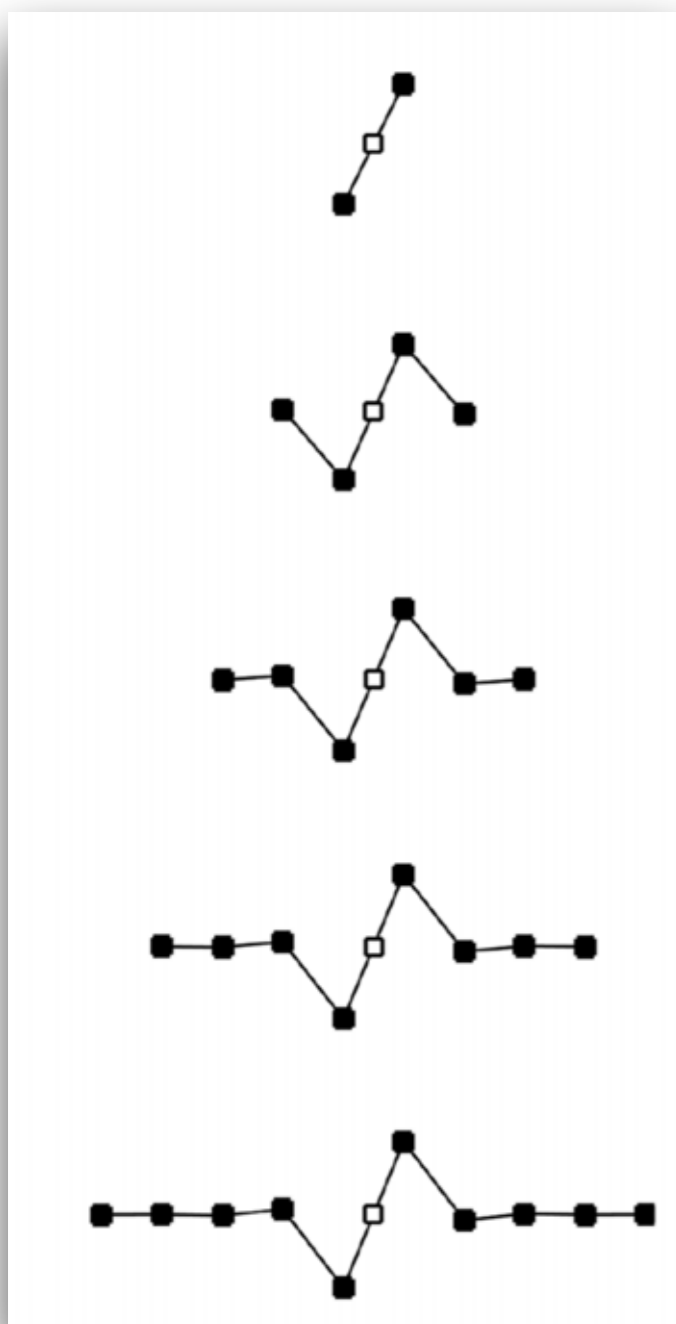
**Forward
difference**

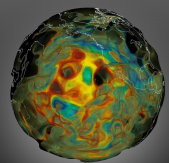
$$\frac{f(x) - f(x - dx)}{dx} = f'(x) + O(dx)$$

**Backward
difference**

$$\frac{f(x + dx) - f(x - dx)}{2dx} = f'(x) + O(dx^2)$$

**Central
difference**





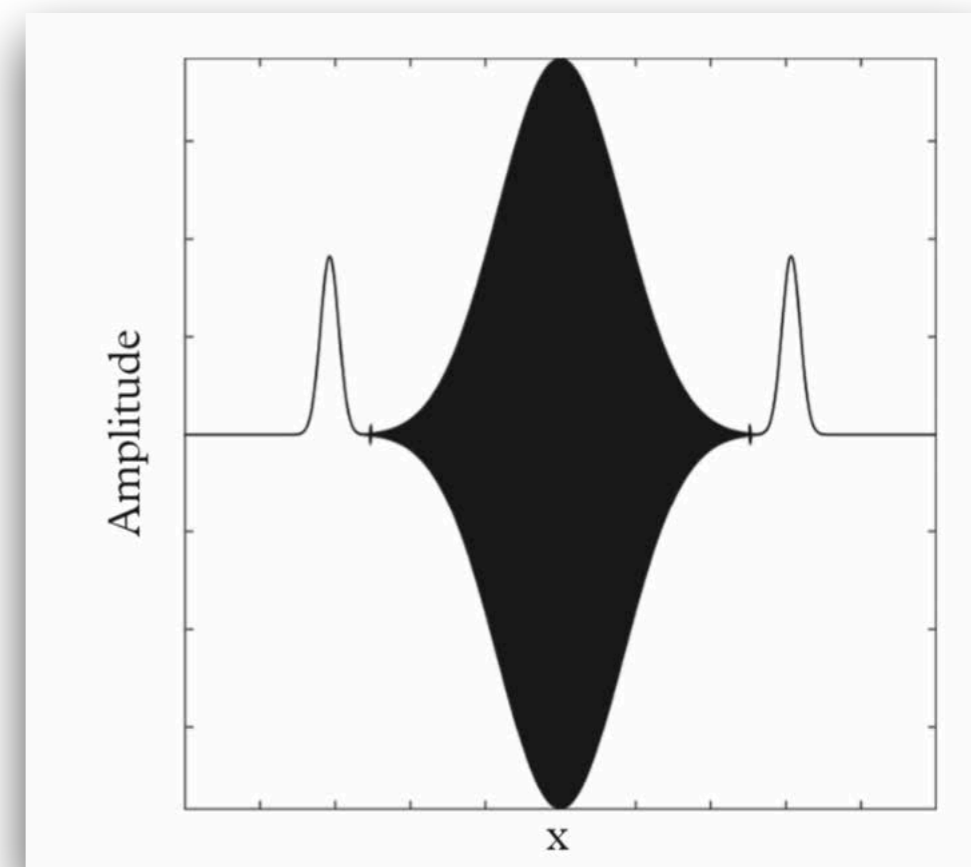
Stability

Von Neumann Analysis of stability:

$$\frac{p_j^{n+1} - 2p_j^n + p_j^{n-1}}{dt^2} = c_j^2 \left[\frac{p_{j+1}^n - 2p_j^n + p_{j-1}^n}{dx^2} \right]$$

with $p_j^n = e^{i(kjdx - \omega ndt)}$

$$\rightarrow \sin\left(\omega \frac{dt}{2}\right) = c \frac{dt}{dx} \sin\left(k \frac{dx}{2}\right)$$



Courant–Friedrichs–Lewy

(CFL) criterion

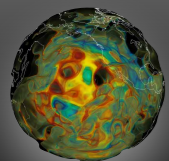
$$\epsilon = c \frac{dt}{dx} \leq 1$$

sampling rate in time

sampling rate in space

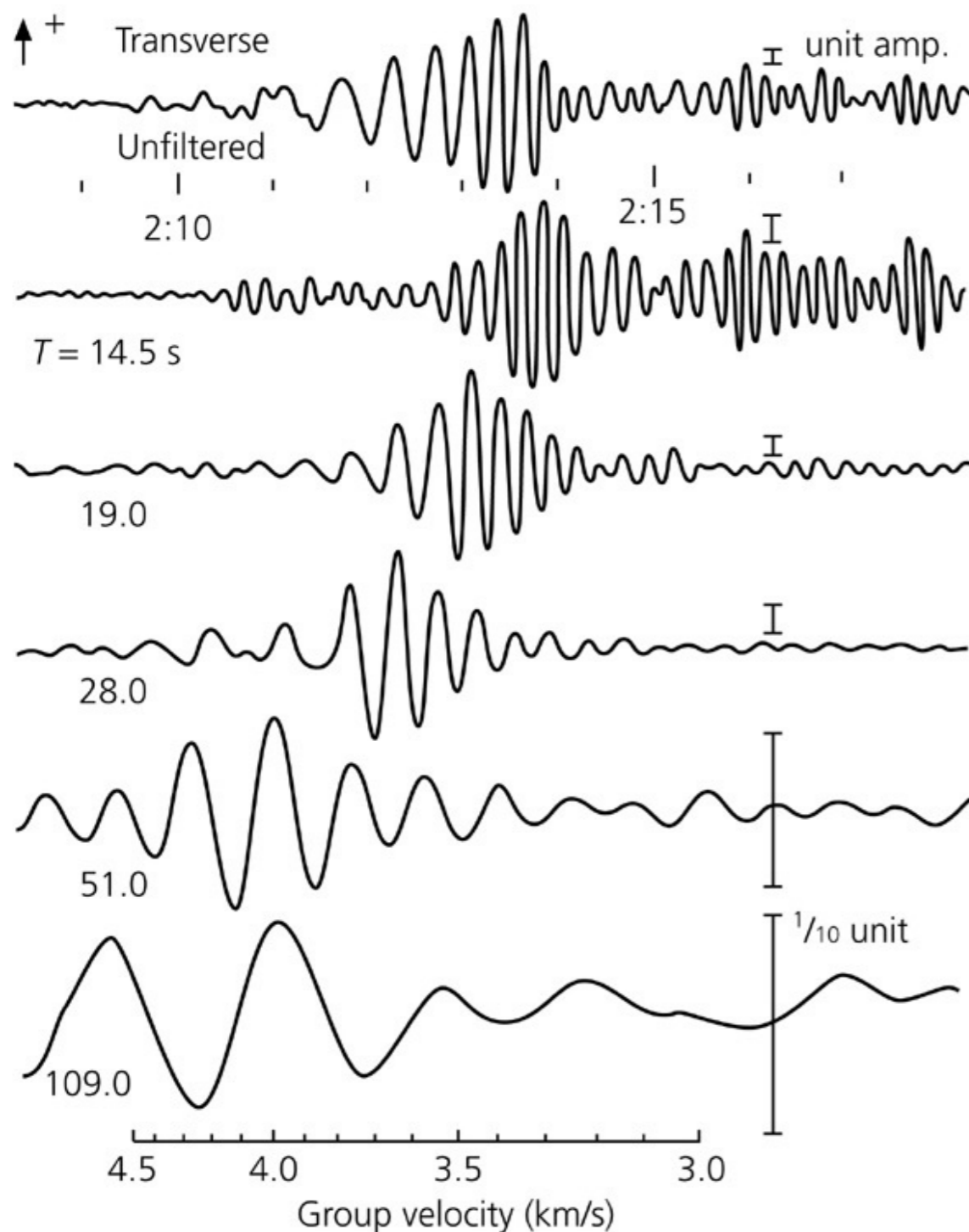
wavespeed

$$\epsilon > 1$$



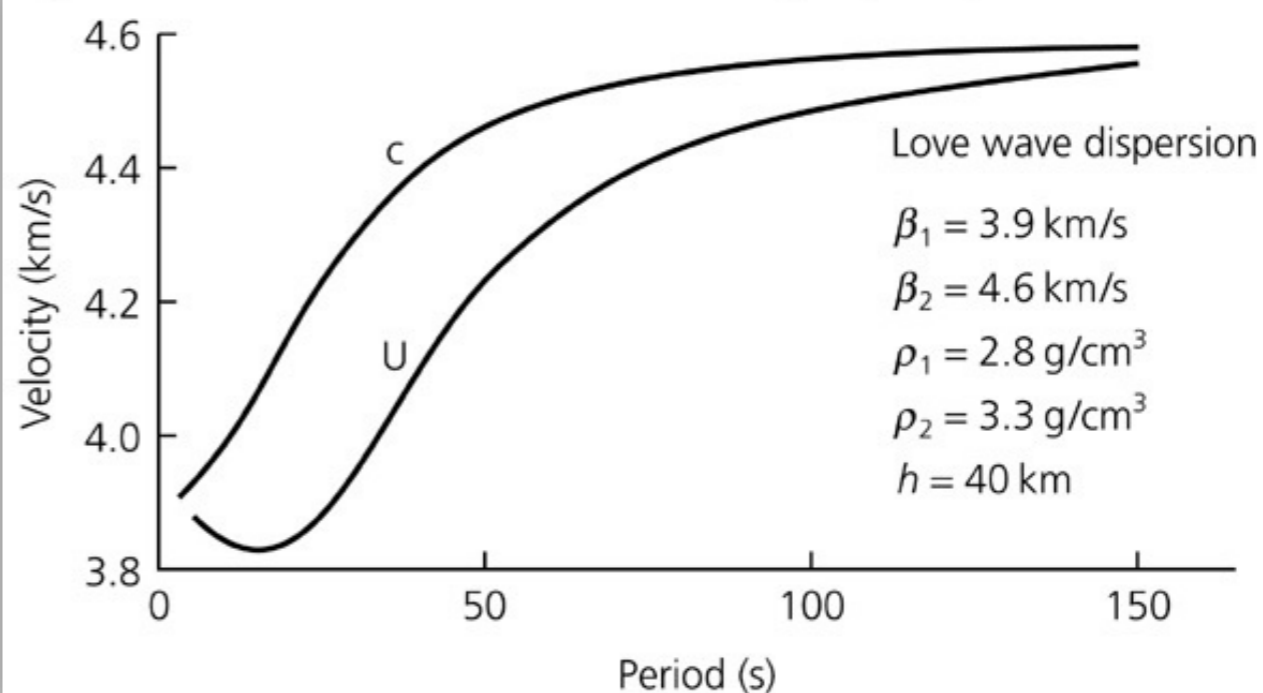
Wave dispersion

Figure 2.8-4: Example of Love wave group velocity dispersion through bandpass filtering.

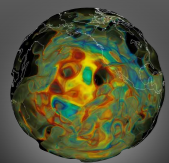


Stein & Wysession 2003

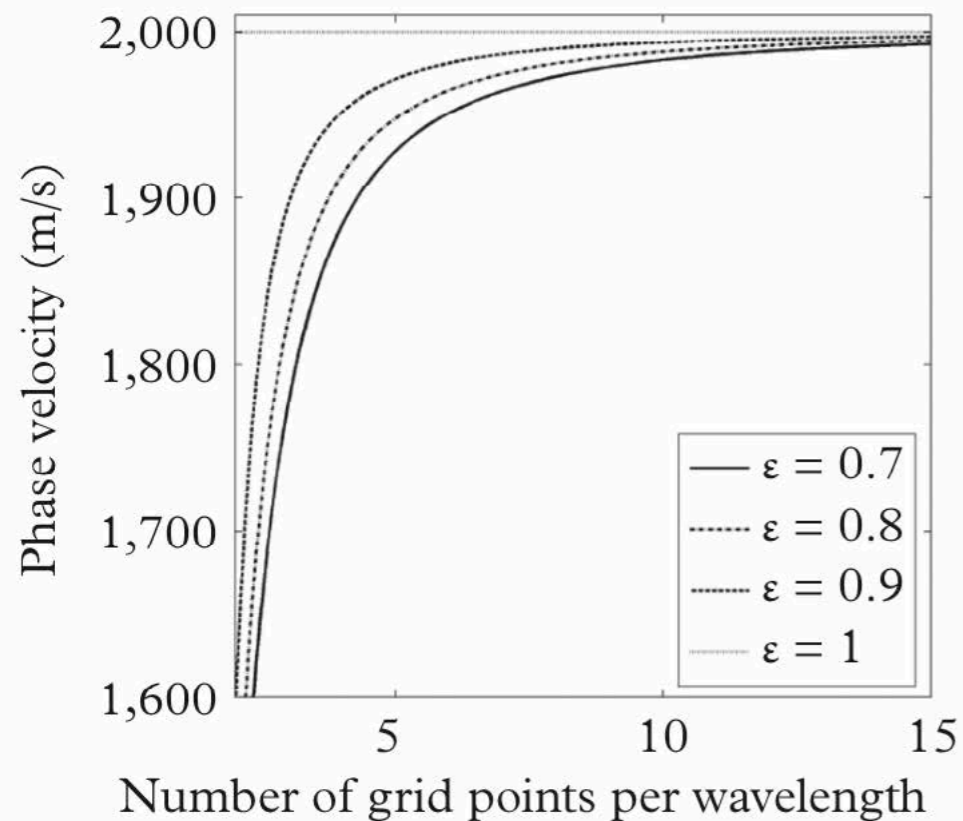
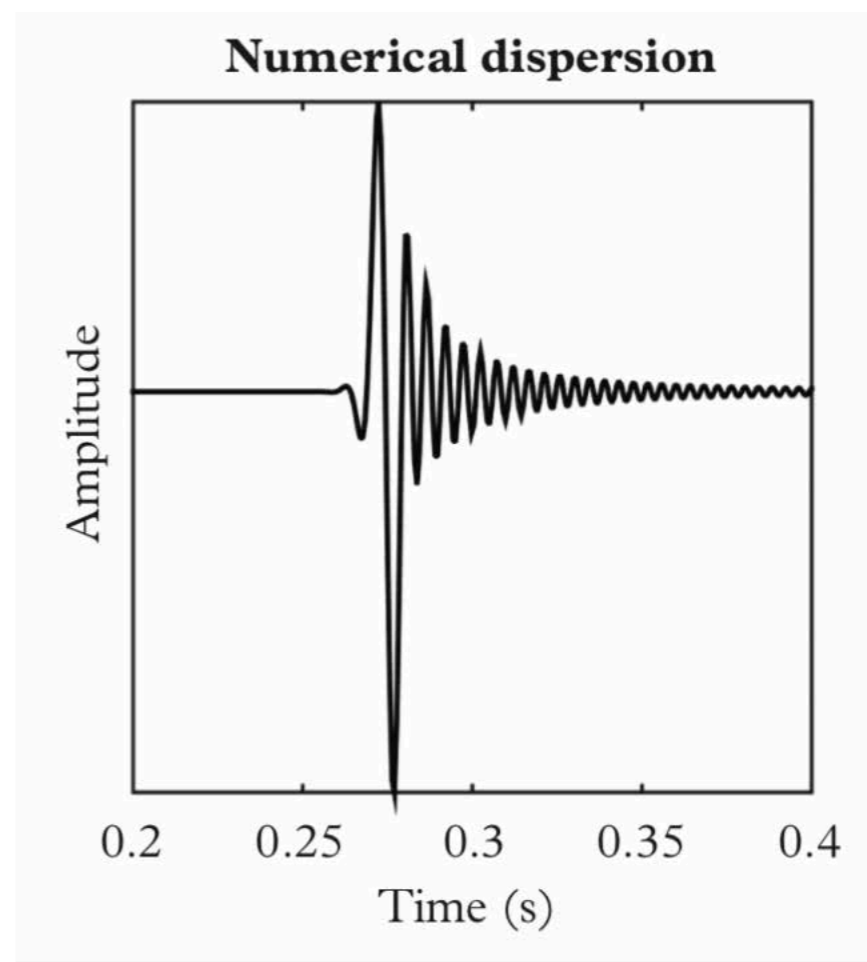
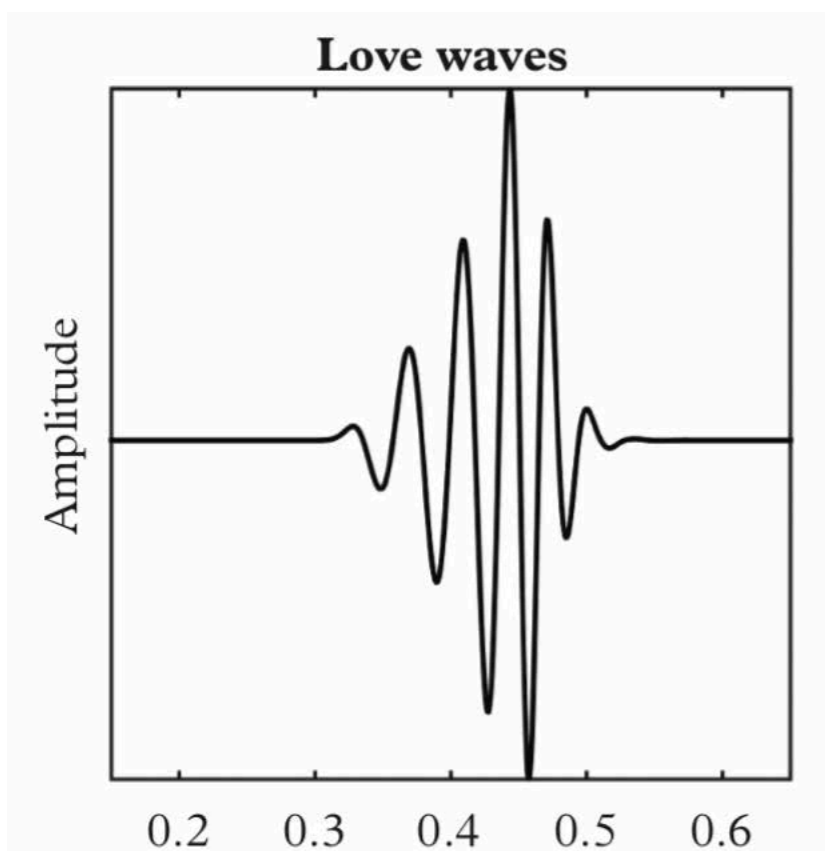
Figure 2.8-2: Fundamental mode Love wave group and phase velocities.



$$c(f) = \lambda(f) f \quad c(f) = \frac{\Delta}{t}$$



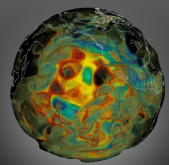
Numerical dispersion



velocity depends on wavenumber!

$$\omega = \frac{2}{dt} \sin^{-1} \left[c \frac{dt}{dx} \sin k \frac{dx}{2} \right]$$

$$c(k) = \frac{\omega}{k} = \frac{2}{k dt} \sin^{-1} \left[c \frac{dt}{dx} \sin k \frac{dx}{2} \right]$$



Weak methods (integral form)

Equation of motion:

$$\rho \partial_t^2 \mathbf{s} - \nabla \cdot \mathbf{T} = \mathbf{f}$$

Strong form:

$$\partial_t \mathbf{v} = \rho^{-1} (\nabla \cdot \mathbf{T} + \mathbf{f}) \quad \partial_t \mathbf{T} = \mathbf{c} : \nabla \mathbf{v}$$

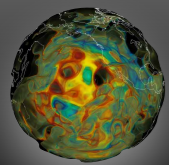
$$\mathbf{v} = \partial_t \mathbf{s}$$

Weak form:

$$\int_{\Omega} \rho \mathbf{w} \cdot \partial_t^2 \mathbf{s} d^3 \mathbf{x} = \int_{\Omega} \nabla \mathbf{w} : \mathbf{T} d^3 \mathbf{x} + \mathbf{M} : \nabla \mathbf{w}(\mathbf{x}_s) S(t)$$

Finite-fault rapture:

$$\mathbf{M} : \nabla \mathbf{w}(\mathbf{x}_s) S(t) \rightarrow \int_{S_s} \mathbf{m}(\mathbf{x}_s, t) : \nabla \mathbf{w}(\mathbf{x}_s) d^2 \mathbf{x}_s$$



Weak methods (integral form)

Equation of motion:

$$\rho \partial_t^2 \mathbf{s} - \nabla \cdot \mathbf{T} = \mathbf{f}$$

Str

- implicitly accounts for free-surface boundary condition

We

- fluid-solid coupling terms may be explicitly added

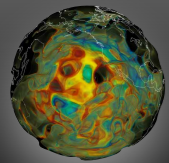
- source term explicitly integrated

- valid for any test vector

(t)

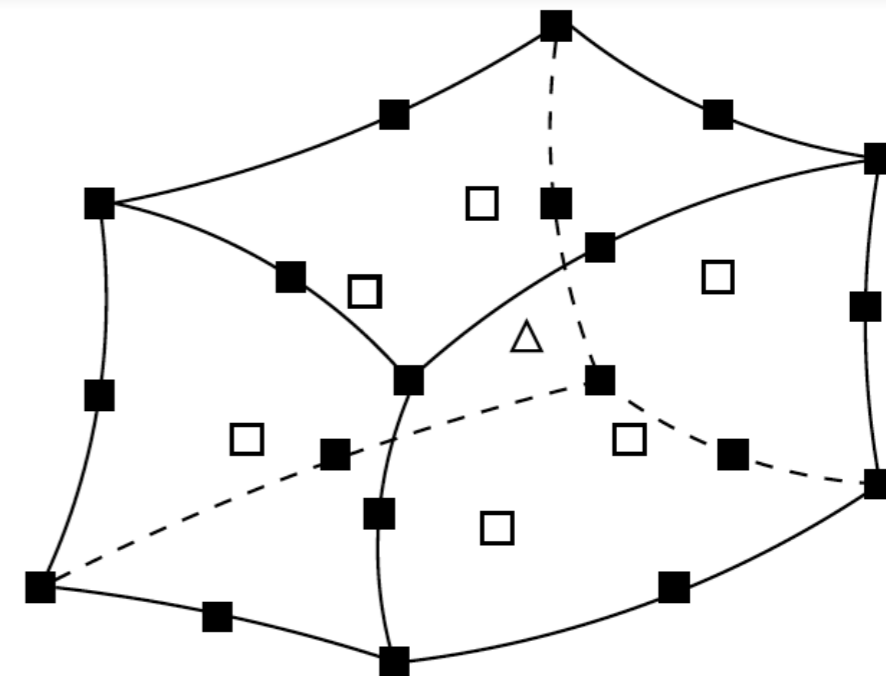
Finite-fault rapture:

$$\mathbf{M} : \nabla \mathbf{w}(\mathbf{x}_s) S(t) \rightarrow \int_{S_s} \mathbf{m}(\mathbf{x}_s, t) : \nabla \mathbf{w}(\mathbf{x}_s) d^2 \mathbf{x}_s$$

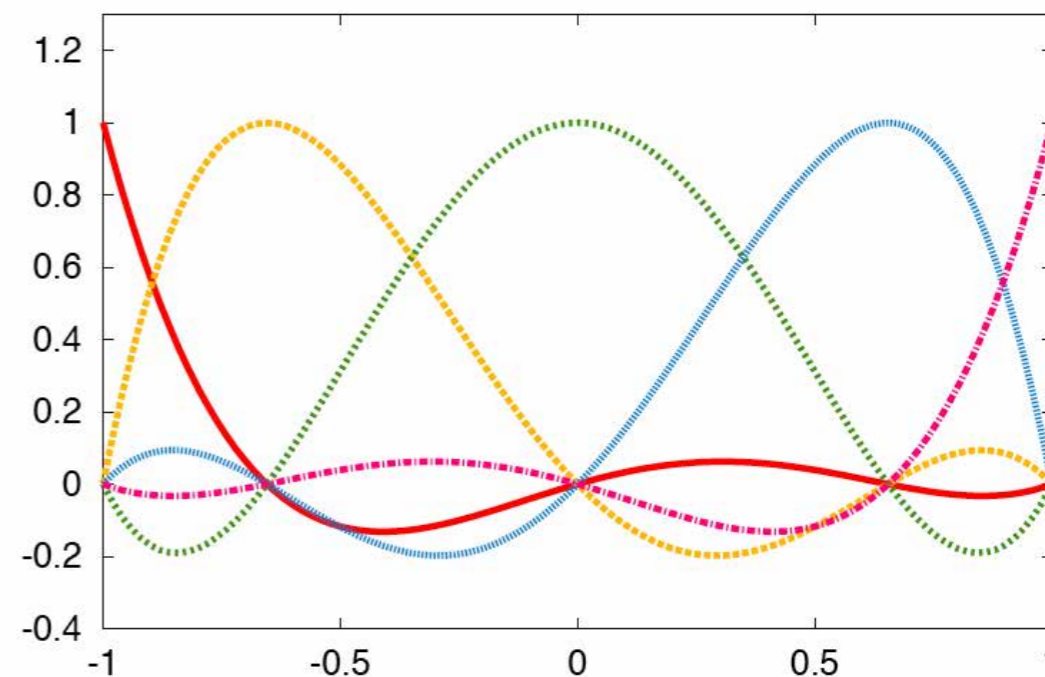


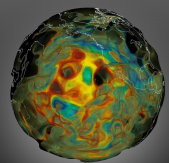
Spectral-element method

- hexahedral elements
- Gauss-Lobatto-Legendre quadrature
- finite elements with high-degree polynomial interpolation: accuracy of pseudo-spectral element methods with the flexibility of finite-element methods
- diagonal mass matrix makes parallel implementation very efficient
- mesh honours major discontinuities
- explicit time-marching scheme



5 Degree 4 Lagrange polynomials



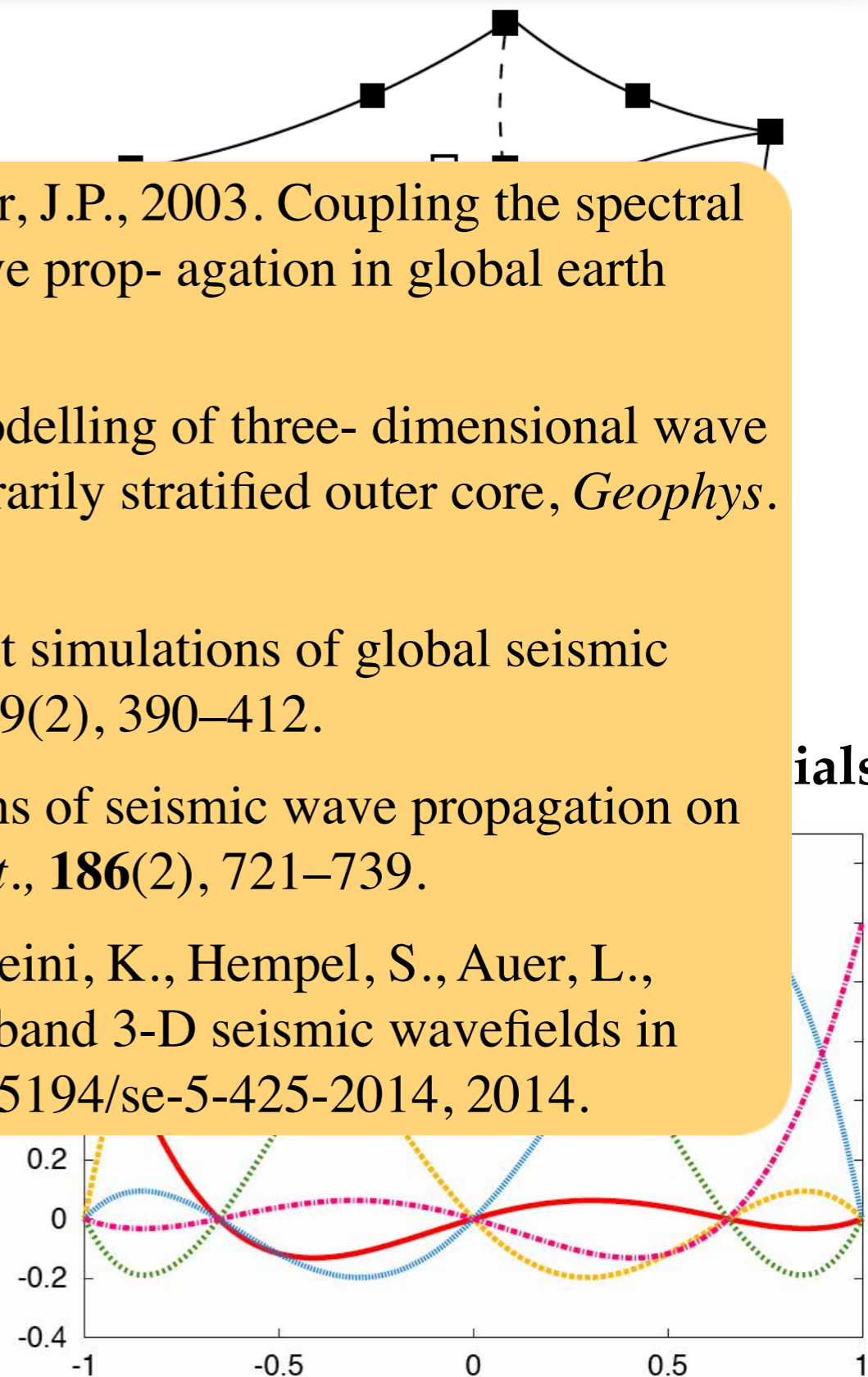


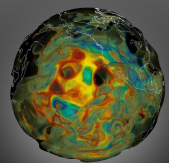
Spectral-element method

- hexahedral elements

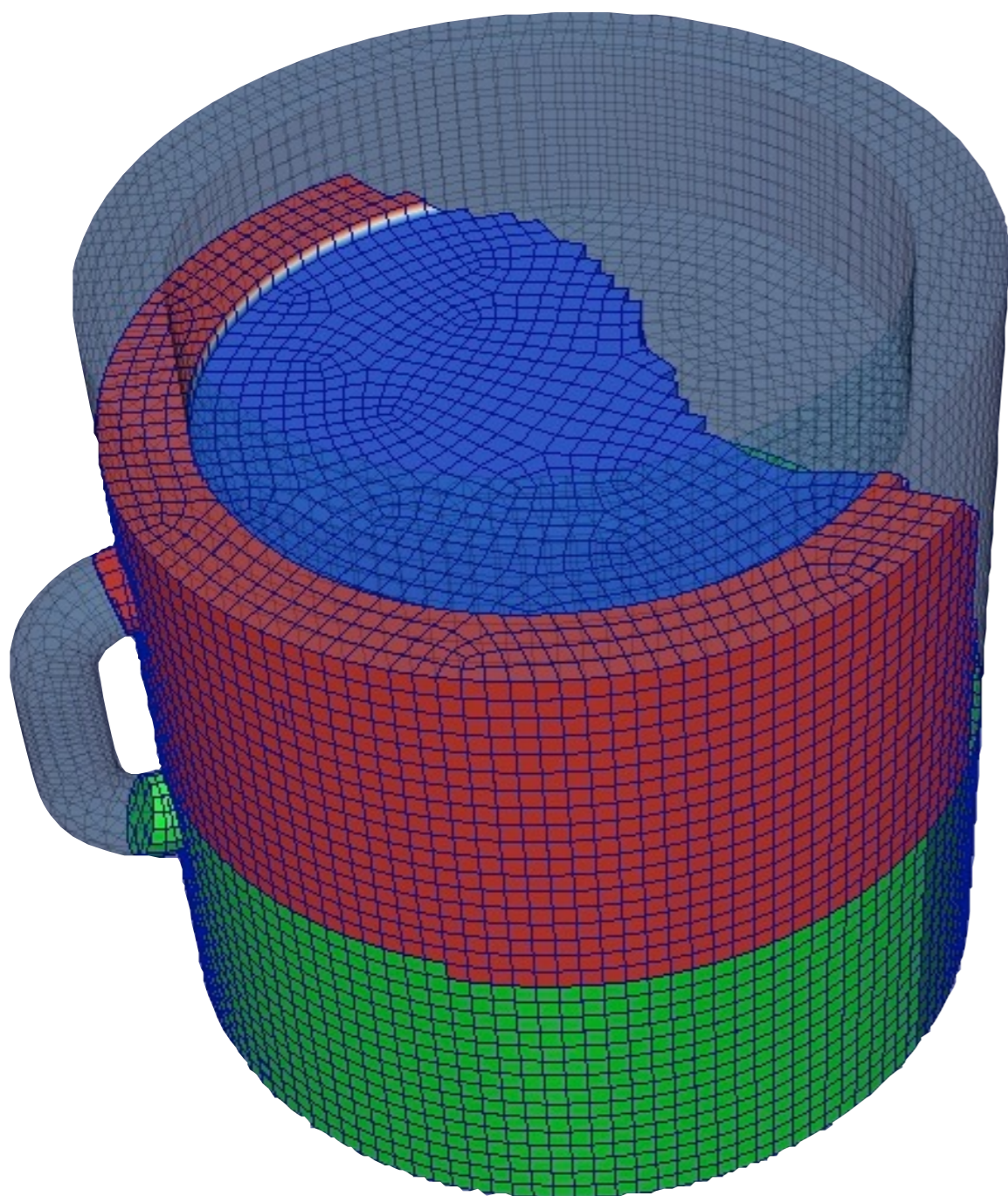
- Capdeville, Y., Chaljub, E., Vilotte, J.P. & Montagner, J.P., 2003. Coupling the spectral element method with a modal solution for elastic wave propagation in global earth models, *Geophys. J. Int.*, 152, 34–67.
- Chaljub, E. & Valette, B., 2004. Spectral element modelling of three-dimensional wave propagation in a self-gravitating Earth with an arbitrarily stratified outer core, *Geophys. J. Int.*, 158, 131–141.
- Komatitsch, D. & Tromp, J., 2002a. Spectral-element simulations of global seismic wave propagation – I. Validation, *Geophys. J. Int.*, 149(2), 390–412.
- Peter, D. *et al.*, 2011. Forward and adjoint simulations of seismic wave propagation on fully unstructured hexahedral meshes, *Geophys. J. Int.*, **186**(2), 721–739.
- Nissen-Meyer, T., van Driel, M., Stähler, S. C., Hosseini, K., Hempel, S., Auer, L., Colombi, A., and Fournier, A., 2014. AxiSEM: broadband 3-D seismic wavefields in axisymmetric media, *Solid Earth*, 5, 425-445, doi:10.5194/se-5-425-2014, 2014.

- mesh honours major discontinuities
- explicit time-marching scheme

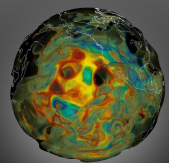




Coffee cup simulation by Daniel

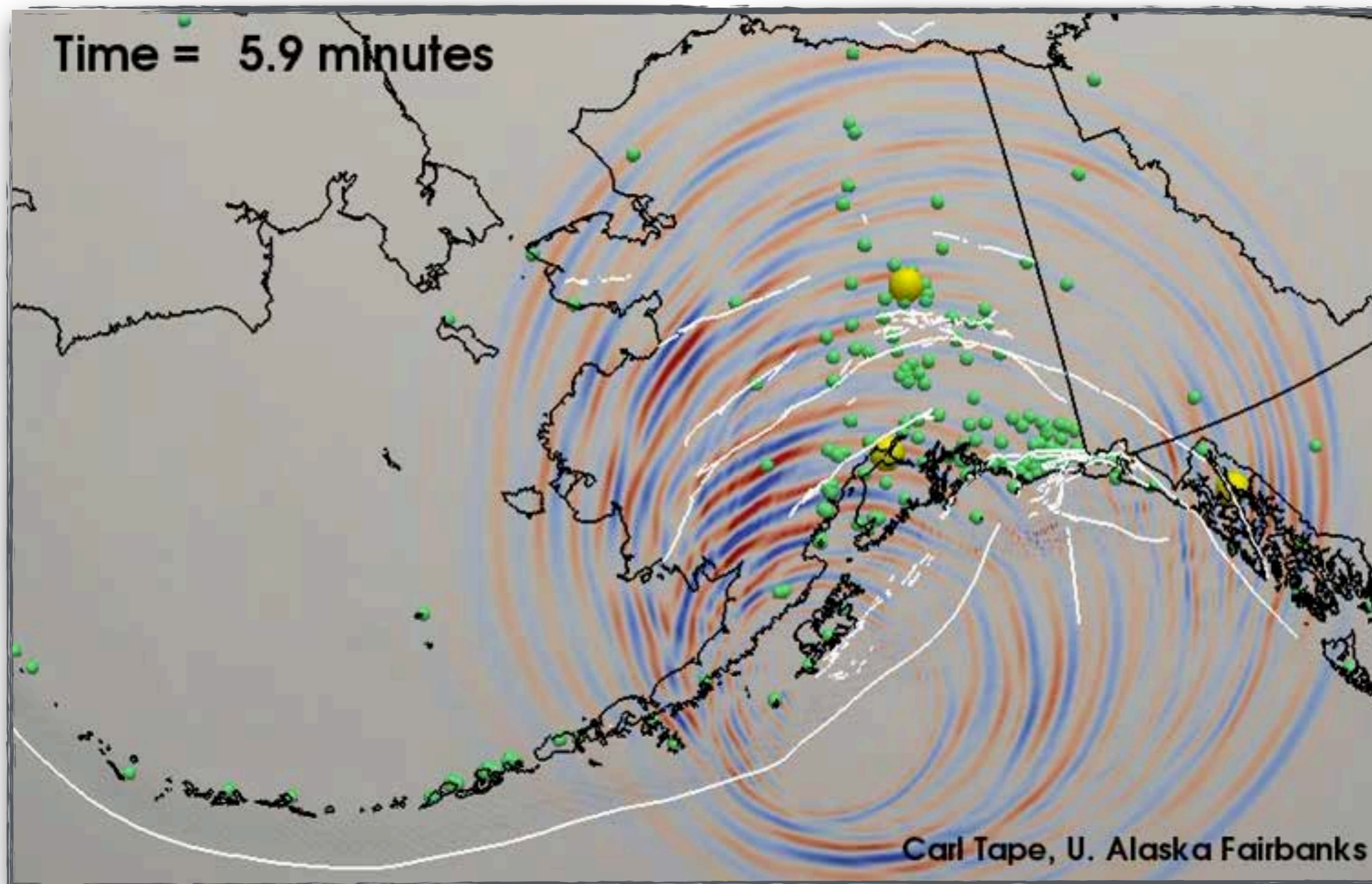
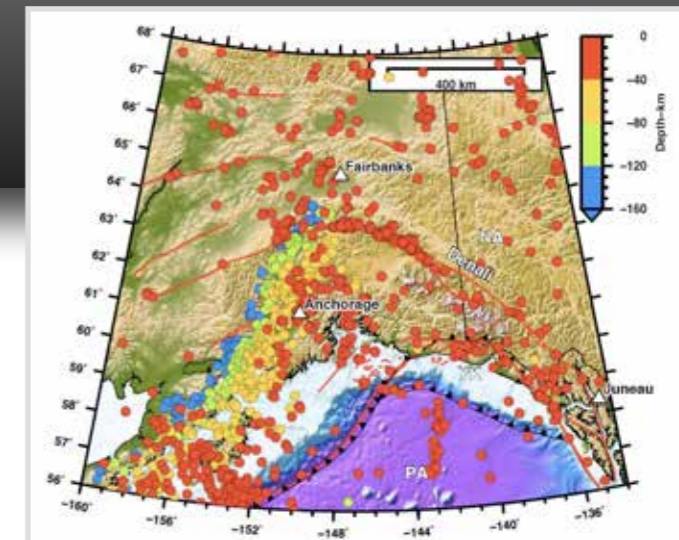


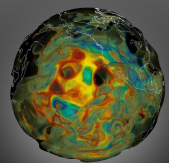
Daniel Peter



Finite-source effects

1964 Alaska Earthquake





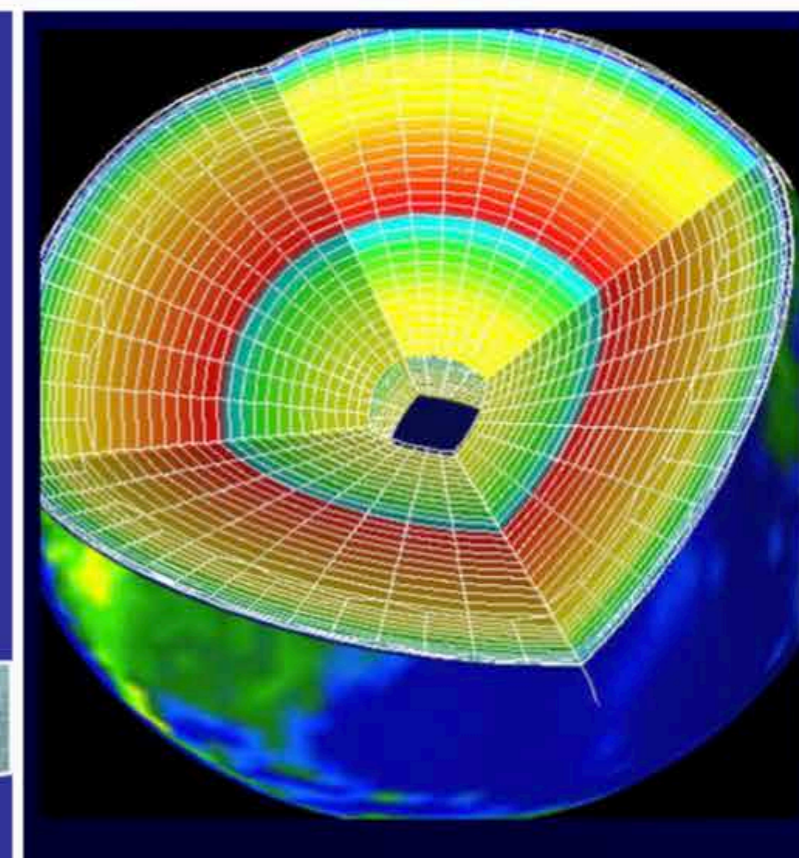
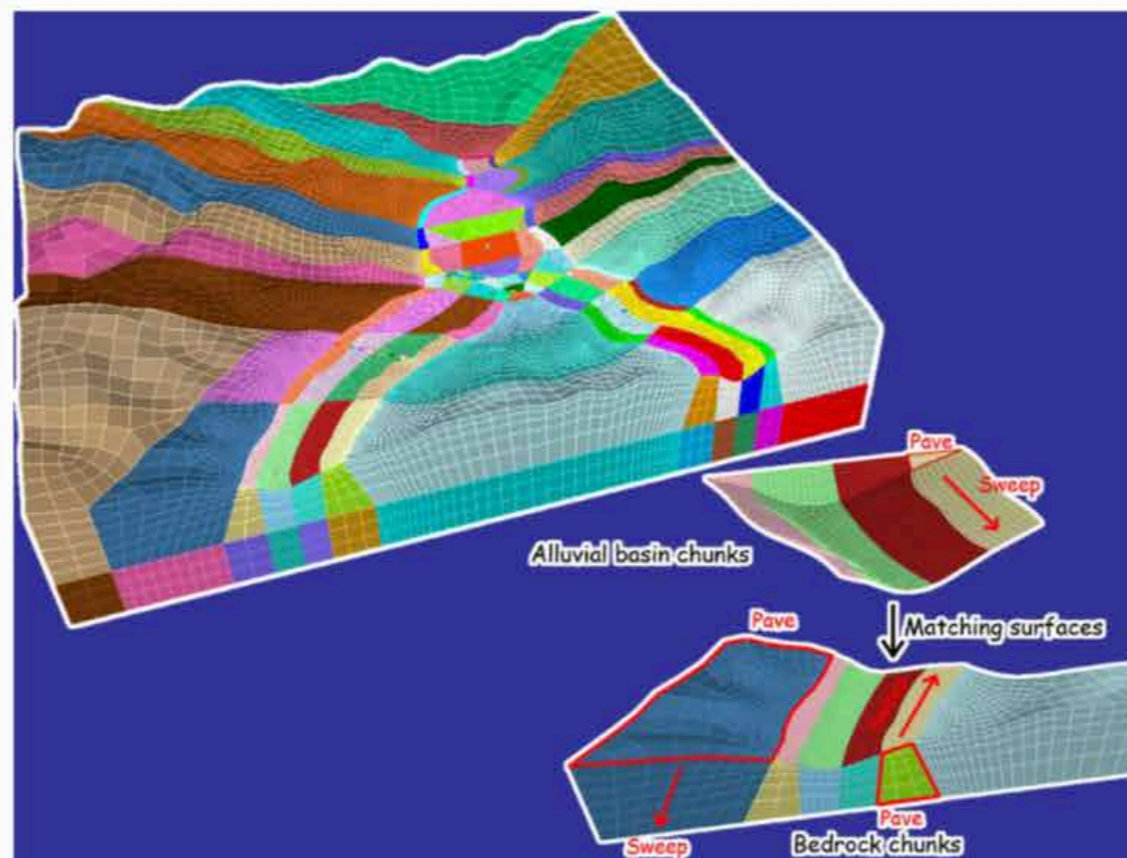
Meshing the Earth

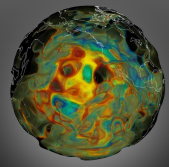
Human time	Simulation workflow	CPU time
15%	Design	0%
80%	Geometry & Meshing	10%
5%	Solver	90%

Source: E. Casarotti, pers. comm.

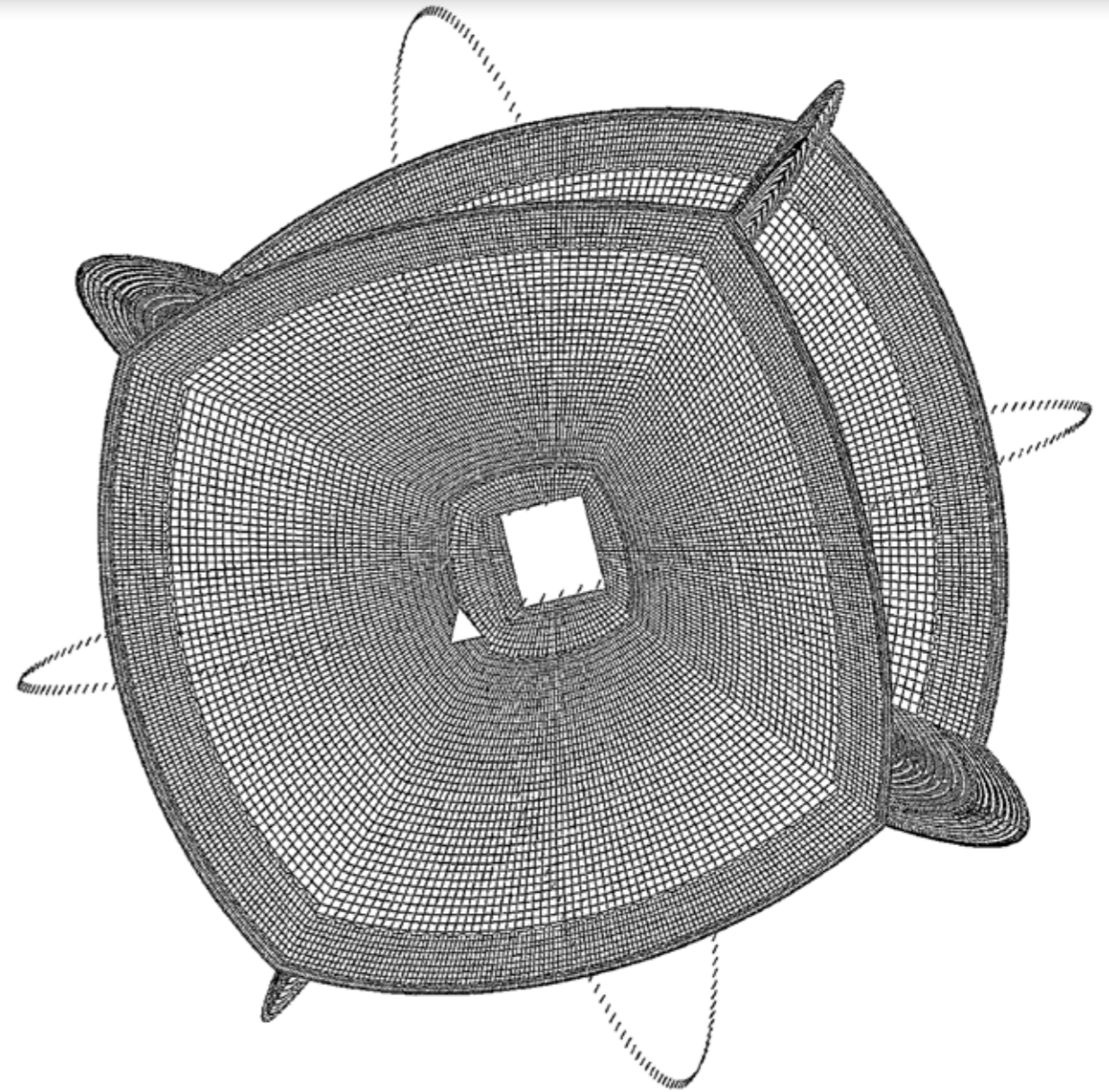
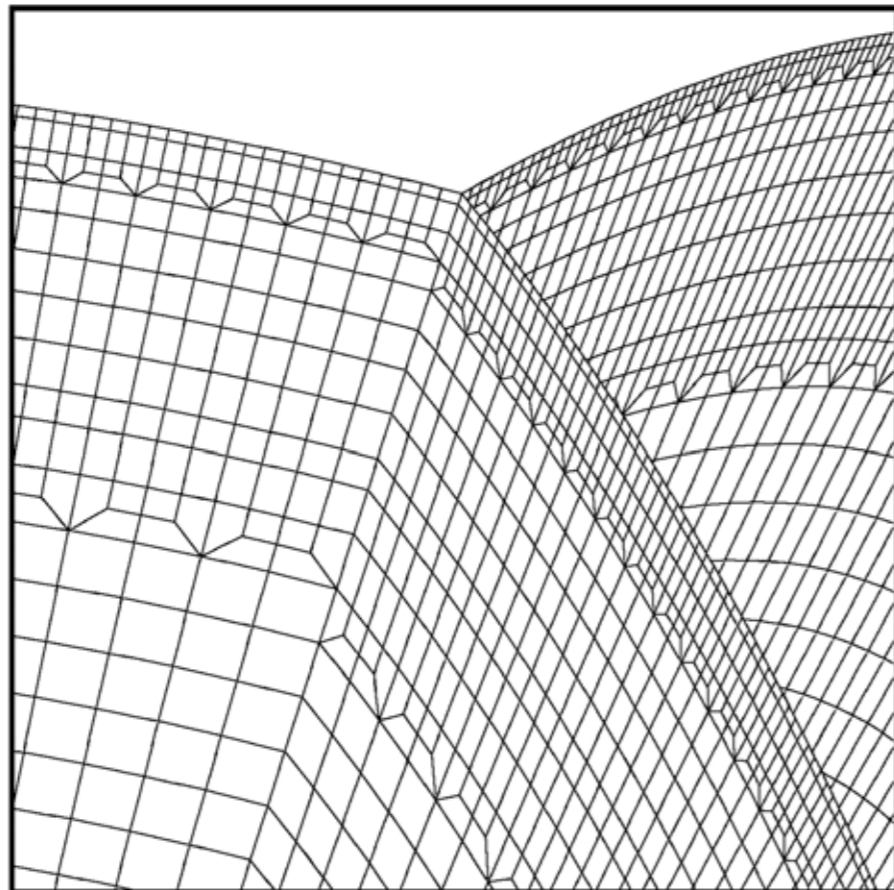
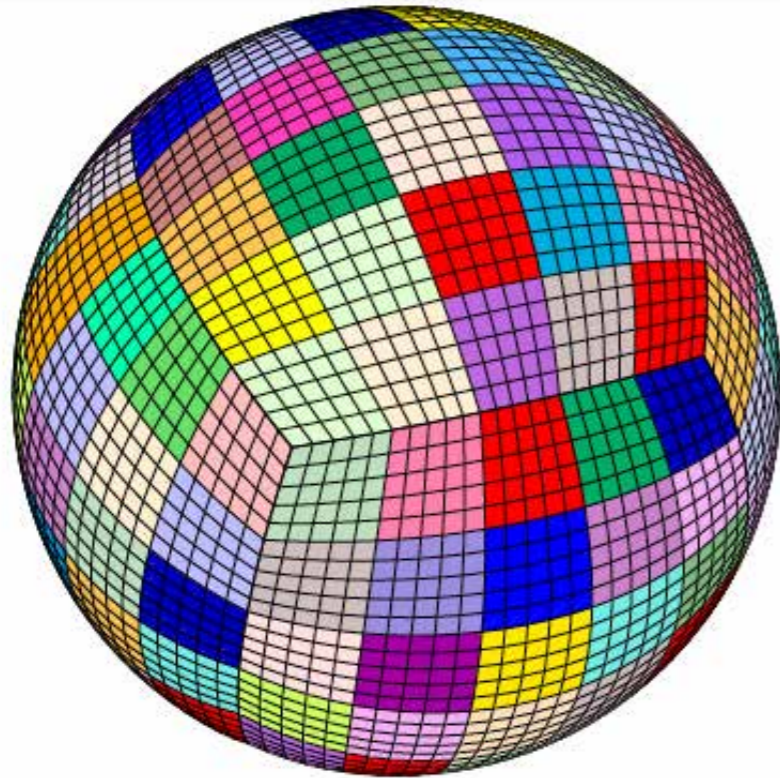
SPECFEM3D & SW4: Regular meshes

Partitioning: balance between computational load and communication.



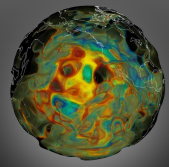


Global mesh

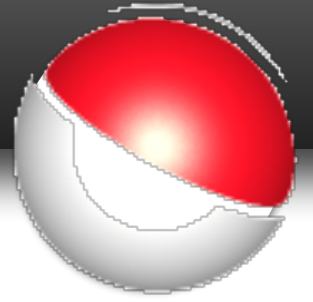


D. Komatitsch & J. Tromp, 2002, GJI

**PREM (Preliminary Reference Model)
(A. Dziewonski & D. Anderson, 1981)**



3D global wave simulations



February 16, 2018, Oaxaca, Mexico

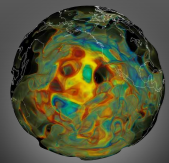
Mw=7.2, depth=19 km

global.shakemovie.princeton.edu/

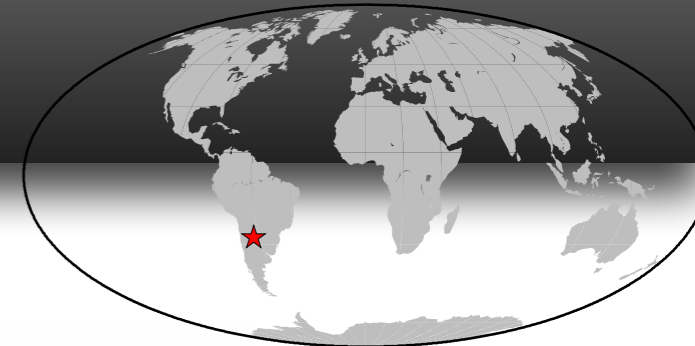


PRINCETON
UNIVERSITY

0:42:52

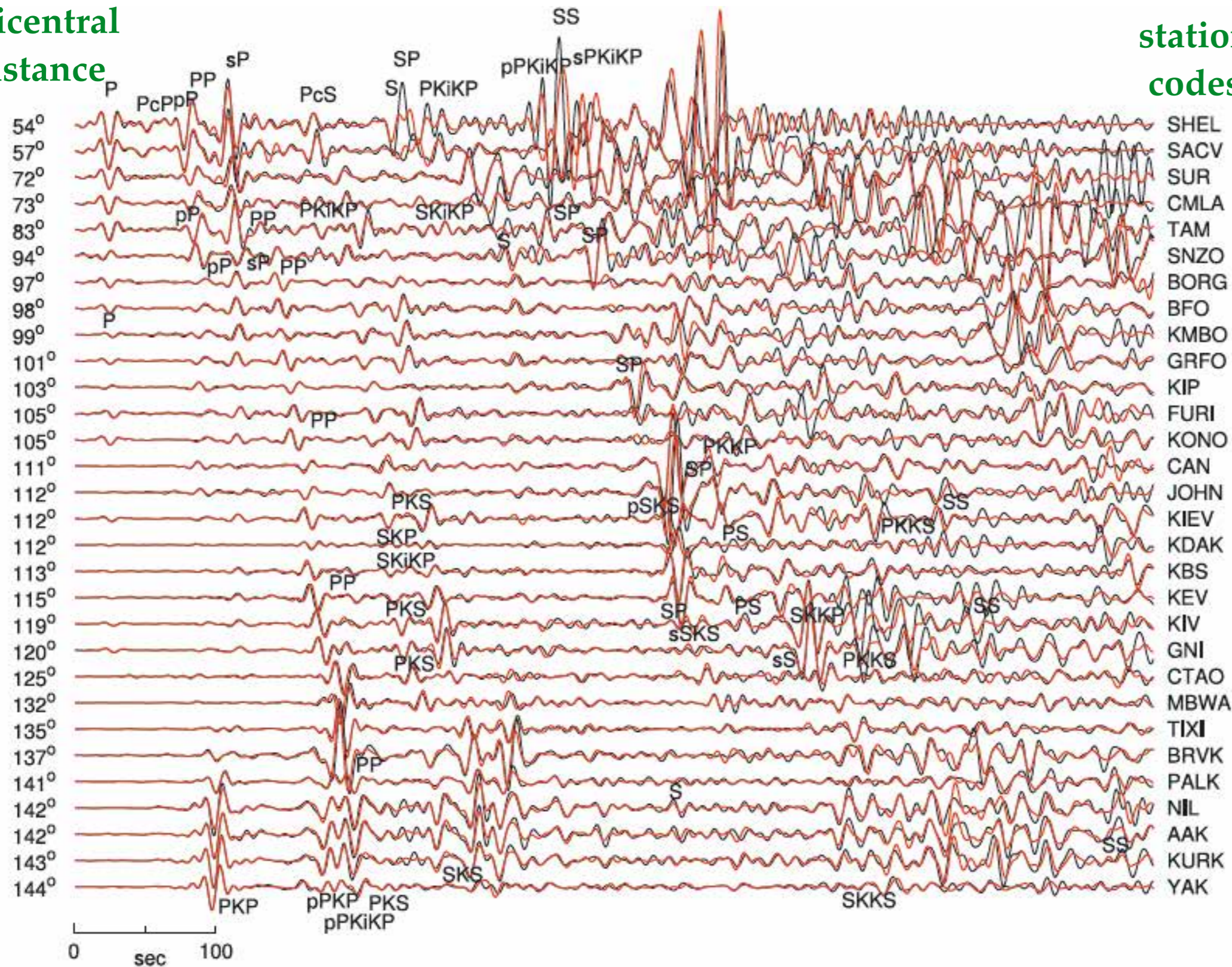


Observed vs. 3D seismograms



epicentral
distance

station
codes

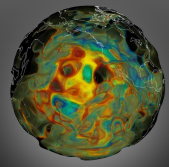


Sept. 3, 2008,
Argentina
Earthquake
 $M_w = 6.3$
depth = 571 km

SPECFEM3D_GLOBE

3D Earth model:
S362ANI + Crust2.0
(Kustowski et al. 2007 &
Bassin et al. 2000)

17 - 60 s

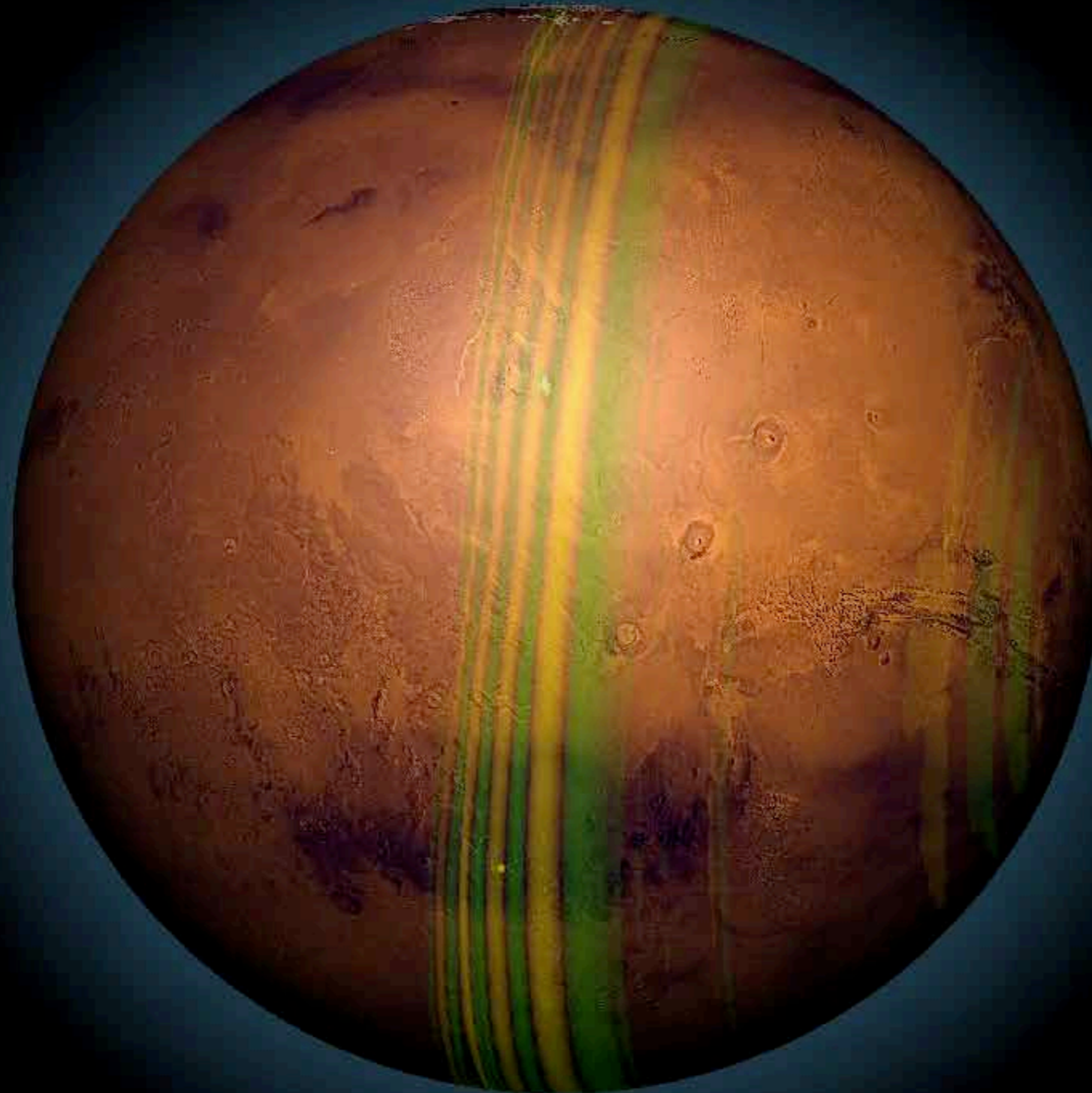


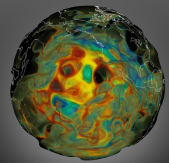
Seismic wave simulations on Mars



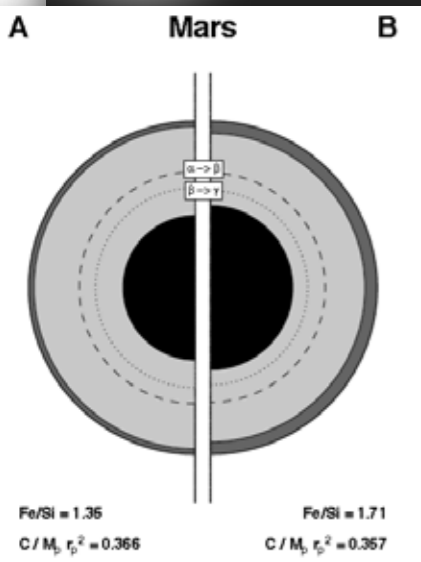
Daniel Peter & Ebru Bozdog

0:22:29



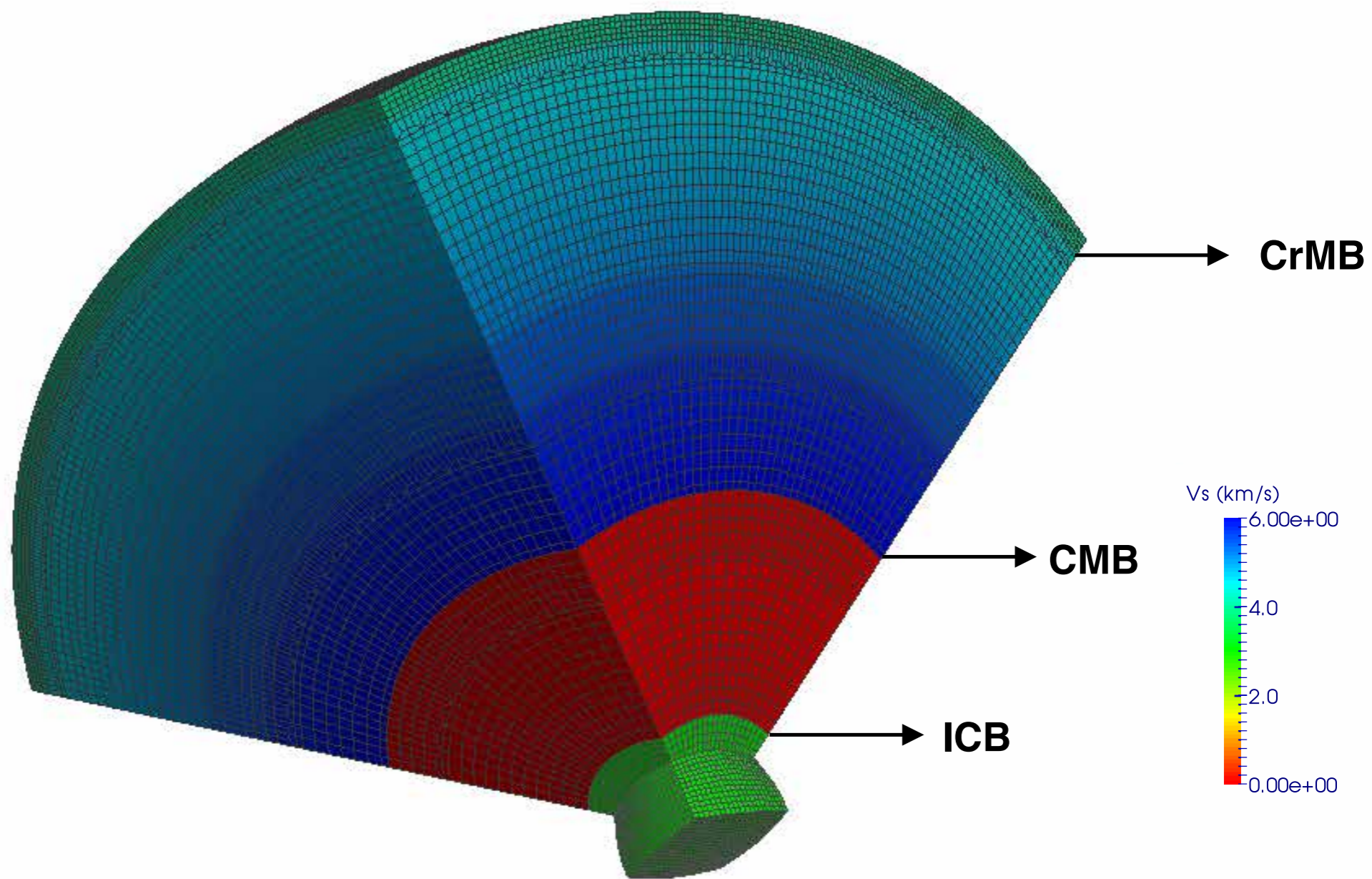
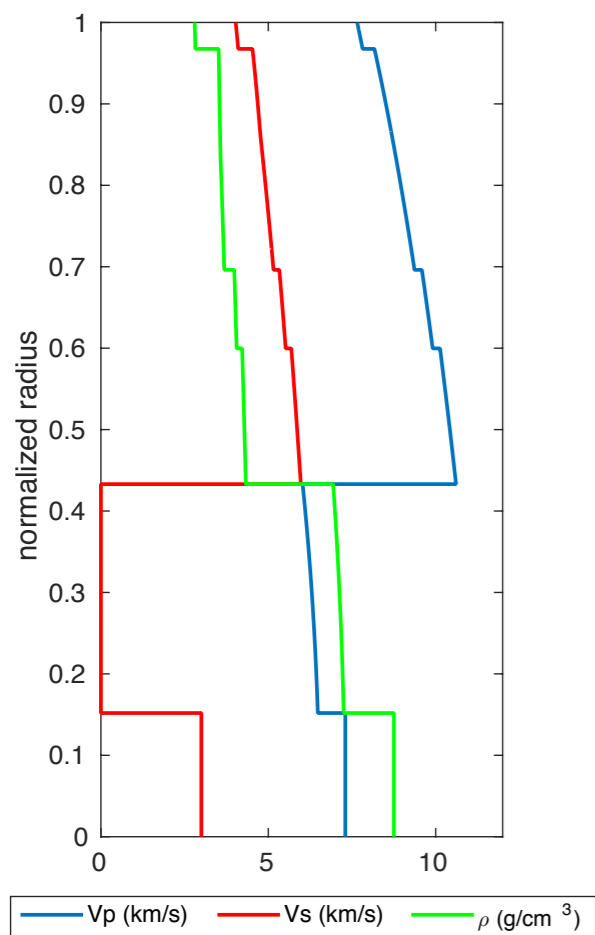


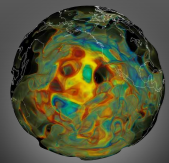
Seismic wave simulations on Mars



Sohl & Spohn (1997)
Model A

E. Bozdag, Y. Ruan et al. 2017, Space Science Rev.

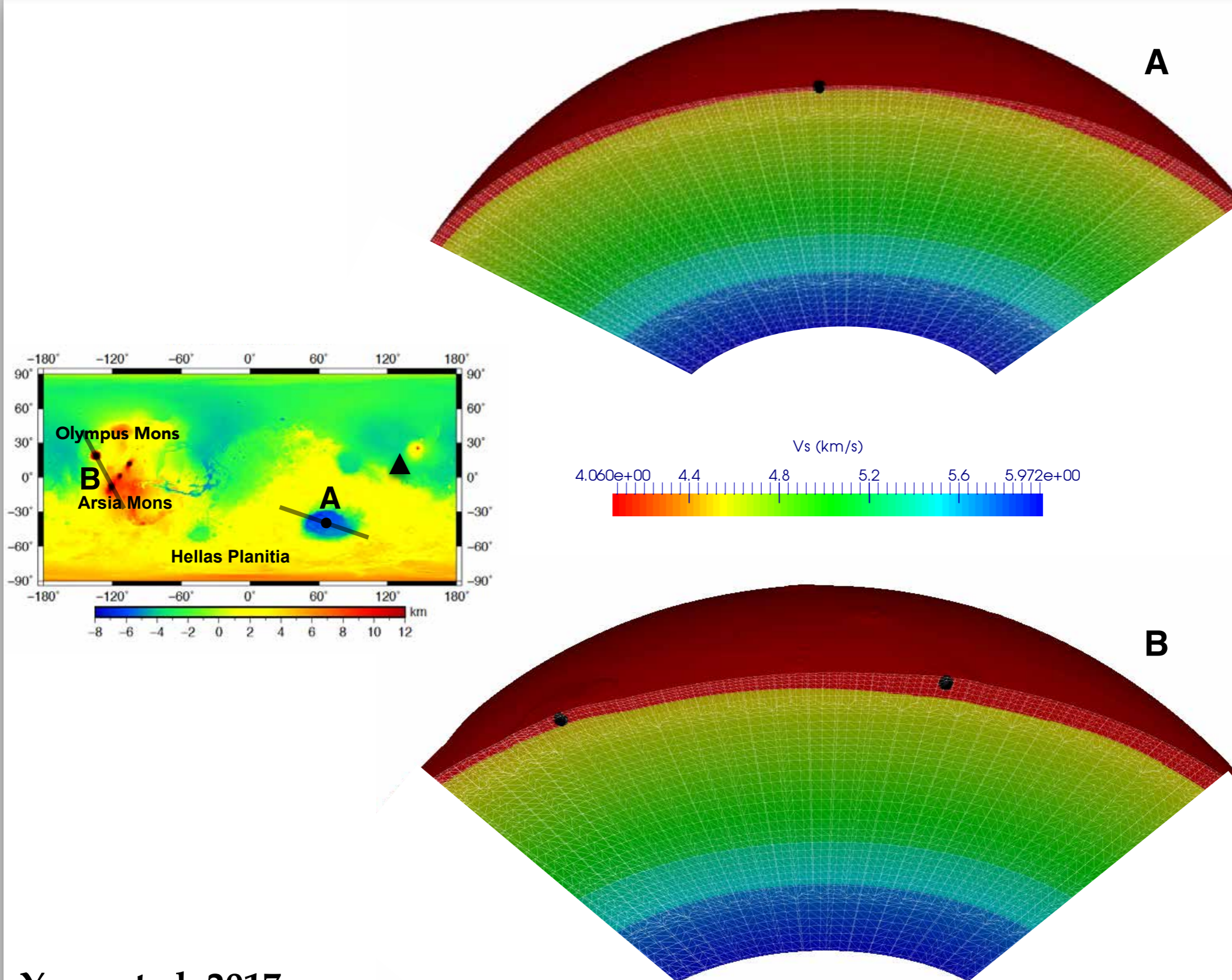


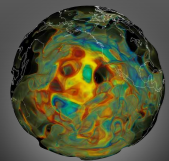


3D crustal mesh



InSight

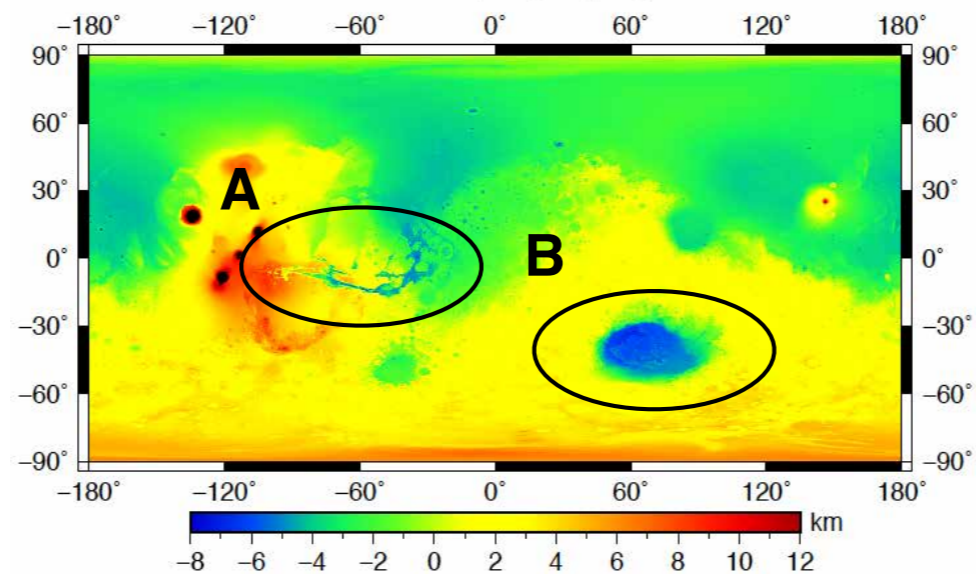




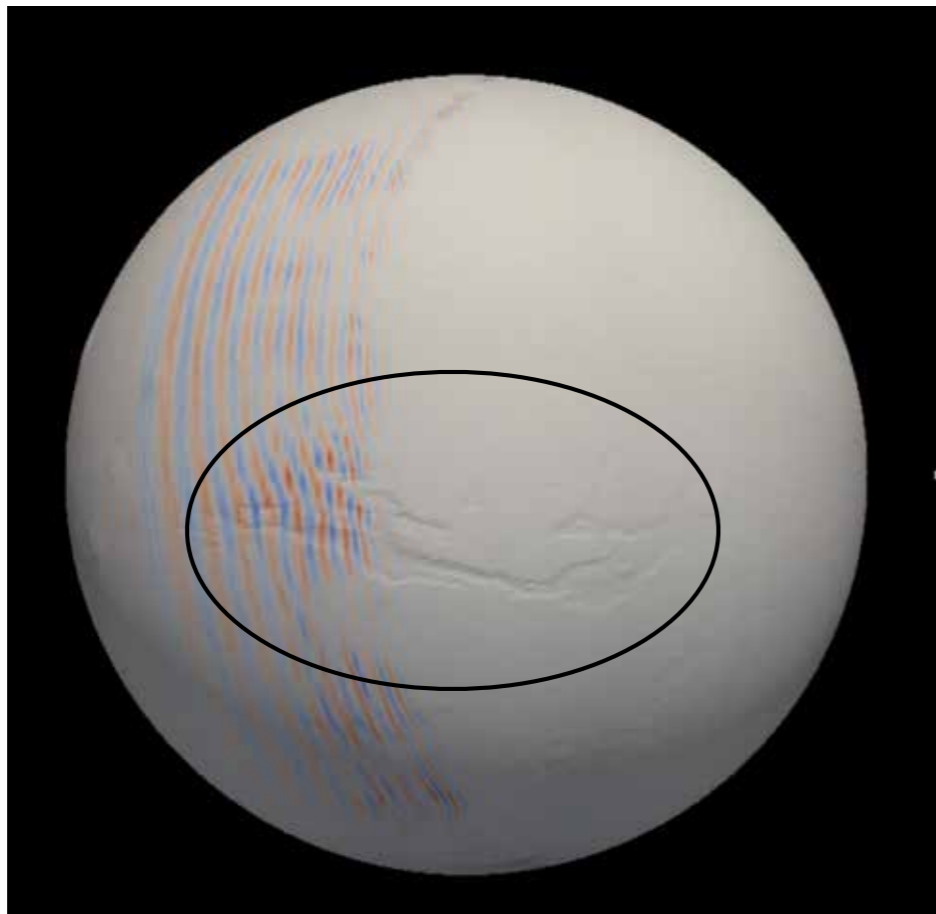
3D crustal effects



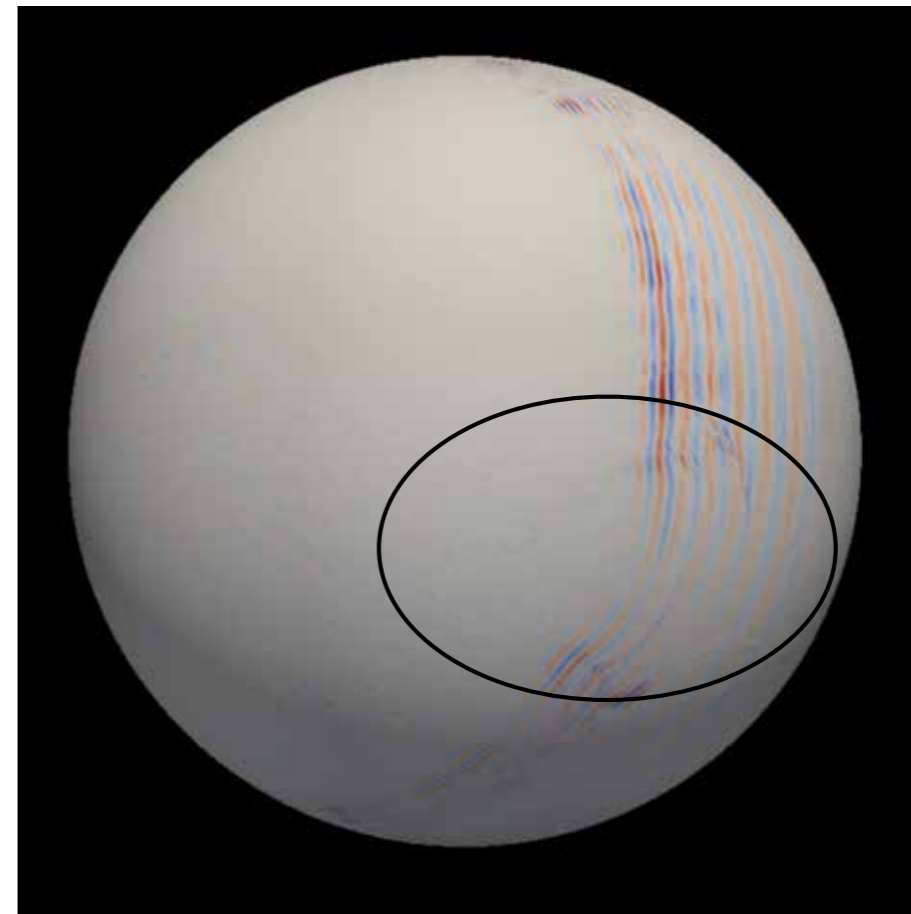
InSight

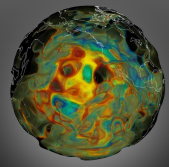


A



B

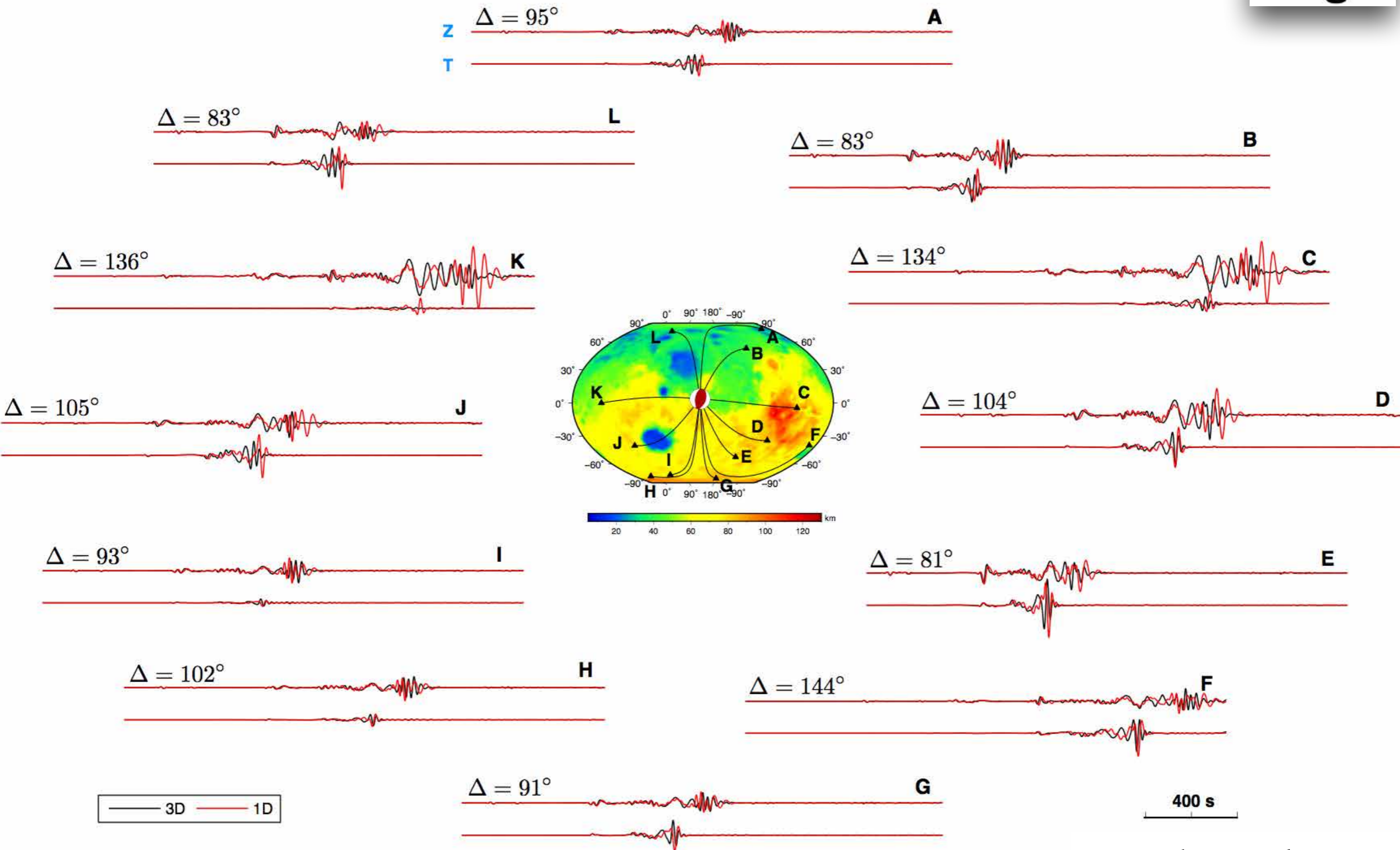


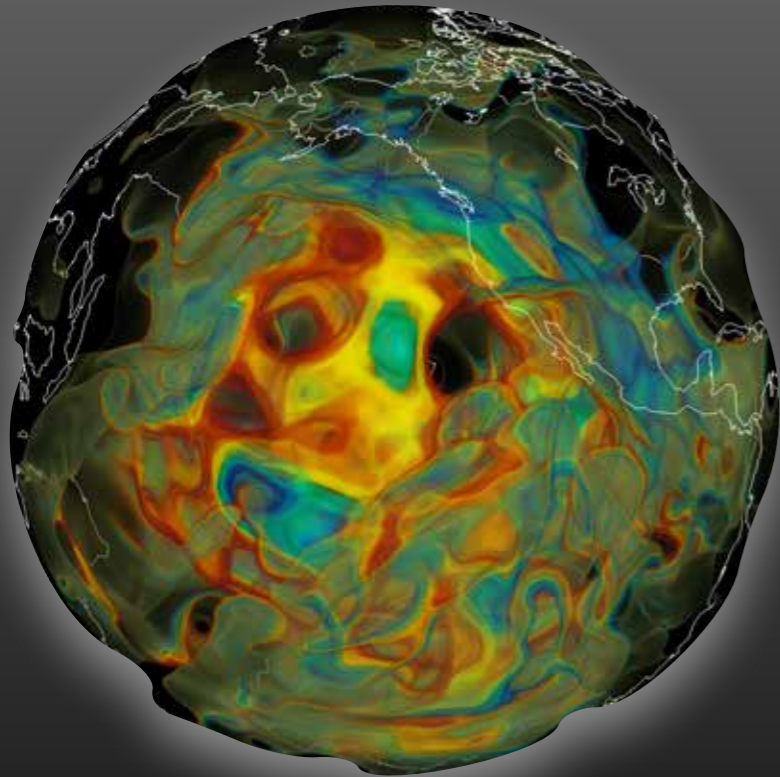


3D vs. 1D Martian seismograms

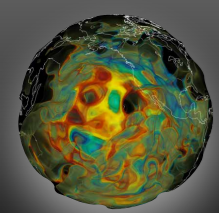


InSight





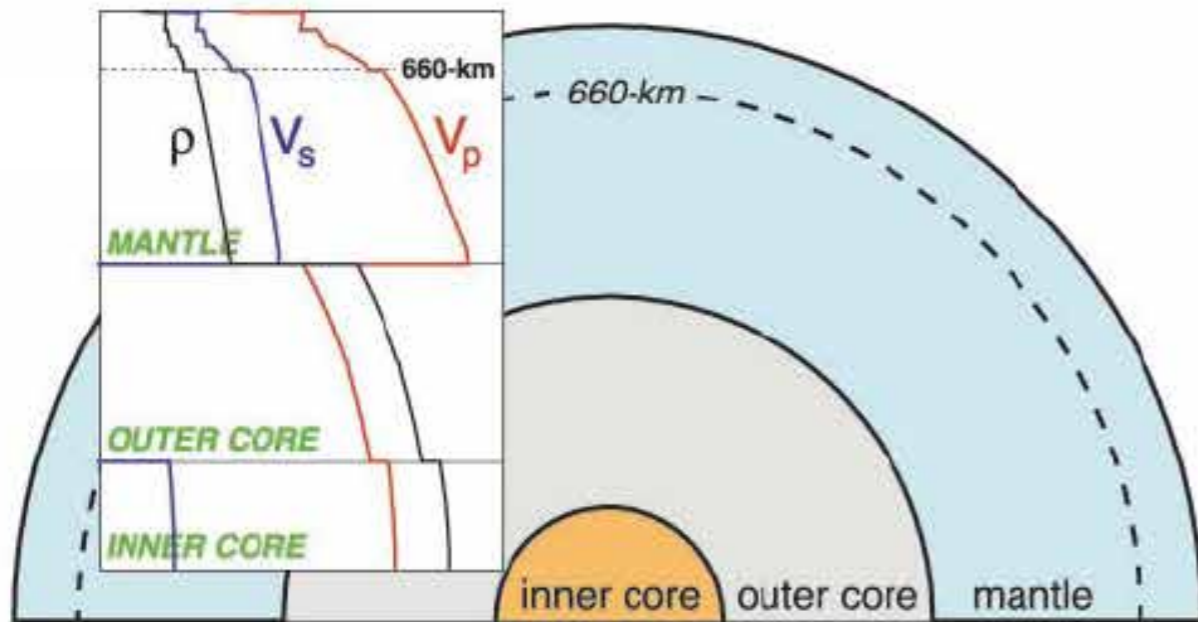
3D wave simulations in seismic tomography



Classical global seismic tomography

PREM

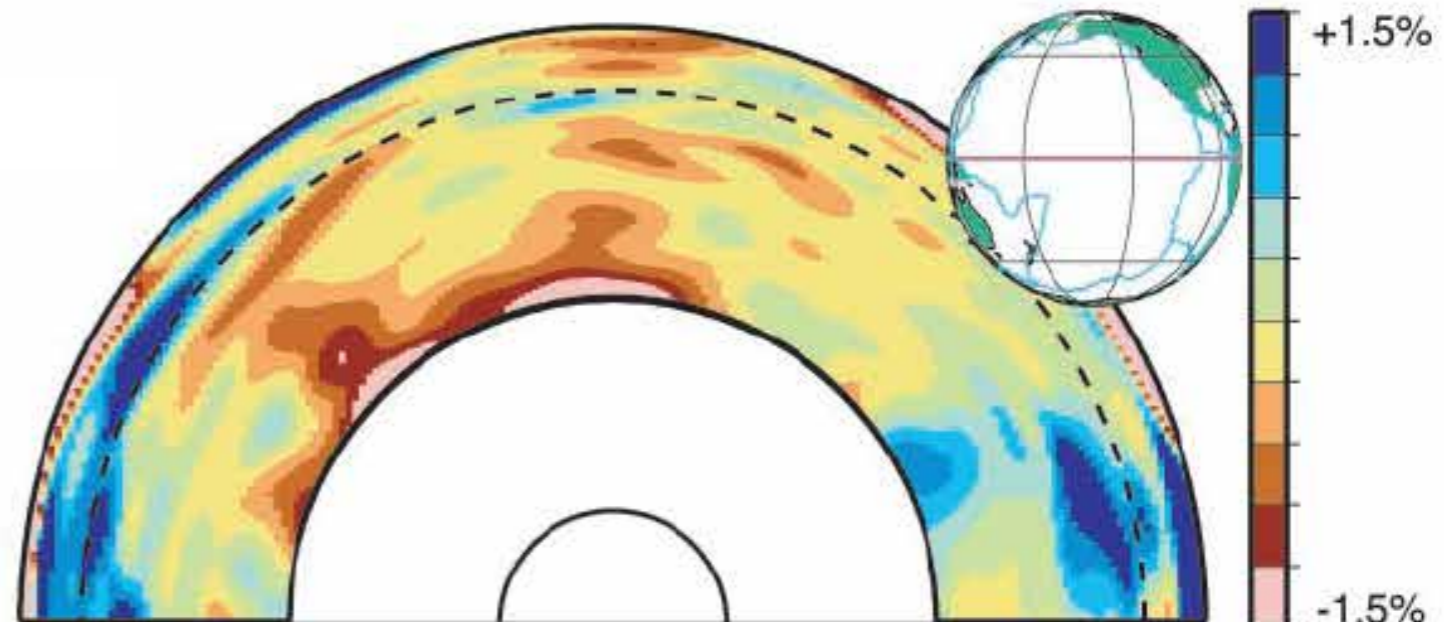
(Dziewonski & Anderson 1981)



$V(r)$
1D reference velocity model

S20RTS

(Ritsema et al. 1999)

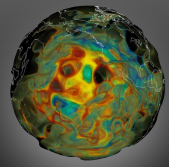


$\delta V(r, \theta, \phi)$
3D velocity structure

seismic
modeling

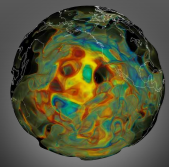
from D. Komatitsch, J. Ritsema & J. Tromp, 2002, Science

- Ray theory, finite-freq. theory based on 1D background models, 3D synthetics with asymptotic kernels
- Data coverage: Combination of different data types
- Crustal corrections



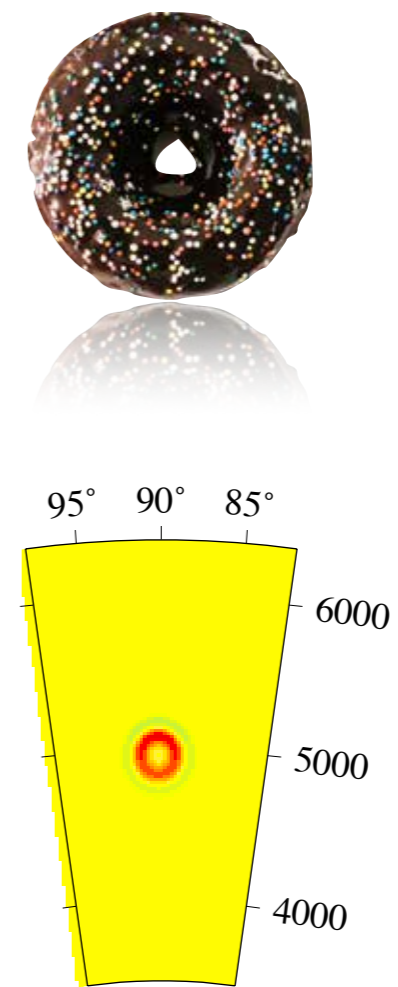
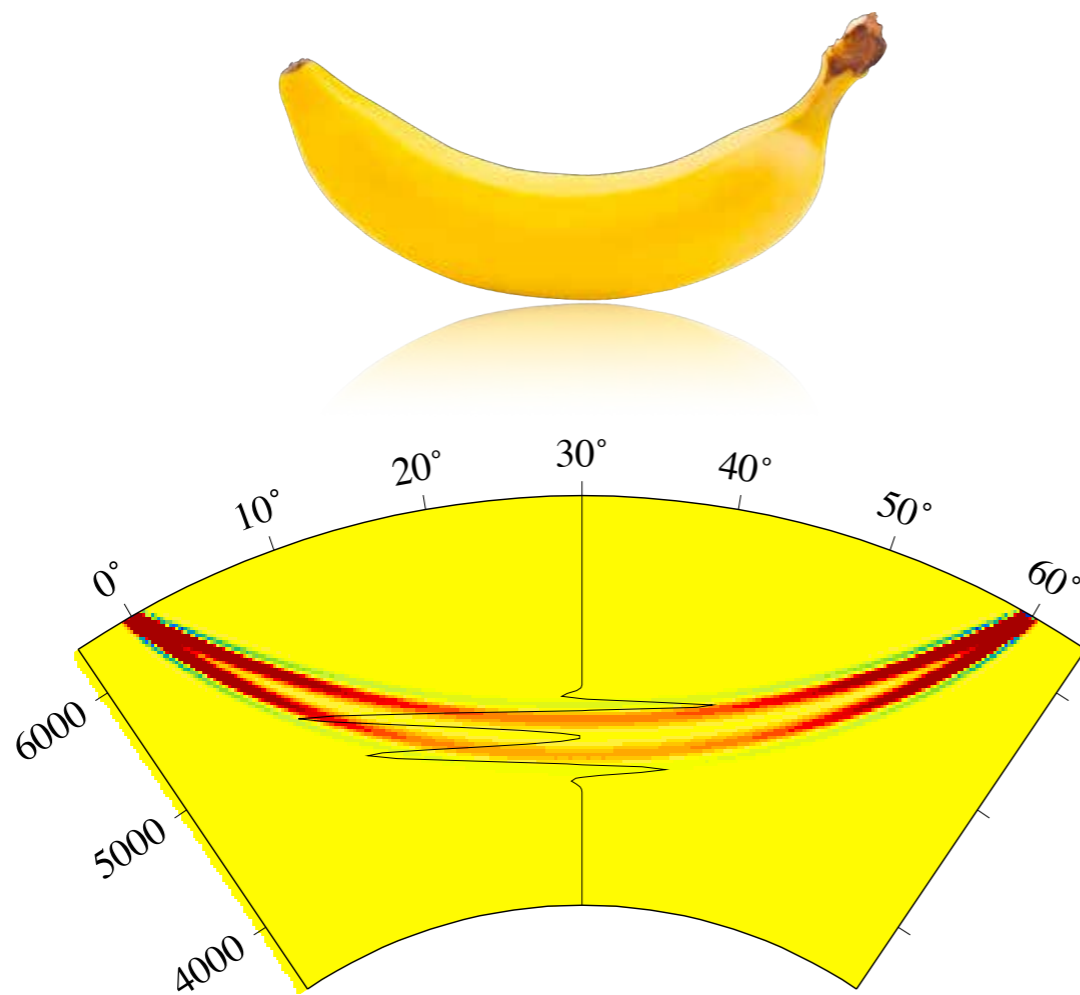
Challenges in global tomography

- **We are at a stage where details matter!**
 - *Avoid approximations as much as possible*

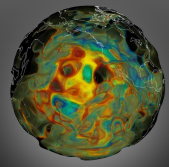


Challenges in global tomography

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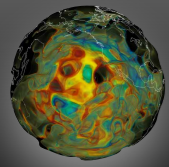


from Hung, Dahlen & Nolet, 2000, GJI



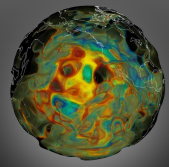
Challenges in global tomography

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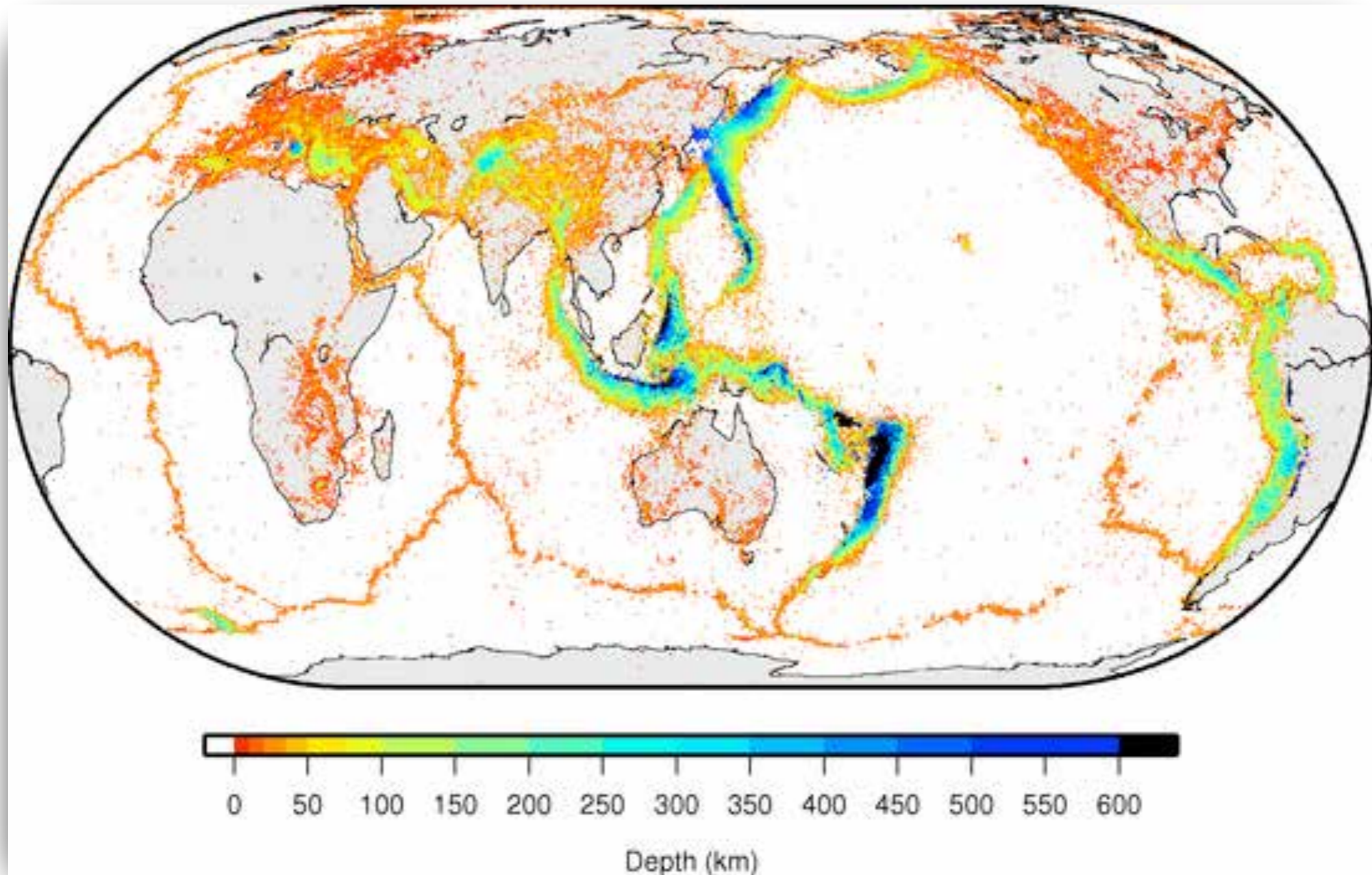
Challenges in global tomography

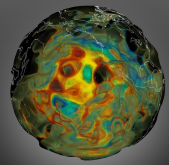
- **We are at a stage where details matter!**
 - *Avoid approximations as much as possible*
 - **Seismograms & Fréchet sensitivity kernels need to be calculated numerically in a realistic 3D Earth model**
- **Data coverage**



Earthquakes

ISC locations since 1960



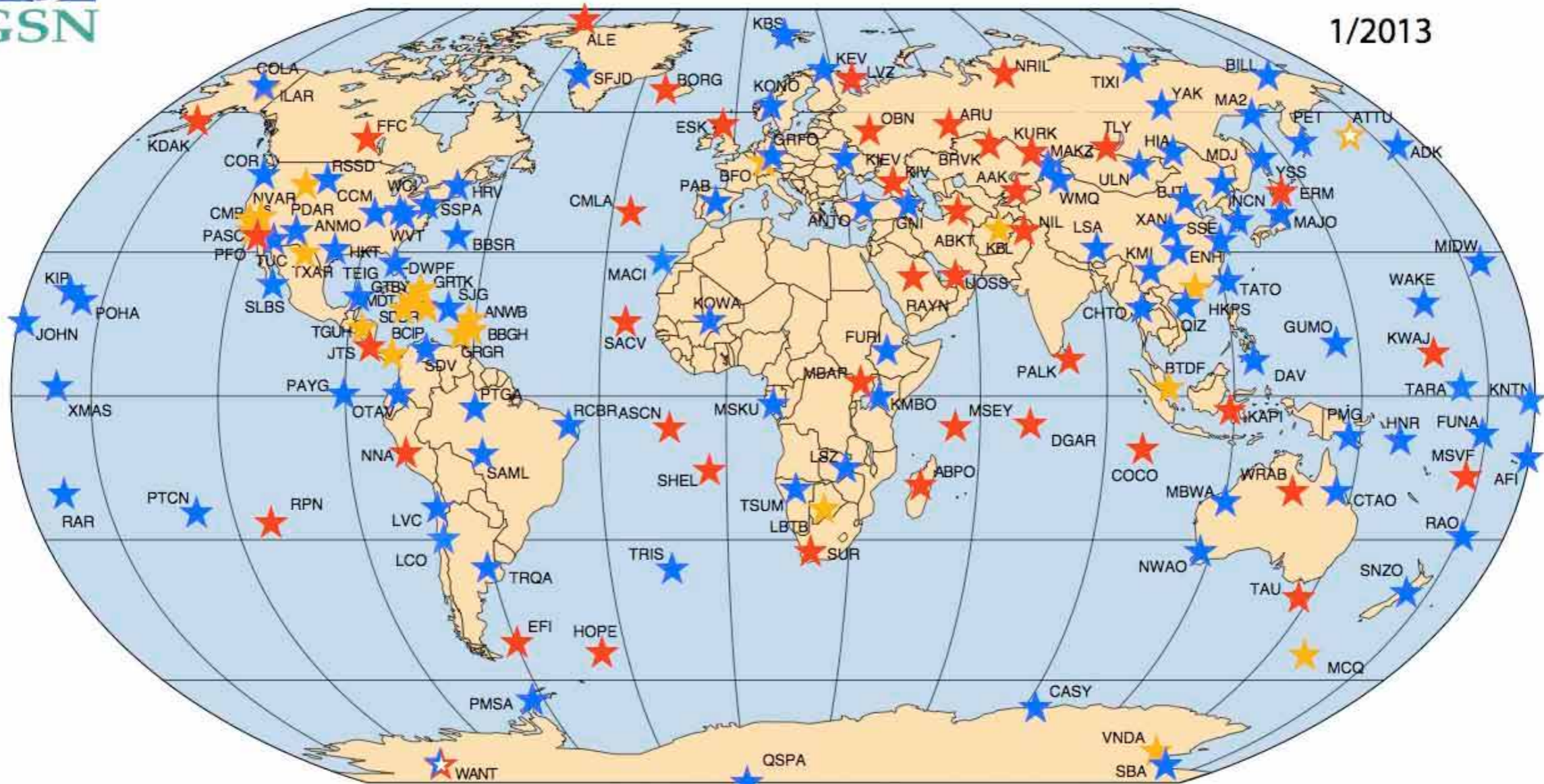


Seismic stations

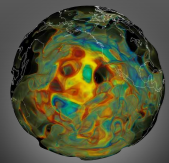


GLOBAL SEISMOGRAPHIC NETWORK

1/2013



- ★ IRIS / IDA Stations
- ★ IRIS / USGS Stations
- ★ Affiliate Stations
- ★ Planned Stations

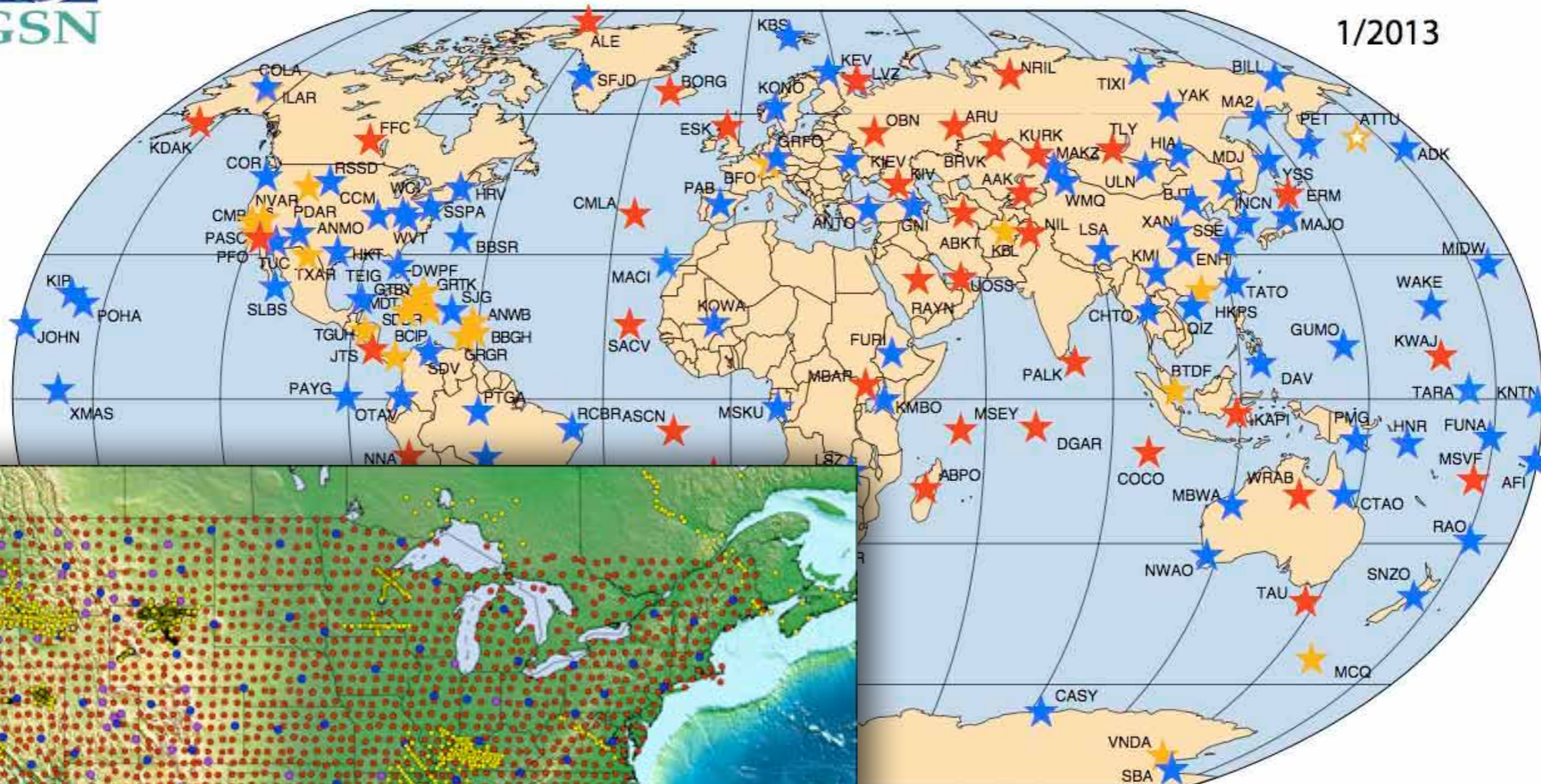


Seismic stations



GLOBAL SEISMOGRAPHIC NETWORK

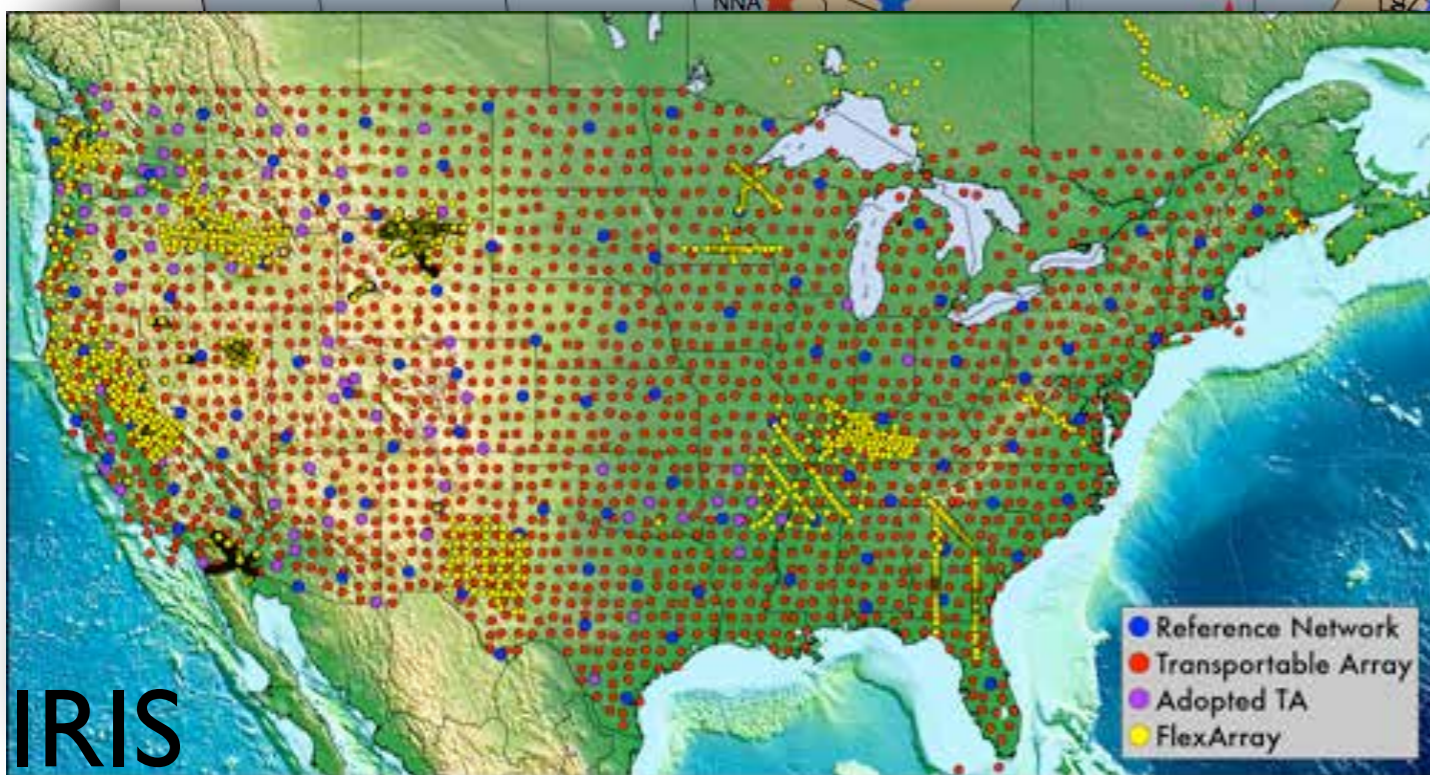
1/2013



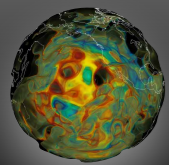
Stations



Affiliate Stations



IRIS



Seismic stations



GLOBAL SEISMIC NETWORK

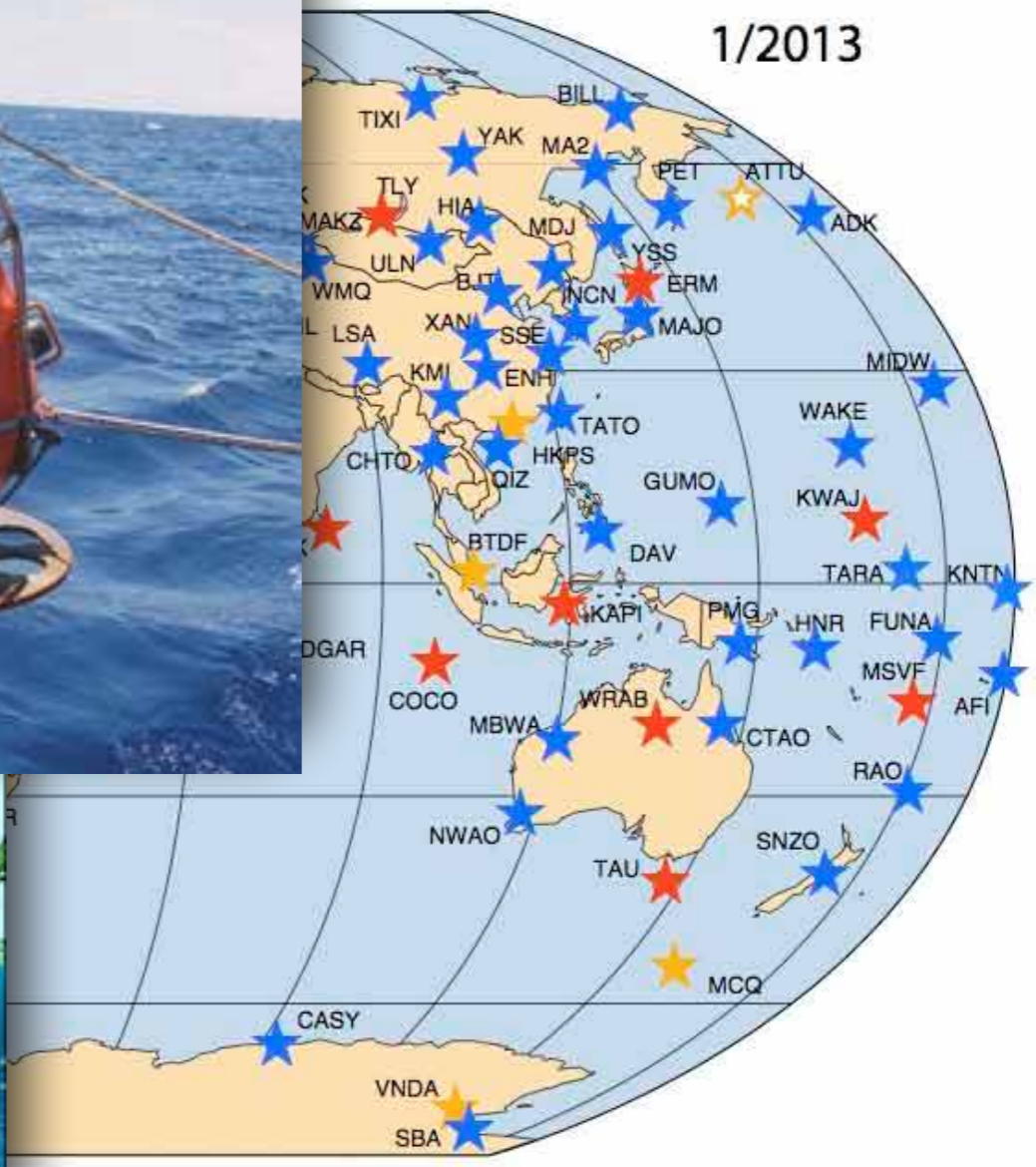


OBS

USGS

GLOBAL SEISMIC NETWORK

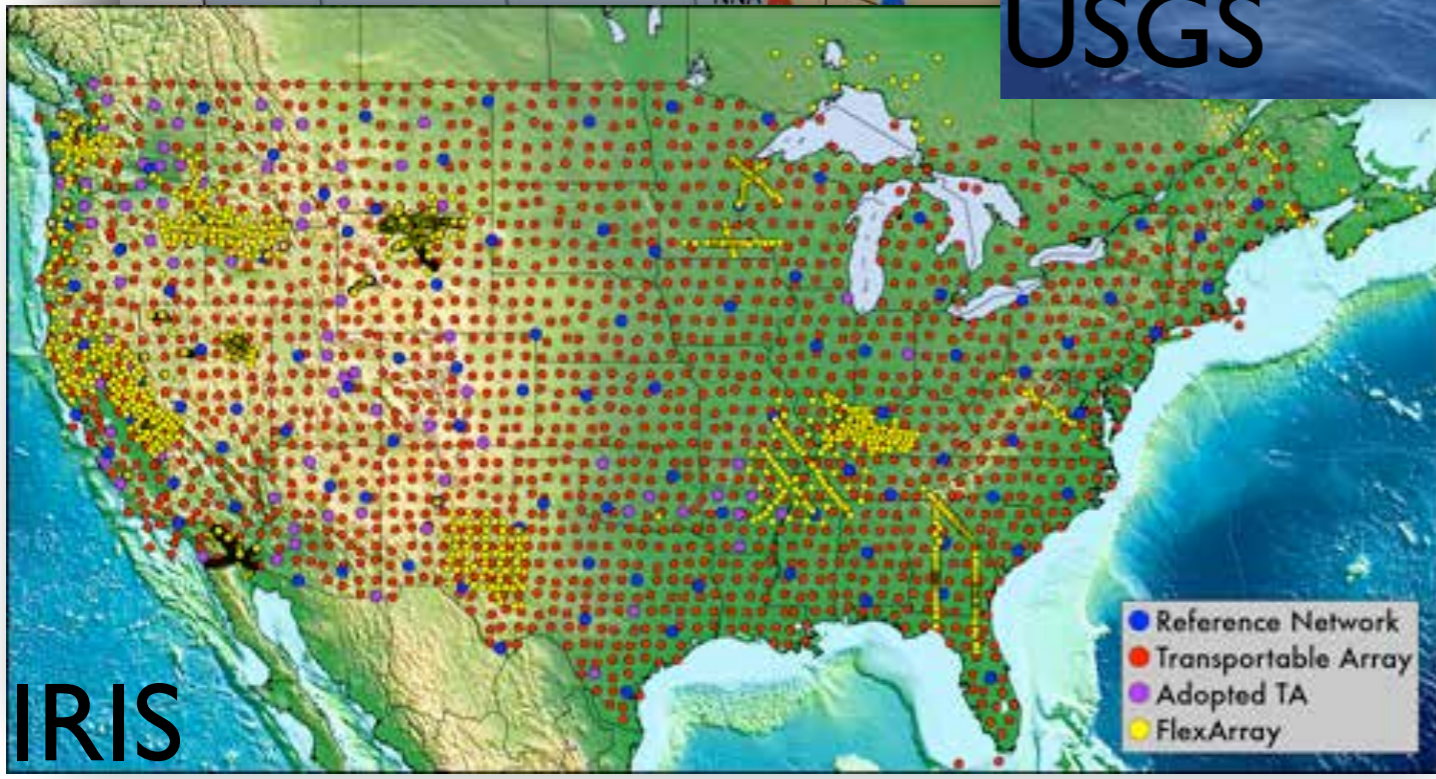
1/2013



Reference Stations

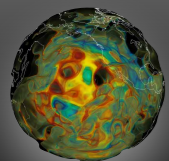


Affiliate Stations



IRIS

- Reference Network
- Transportable Array
- Adopted TA
- FlexArray



Seismic stations

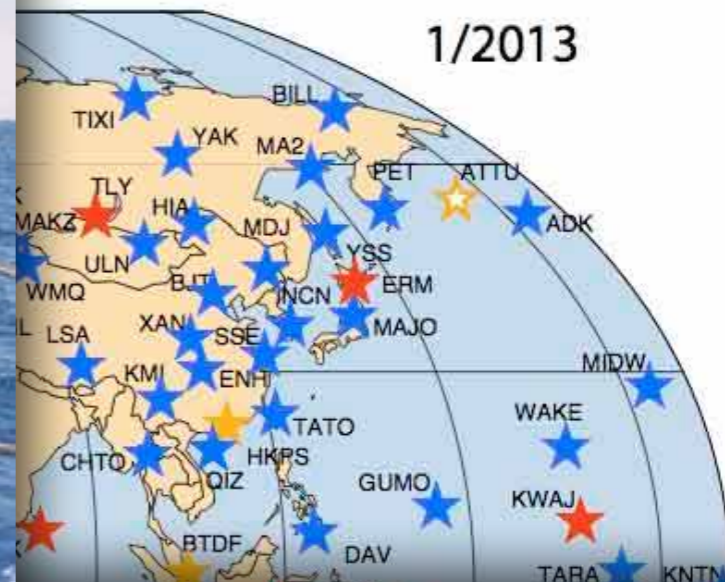


GLOBAL SEISMIC NETWORK



OBS

NETWORK

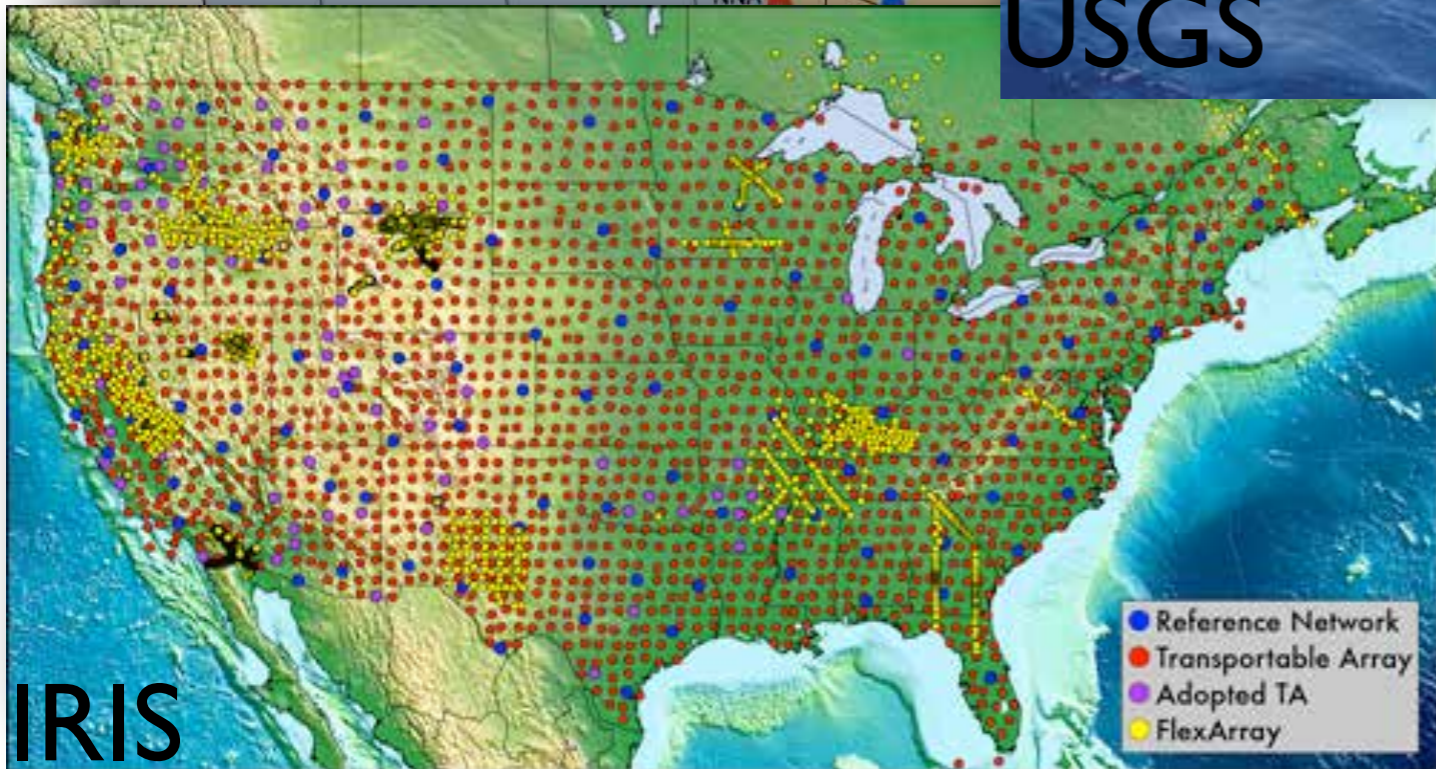


MERMAIDS

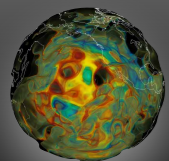


Courtesy Guust Nolet

USGS



IRIS

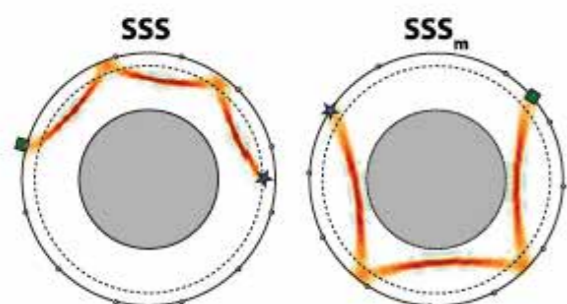


Data used in classical global tomography

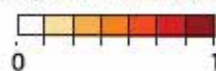
Body-wave traveltimes

Normal modes

Ritsema et al. (2011)



Ray sensitivity in S40RTS parameterization

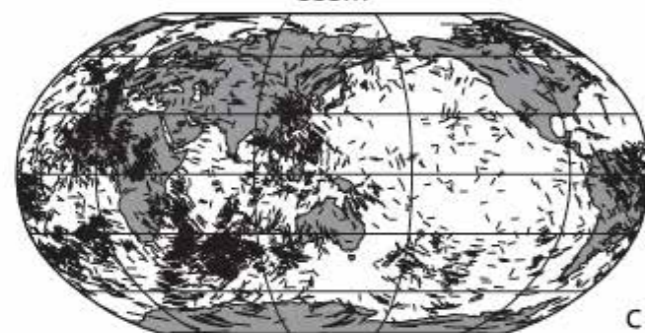


SSS



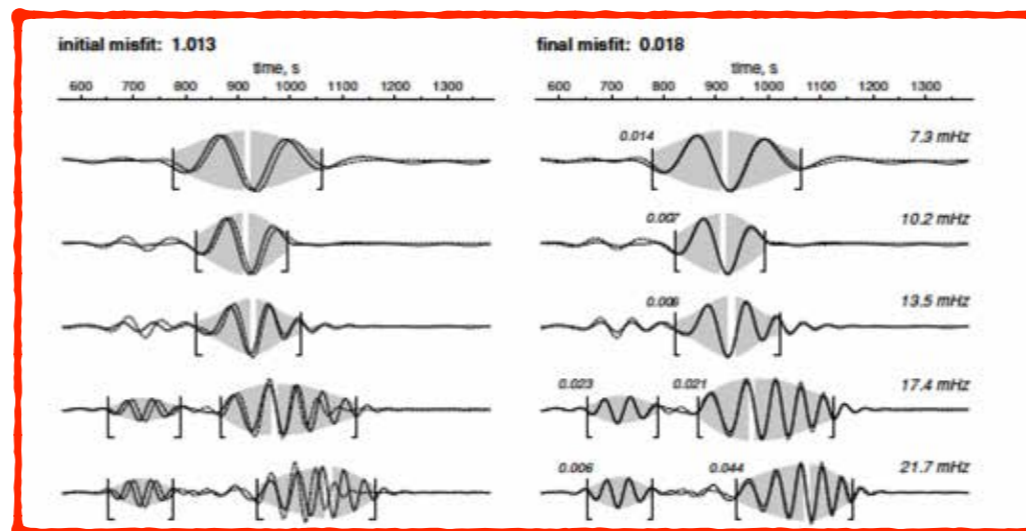
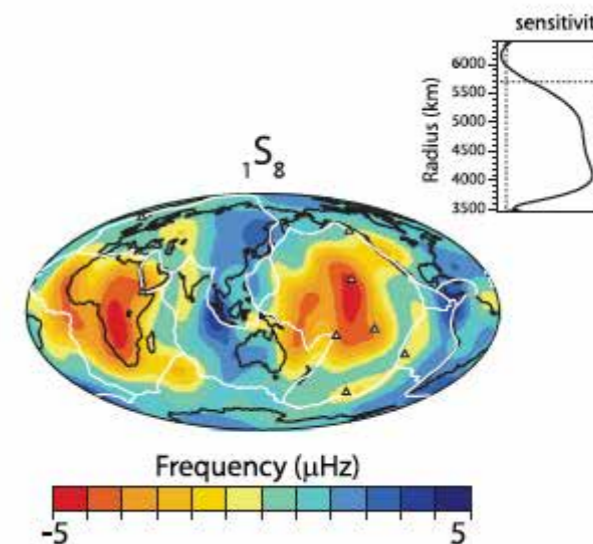
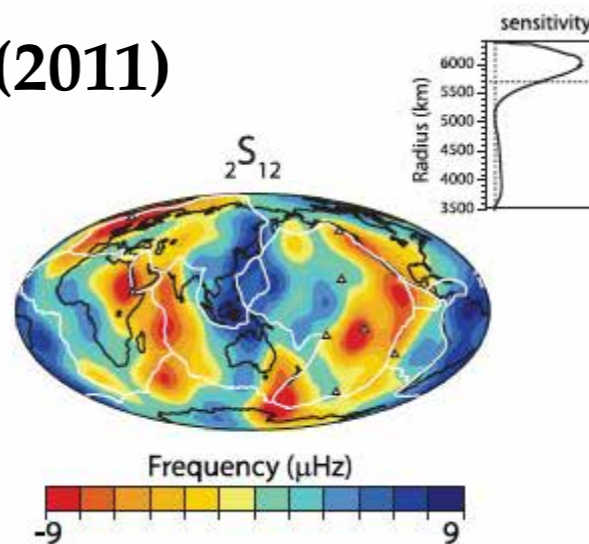
A

SSSm



B

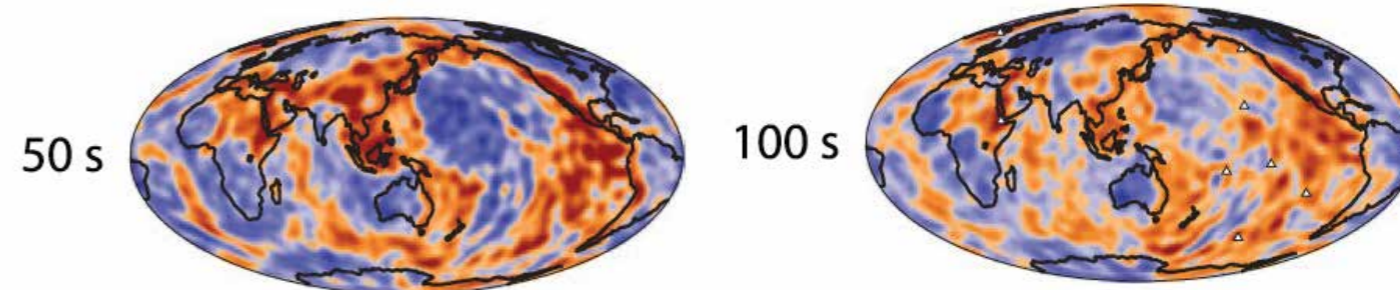
C

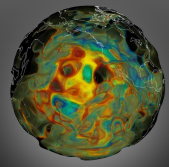


Waveforms

Lebedev et al. (2005)

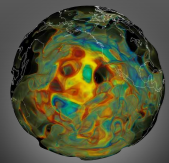
Minor- & major-arc surface-wave phase/group speeds



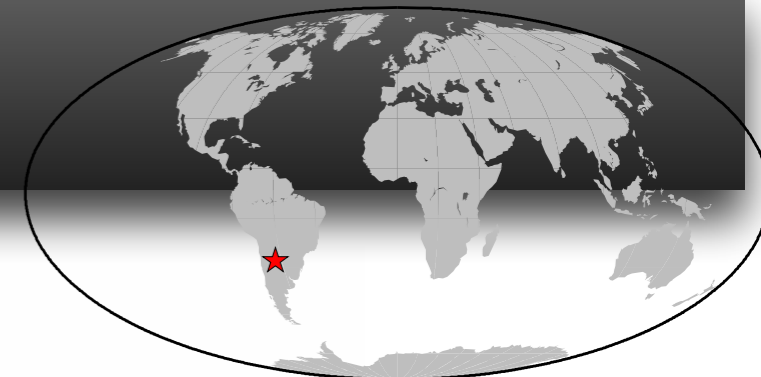


Challenges in global tomography

- **We are at a stage where details matter!**
 - *Avoid approximations as much as possible*
 - **Seismograms & Fréchet sensitivity kernels need to be calculated numerically in a realistic 3D Earth model**
- **Data coverage**
 - **Source-station distribution can not be changed, but we can change the amount of information that we get from every seismogram**



We can use complete seismograms

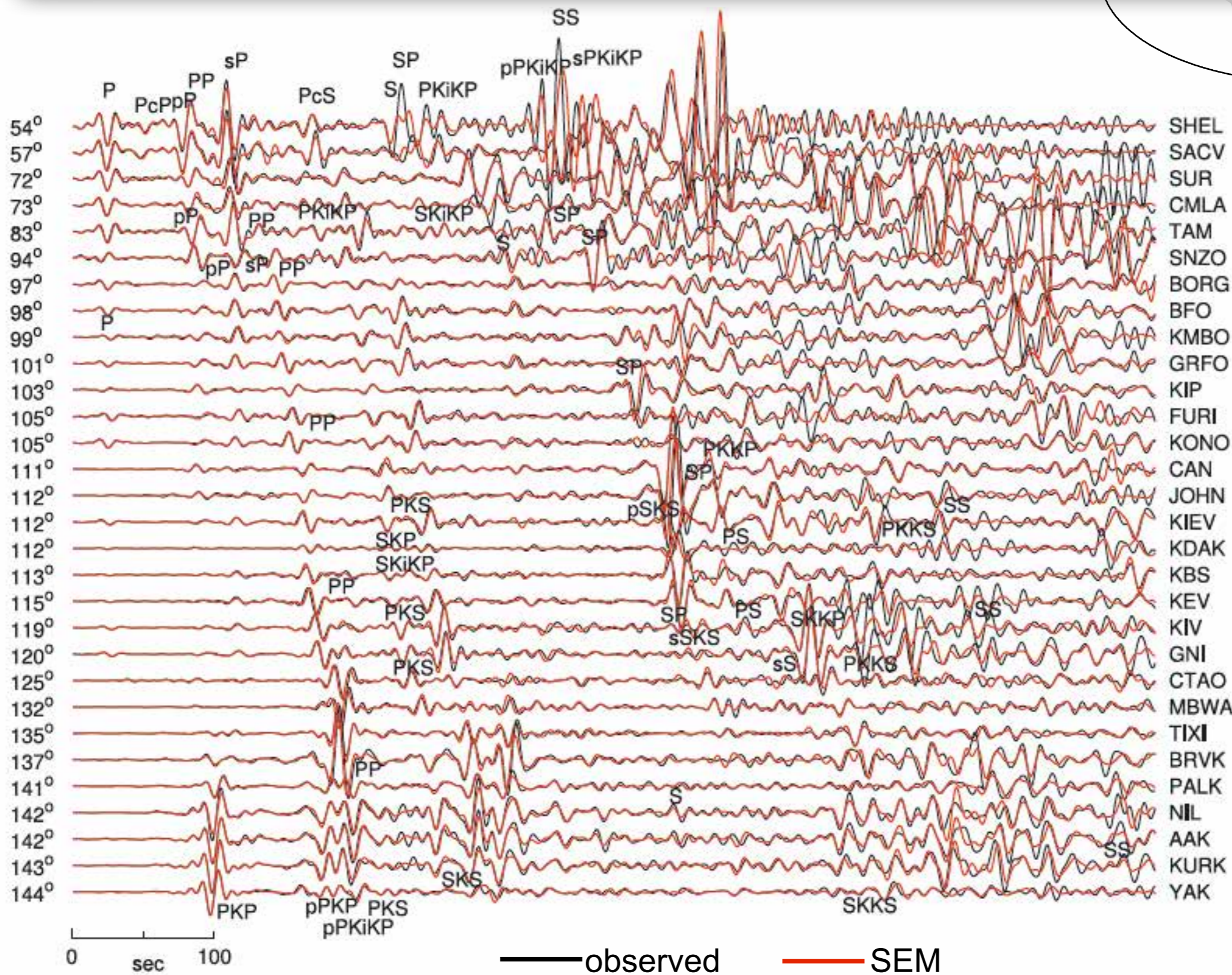


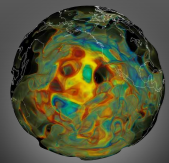
Sept. 3, 2008,
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SPECFEM3D_GLOBE

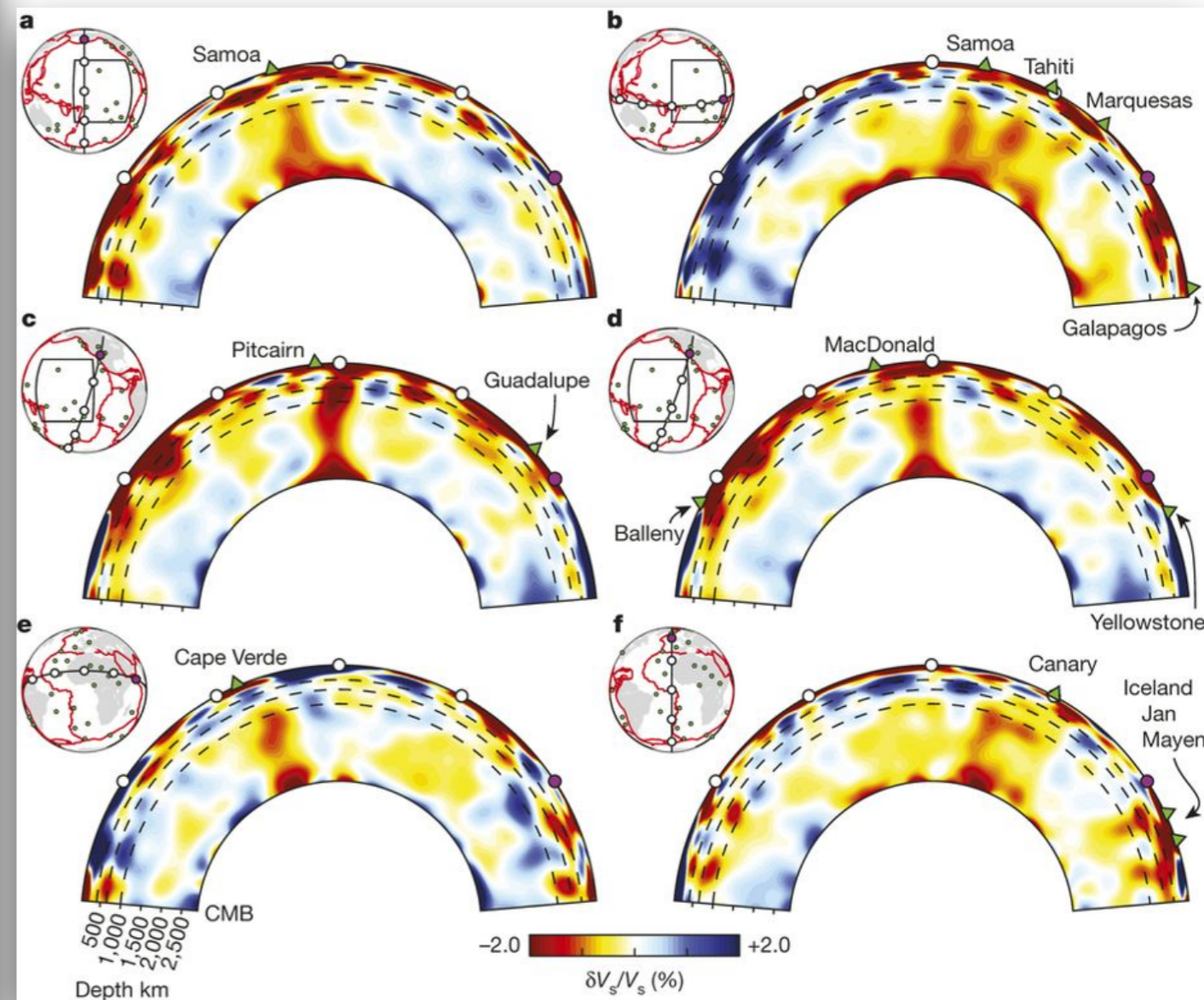
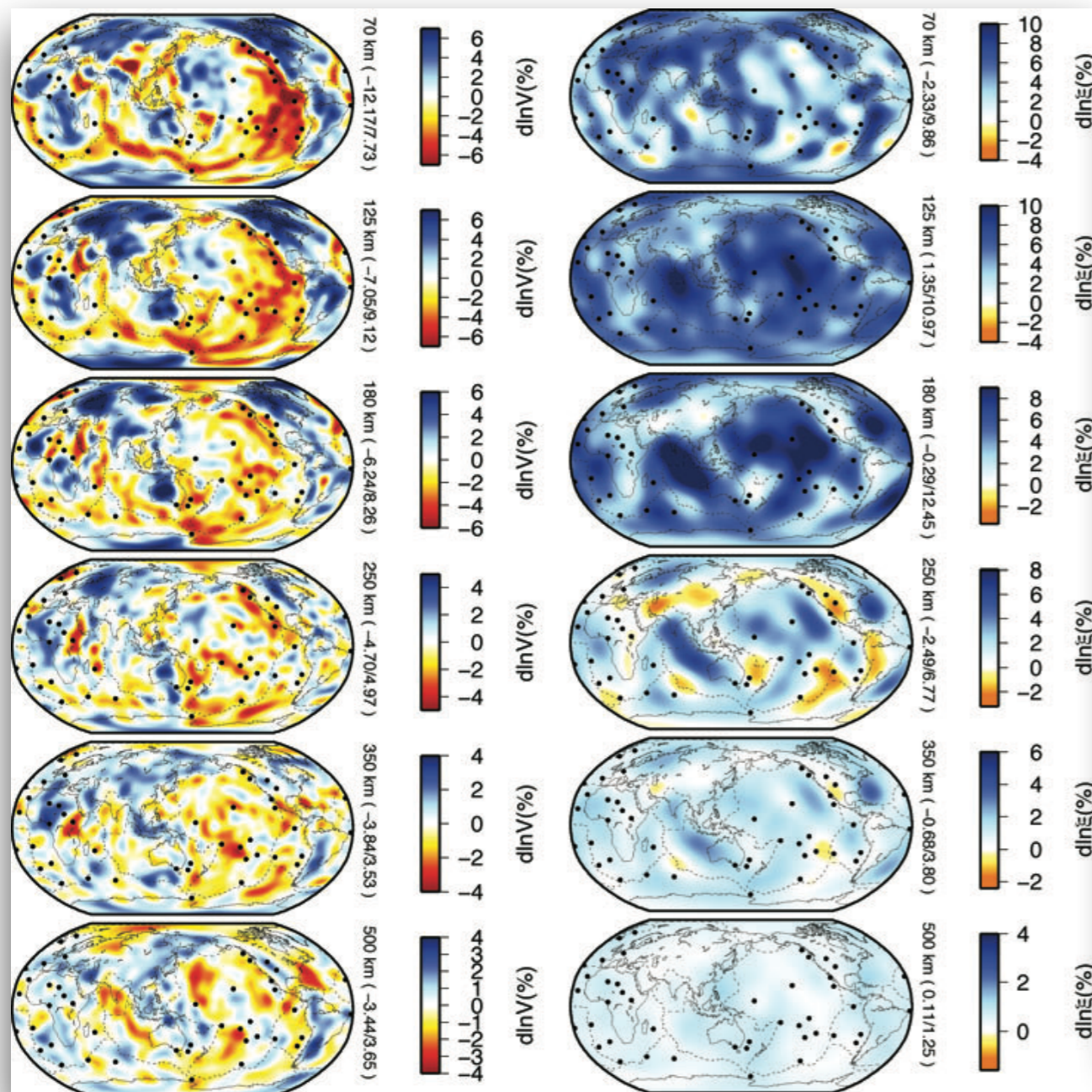
3D Earth model:
S362ANI +
Crust2.0

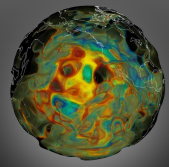
17 - 60 s





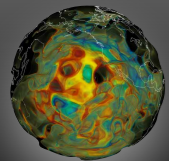
3D seismograms with NACT kernels



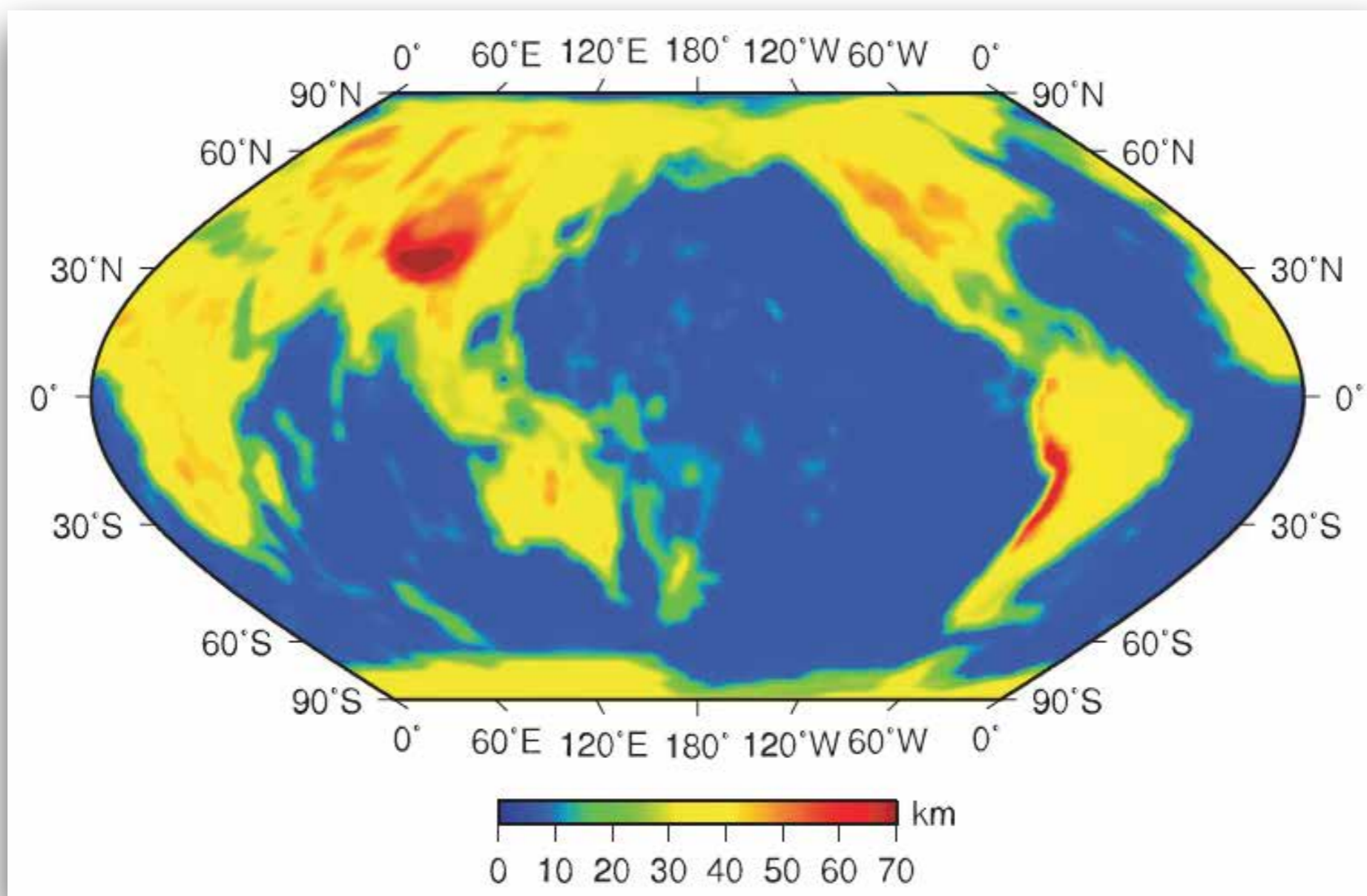


Challenges in global tomography

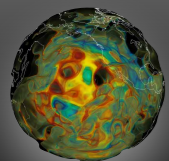
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- **Data coverage**
 - **Source-station distribution can not be changed, but we can change the amount of information that we get from every seismogram**
- **Crust can be highly nonlinear**



Smoothed Moho depths



Crust2.0 (Bassin et al. 2000)



Radial anisotropy

$$\xi = \frac{V_{SH}^2}{V_{SV}^2}$$

Perturbations of radial anisotropy (ξ)

S362WMANI+M

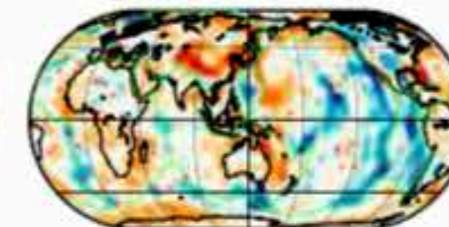
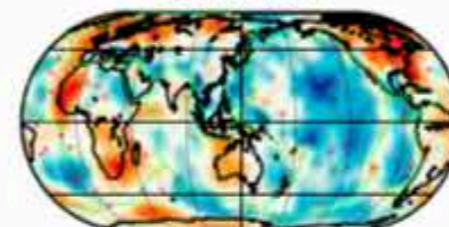
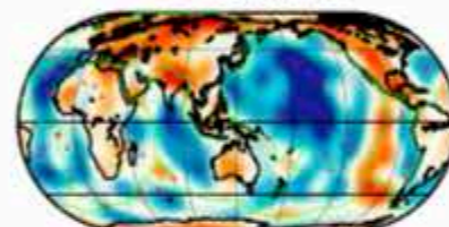
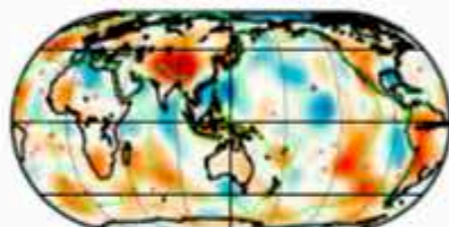
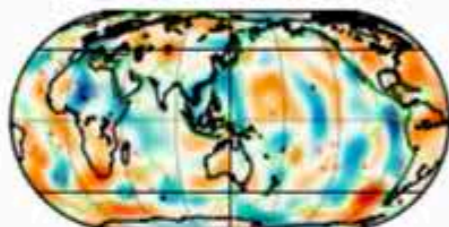
SAW642ANb

SEMum2

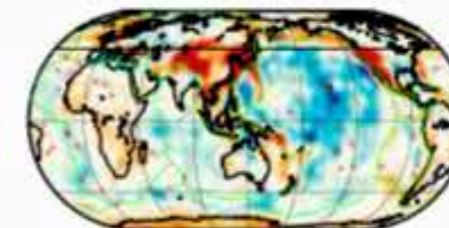
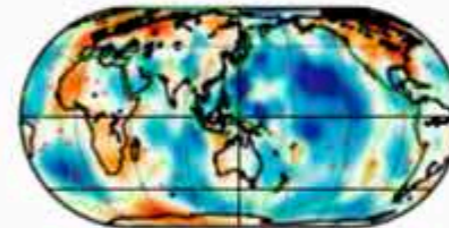
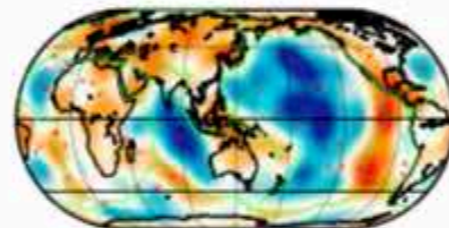
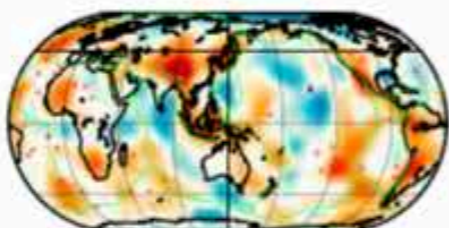
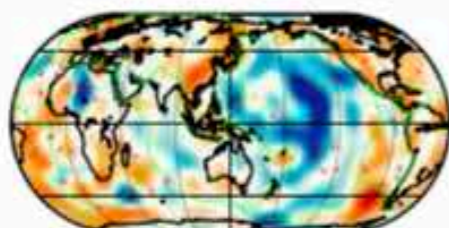
SAVANI

SGLOBE-rani

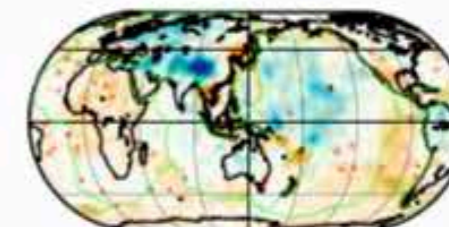
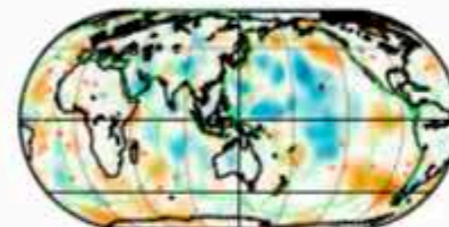
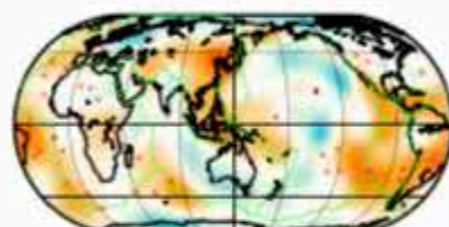
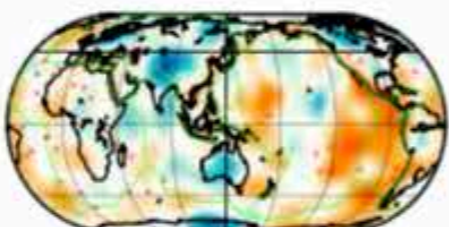
100 km
(max=5%)



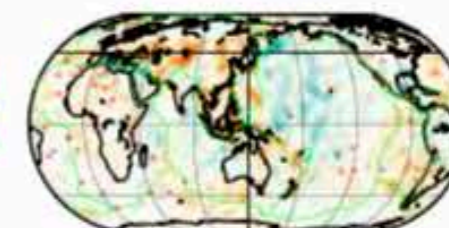
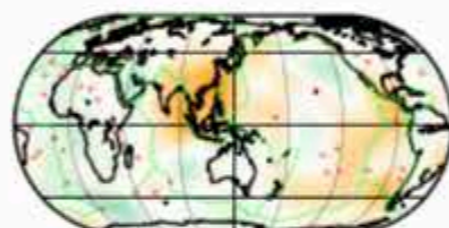
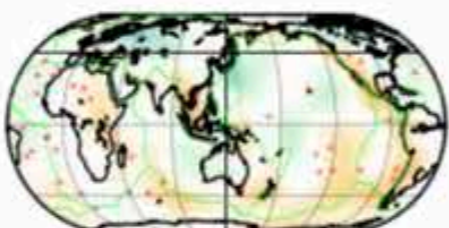
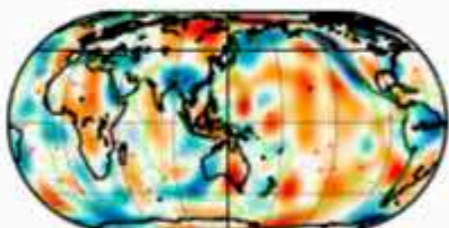
150 km
(max=5%)



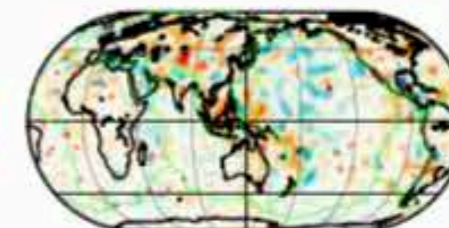
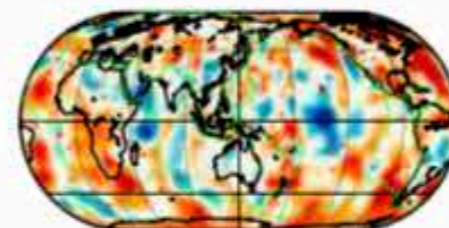
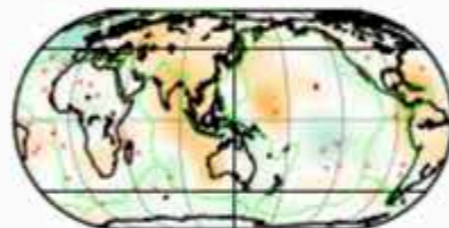
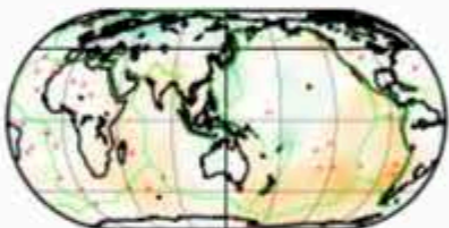
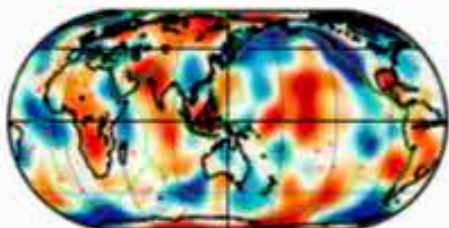
250 km
(max=5%)



400 km
(max=3%)

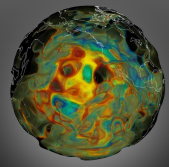


600 km
(max=2%)



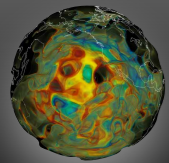
$d\xi/\xi$



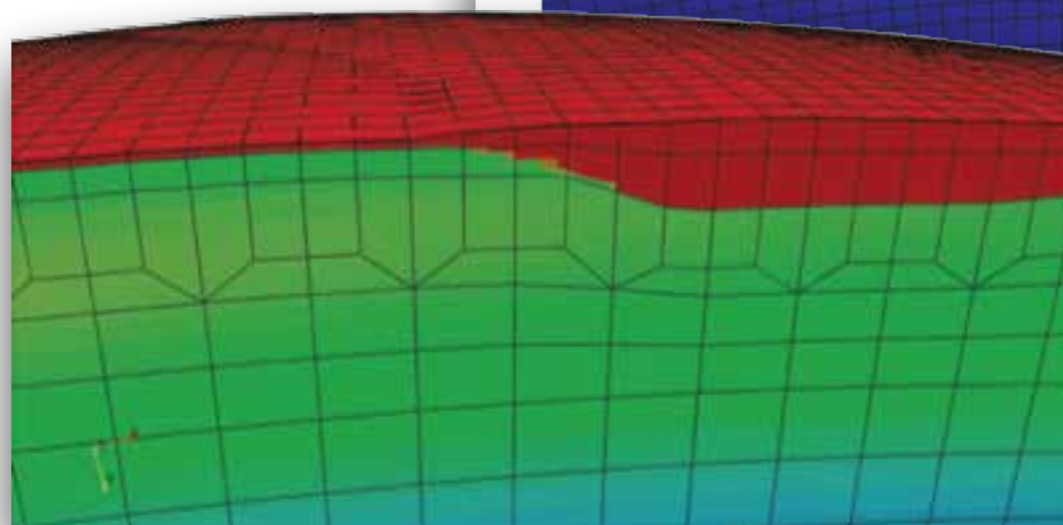
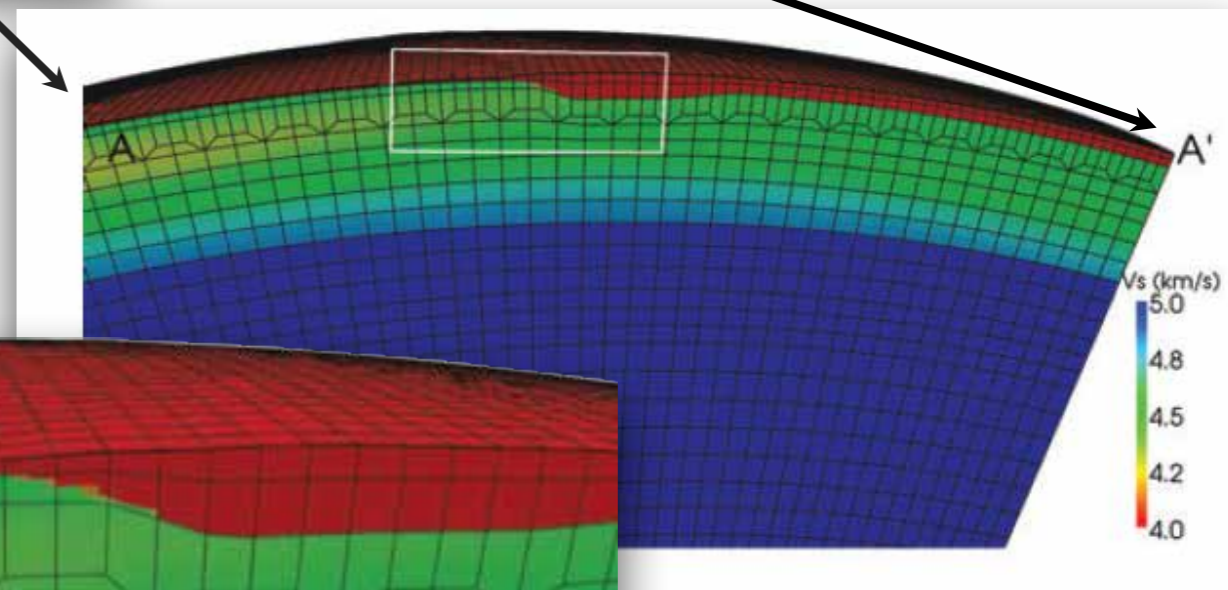
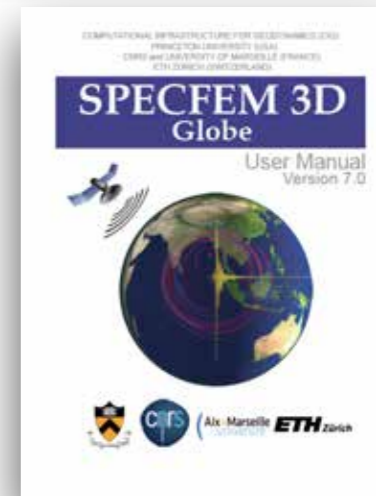
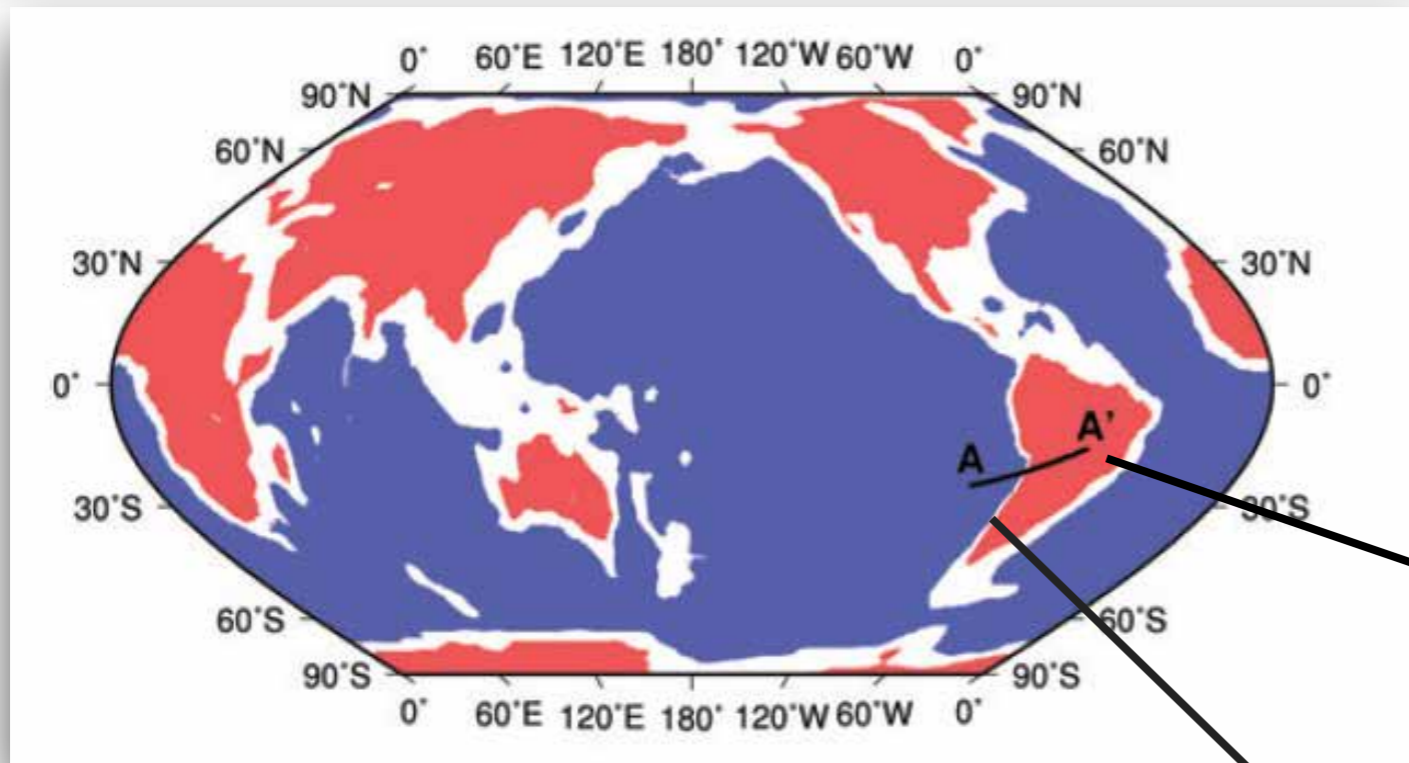


Challenges in global tomography

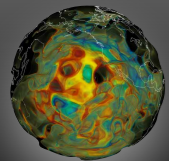
- **We are at a stage where details matter!**
 - *Avoid approximations as much as possible*
 - **Seismograms & Fréchet sensitivity kernels need to be calculated numerically in a realistic 3D Earth model**
- **Data coverage**
 - **Source-station distribution can not be changed, but we can change the amount of information that we get from every seismogram**
- **Crust can be highly nonlinear**
 - **Crustal corrections are questionable**
 - **Crust must be addressed properly in inversions**



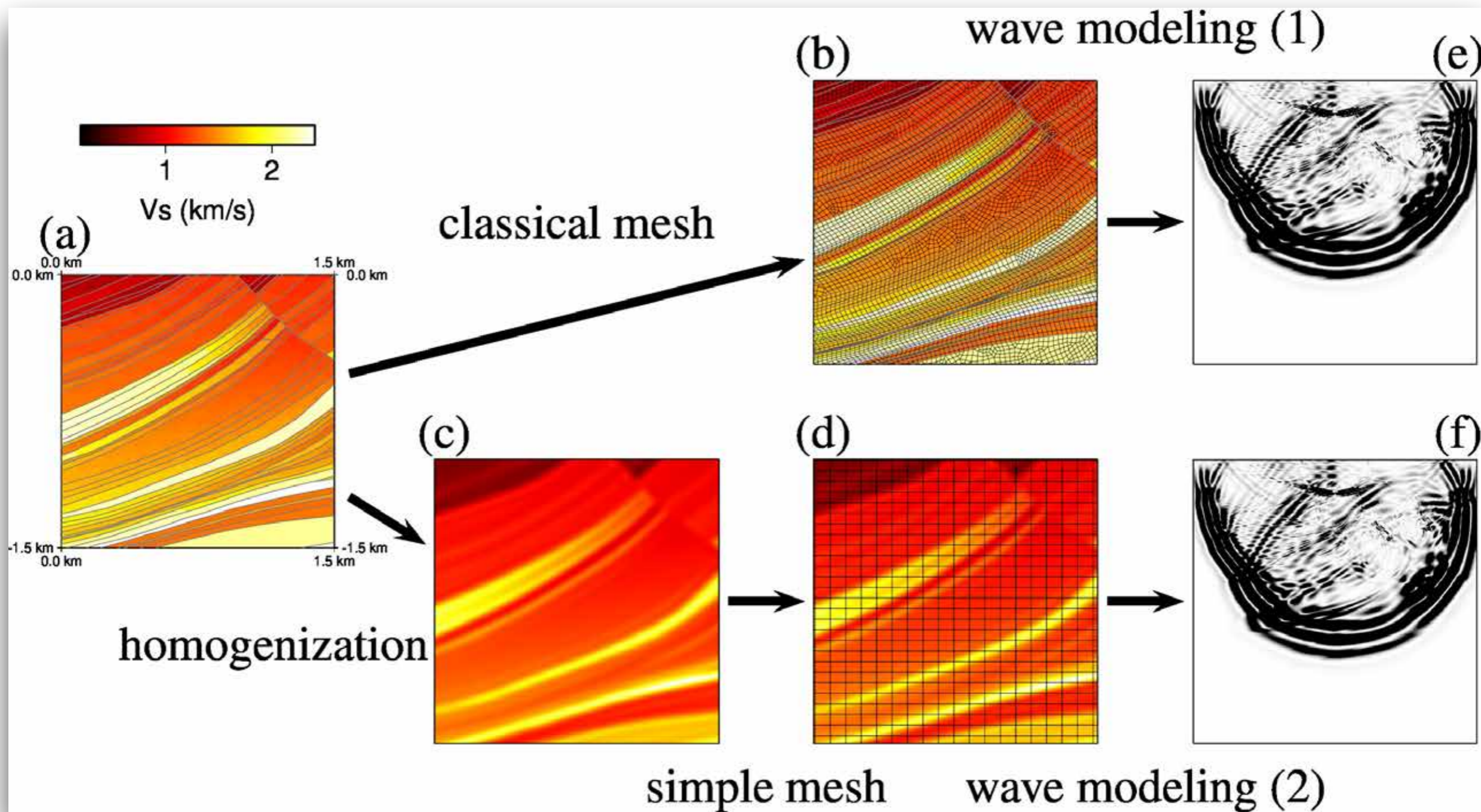
Dealing with crust in global simulations

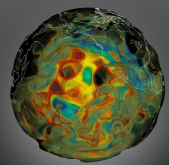


Moho is honored
if crustal thickness is
 ≤ 15 km and ≥ 35 km.



Homogenization





Non-linear inverse problems

\mathbf{m} is an initial model

$\mathbf{m} + \delta\mathbf{m}$ let's perturb it a bit!

χ a misfit function (waveform difference, travelttime difference, etc.)

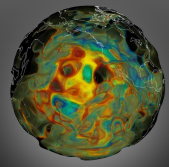
Taylor expansion $\chi(\mathbf{m} + \delta\mathbf{m}) \approx \chi(\mathbf{m}) + \mathbf{g}(\mathbf{m})^T \delta\mathbf{m} + \frac{1}{2} \delta\mathbf{m}^T \mathbf{H}(\mathbf{m}) \delta\mathbf{m}$

gradient $\leftarrow \mathbf{g}(\mathbf{m}) = \left. \frac{\partial \chi}{\partial \mathbf{m}} \right|_{\mathbf{m}}$

Hessian $\leftarrow \mathbf{H}(\mathbf{m}) = \left. \frac{\partial^2 \chi}{\partial \mathbf{m} \partial \mathbf{m}} \right|_{\mathbf{m}}$

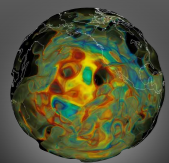
$$\mathbf{g}(\mathbf{m} + \delta\mathbf{m}) \approx \mathbf{g}(\mathbf{m}) + \mathbf{H}(\mathbf{m}) \delta\mathbf{m}$$

- If we have the Hessian & the gradient ==> Gauss-Newton method
- What if we do not have access to the Hessian?



Gradient-based optimizations

- **Steepest descent**
- **Conjugate gradient**
- **BFGS (Broyden–Fletcher–Goldfarb–Shanno)**
- **L-BFGS (Limited-memory BFGS)**
- ...



Gradient-based optimizations

calculate misfit

$\chi(\mathbf{m}^0) \rightarrow$ initial model

calculate gradient

$$\mathbf{g}^0 = \partial \chi / \partial \mathbf{m} (\mathbf{m}^0)$$

calculate gradient search
directions

$$\mathbf{p}^0 = -\mathbf{g}^0$$

perform a line search

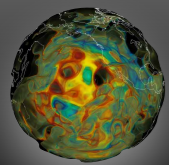
$$\tilde{\chi}^k(v) = \chi(\mathbf{m}^k + v\mathbf{p}^k)$$

$$\tilde{\mathbf{g}}^k(v) = \frac{\partial \tilde{\chi}^k}{\partial v} = \frac{\partial \chi}{\partial \mathbf{m}} (\mathbf{m}^k + v\mathbf{p}^k) \cdot \mathbf{p}^k$$

calculate test models

$$\mathbf{m}_t^k = \mathbf{m}^k + v_t^k \mathbf{p}^k$$

calculate $\chi(\mathbf{m}_t^k)$ for cubic interpolation



Gradient-based optimizations

update the model

$$\mathbf{m}^{k+1} = \mathbf{m}^k + \nu^k \mathbf{p}^k$$

recalculate the gradient

$$\mathbf{g}^{k+1} = \partial \chi / \partial \mathbf{m} (\mathbf{m}^{k+1})$$

recalculate the search direction

$$\mathbf{p}^{k+1} = -\mathbf{g}^{k+1} \longrightarrow \text{steepest descent}$$

conjugate gradient

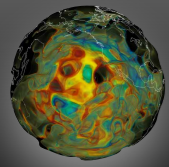


$$\mathbf{p}^{k+1} = -\mathbf{g}^{k+1} + \beta^{k+1} \mathbf{p}^k$$

$$\beta^{k+1} = \mathbf{g}^{k+1} \cdot (\mathbf{g}^{k+1} - \mathbf{g}^k) / (\mathbf{g}^k \cdot \mathbf{g}^k)$$

L-BFGS (Limited-memory BFGS)

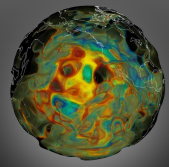
If $\|\mathbf{p}^{k+1}\| < \epsilon \longrightarrow \mathbf{m}^{k+1}$ is the desired model



Adjoint tomography

misfit function

$$\chi = \frac{1}{2} \sum_{r=1}^N \int_0^T [\mathbf{s}(\mathbf{x}_r, t, \mathbf{m}) - \mathbf{d}(\mathbf{x}_r, t)]^2 dt$$



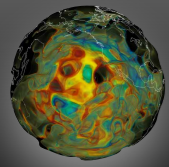
Adjoint tomography

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$$\chi = \frac{1}{2} \sum_{r=1}^N \int_0^T [\mathbf{s}(\mathbf{x}_r, t, \mathbf{m}) - \mathbf{d}(\mathbf{x}_r, t)]^2 dt$$

gradient

$$\delta\chi = \sum_{r=1}^N \int_0^T [\mathbf{s}(\mathbf{x}_r, t, \mathbf{m}) - \mathbf{d}(\mathbf{x}_r, t)] \delta\mathbf{s}(\mathbf{x}_r, t, \mathbf{m}) dt$$



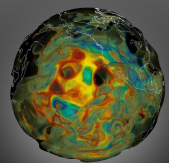
Adjoint tomography

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Adjoint tomography

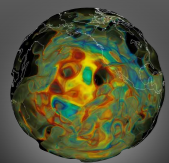
misfit function

$$\chi = \frac{1}{2} \sum_{r=1}^N \int_0^T [\mathbf{s}(\mathbf{x}_r, t, \mathbf{m}) - \mathbf{d}(\mathbf{x}_r, t)]^2 dt$$

gradient

$$\delta\chi = \sum_{r=1}^N \int_0^T [\mathbf{s}(\mathbf{x}_r, t, \mathbf{m}) - \mathbf{d}(\mathbf{x}_r, t)] \delta\mathbf{s}(\mathbf{x}_r, t, \mathbf{m}) dt$$

$$\delta\chi = - \sum_{r=1}^N \int_0^T [s_i(\mathbf{x}_r, t) - d_i(\mathbf{x}_r, t)] \int_0^t \int_V [\delta\rho(\mathbf{x}') G_{ij}(\mathbf{x}_r, \mathbf{x}'; t - t') \partial_{t'}^2 s_j(\mathbf{x}', t') + \delta c_{jklm}(\mathbf{x}') \partial'_k G_{ij}(\mathbf{x}_r, \mathbf{x}'; t - t') \partial'_l s_m(\mathbf{x}', t')] d^3\mathbf{x}' dt' dt.$$



Adjoint tomography

misfit function

$$\chi = \frac{1}{2} \sum_{r=1}^N \int_0^T [\mathbf{s}(\mathbf{x}_r, t, \mathbf{m}) - \mathbf{d}(\mathbf{x}_r, t)]^2 dt$$

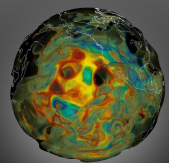
gradient

$$\delta\chi = \sum_{r=1}^N \int_0^T [\mathbf{s}(\mathbf{x}_r, t, \mathbf{m}) - \mathbf{d}(\mathbf{x}_r, t)] \delta\mathbf{s}(\mathbf{x}_r, t, \mathbf{m}) dt$$

$$\delta\chi = - \sum_{r=1}^N \int_0^T [s_i(\mathbf{x}_r, t) - d_i(\mathbf{x}_r, t)] \int_0^t \int_V [\delta\rho(\mathbf{x}') G_{ij}(\mathbf{x}_r, \mathbf{x}'; t - t') \partial_{t'}^2 s_j(\mathbf{x}', t') + \delta c_{jklm}(\mathbf{x}') \partial'_k G_{ij}(\mathbf{x}_r, \mathbf{x}'; t - t') \partial'_l s_m(\mathbf{x}', t')] d^3\mathbf{x}' dt' dt.$$

$$G_{ij}(\mathbf{x}, \mathbf{x}'; t - t') = G_{ji}(\mathbf{x}', \mathbf{x}; t' - t)$$

$$t \longrightarrow T - t$$



Adjoint tomography

misfit function

$$\chi = \frac{1}{2} \sum_{r=1}^N \int_0^T [\mathbf{s}(\mathbf{x}_r, t, \mathbf{m}) - \mathbf{d}(\mathbf{x}_r, t)]^2 dt$$

gradient

$$\delta\chi = \sum_{r=1}^N \int_0^T [\mathbf{s}(\mathbf{x}_r, t, \mathbf{m}) - \mathbf{d}(\mathbf{x}_r, t)] \delta\mathbf{s}(\mathbf{x}_r, t, \mathbf{m}) dt$$

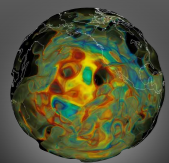
$$\delta\chi = - \sum_{r=1}^N \int_0^T [s_i(\mathbf{x}_r, t) - d_i(\mathbf{x}_r, t)] \int_0^t \int_V [\delta\rho(\mathbf{x}') G_{ij}(\mathbf{x}_r, \mathbf{x}'; t - t') \partial_{t'}^2 s_j(\mathbf{x}', t') + \delta c_{jklm}(\mathbf{x}') \partial'_k G_{ij}(\mathbf{x}_r, \mathbf{x}'; t - t') \partial'_l s_m(\mathbf{x}', t')] d^3\mathbf{x}' dt' dt.$$

$$G_{ij}(\mathbf{x}, \mathbf{x}'; t - t') = G_{ji}(\mathbf{x}', \mathbf{x}; t' - t)$$

$$t \longrightarrow T - t$$

adjoint
wavefield

$$s_k^\dagger(\mathbf{x}', t') = \sum_{r=1}^N \int_0^{t'} \int_V G_{ki}(\mathbf{x}', \mathbf{x}_r; t' - t) [s_i(\mathbf{x}_r, T - t) - d_i(\mathbf{x}_r, T - t)] d^3\mathbf{x} dt$$



Adjoint tomography

misfit function

$$\chi = \frac{1}{2} \sum_{r=1}^N \int_0^T [\mathbf{s}(\mathbf{x}_r, t, \mathbf{m}) - \mathbf{d}(\mathbf{x}_r, t)]^2 dt$$

gradient

$$\delta\chi = \sum_{r=1}^N \int_0^T [\mathbf{s}(\mathbf{x}_r, t, \mathbf{m}) - \mathbf{d}(\mathbf{x}_r, t)] \delta\mathbf{s}(\mathbf{x}_r, t, \mathbf{m}) dt$$

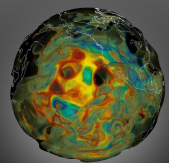
$$\begin{aligned} \delta\chi = & - \sum_{r=1}^N \int_0^T [s_i(\mathbf{x}_r, t) - d_i(\mathbf{x}_r, t)] \int_0^t \int_V [\delta\rho(\mathbf{x}') G_{ij}(\mathbf{x}_r, \mathbf{x}'; t - t') \partial_{t'}^2 s_j(\mathbf{x}', t') \\ & + \delta c_{jklm}(\mathbf{x}') \partial'_k G_{ij}(\mathbf{x}_r, \mathbf{x}'; t - t') \partial'_l s_m(\mathbf{x}', t')] d^3\mathbf{x}' dt' dt. \end{aligned}$$

$$G_{ij}(\mathbf{x}, \mathbf{x}'; t - t') = G_{ji}(\mathbf{x}', \mathbf{x}; t' - t)$$

$$t \longrightarrow T - t$$

adjoint
wavefield

$$s_k^\dagger(\mathbf{x}', t') = \sum_{r=1}^N \int_0^{t'} \int_V G_{ki}(\mathbf{x}', \mathbf{x}_r; t' - t) [s_i(\mathbf{x}_r, T - t) - d_i(\mathbf{x}_r, T - t)] d^3\mathbf{x} dt$$



Adjoint tomography

misfit function

$$\chi = \frac{1}{2} \sum_{r=1}^N \int_0^T [\mathbf{s}(\mathbf{x}_r, t, \mathbf{m}) - \mathbf{d}(\mathbf{x}_r, t)]^2 dt$$

gradient

$$\delta\chi = \sum_{r=1}^N \int_0^T [\mathbf{s}(\mathbf{x}_r, t, \mathbf{m}) - \mathbf{d}(\mathbf{x}_r, t)] \delta\mathbf{s}(\mathbf{x}_r, t, \mathbf{m}) dt$$

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$$G_{ij}(\mathbf{x}, \mathbf{x}'; t - t') = G_{ji}(\mathbf{x}', \mathbf{x}; t' - t)$$

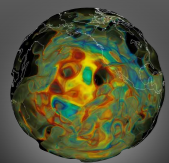
$$t \longrightarrow T - t$$

adjoint
wavefield

$$s_k^\dagger(\mathbf{x}', t') = \sum_{r=1}^N \int_0^{t'} \int_V G_{ki}(\mathbf{x}', \mathbf{x}_r; t' - t) [s_i(\mathbf{x}_r, T - t) - d_i(\mathbf{x}_r, T - t)] d^3\mathbf{x} dt$$

adjoint source

$$f_i^\dagger(\mathbf{x}, t) = \sum_{r=1}^N [s_i(\mathbf{x}_r, T - t) - d_i(\mathbf{x}_r, T - t)] \delta(\mathbf{x} - \mathbf{x}_r)$$

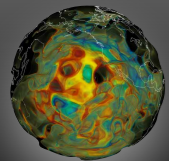


How to compute adjoint kernels?

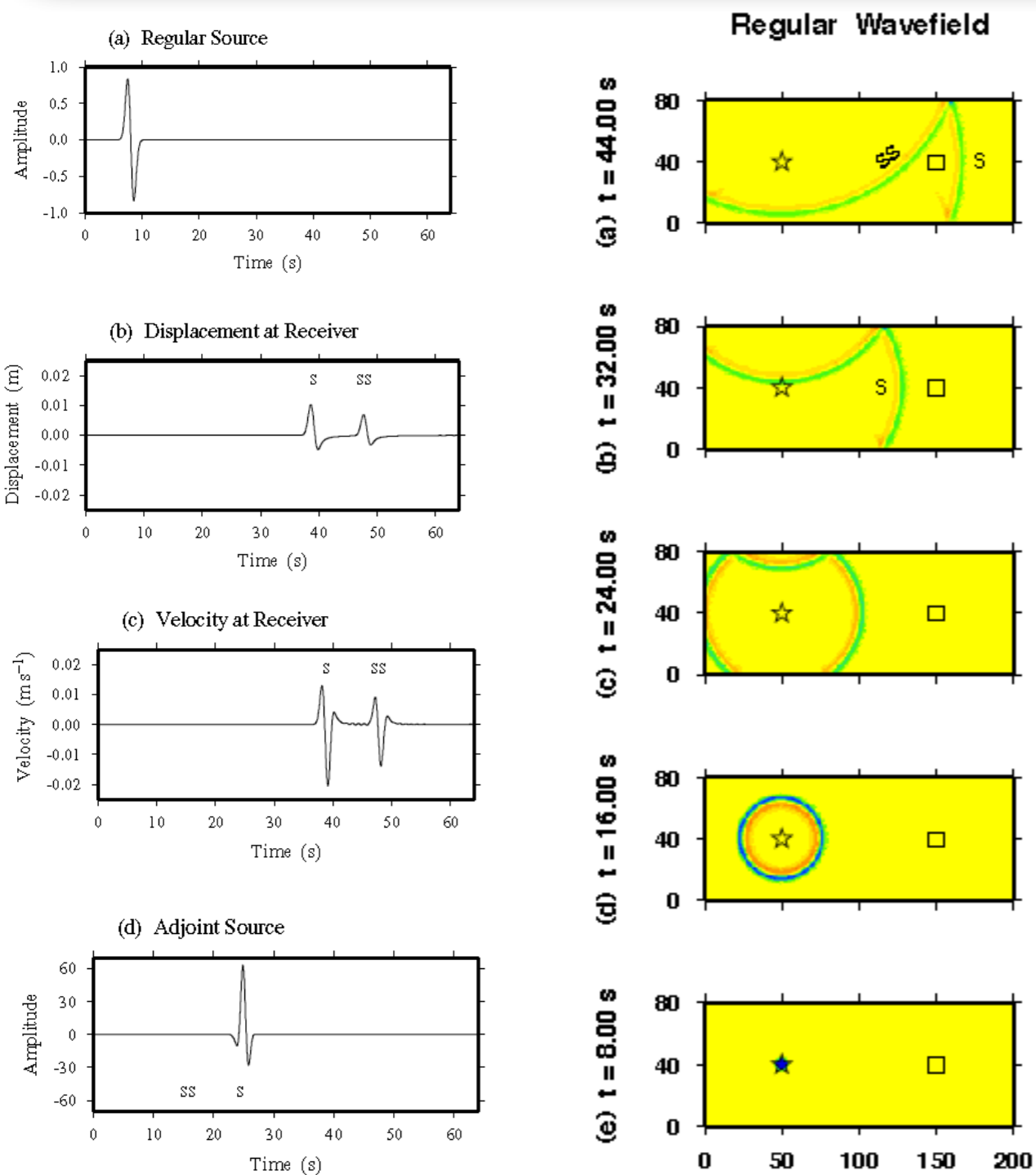
$$\delta\chi = \int_V [K_\rho(\mathbf{x})\delta \ln \rho(\mathbf{x}) + K_{c_{jklm}}(\mathbf{x})\delta \ln c_{jklm}(\mathbf{x})] d^3\mathbf{x}$$

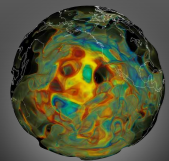
$$K_\rho(\mathbf{x}) = - \int_0^T \rho(\mathbf{x}) \mathbf{s}^\dagger(\mathbf{x}, T-t) \cdot \partial_t^2 \mathbf{s}(\mathbf{x}, t) dt$$

$$K_{c_{jklm}}(\mathbf{x}) = - \int_0^T \epsilon_{jk}^\dagger(\mathbf{x}, T-t) c_{jklm}(\mathbf{x}) \epsilon_{lm}(\mathbf{x}, t) dt$$

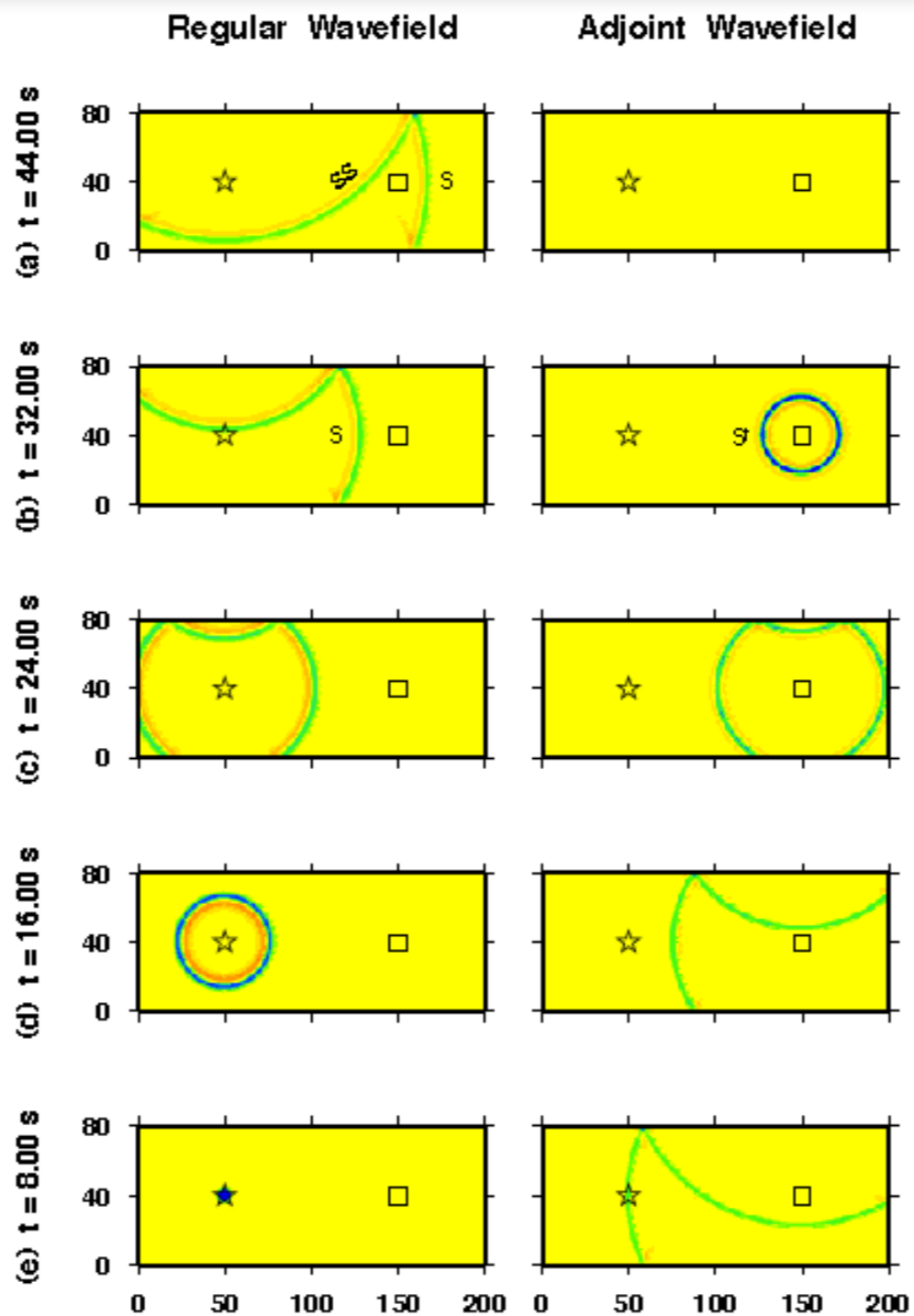
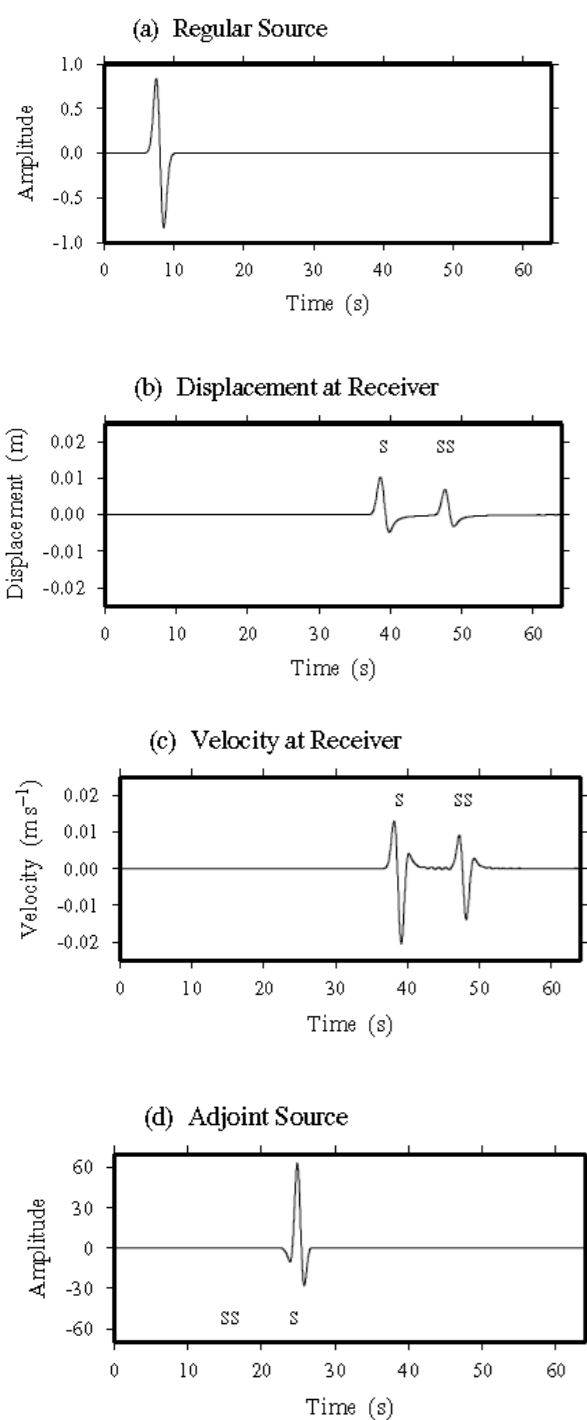


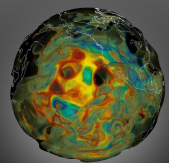
How to compute adjoint kernels?



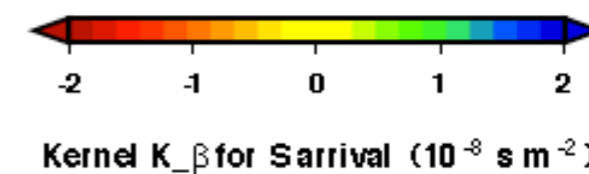
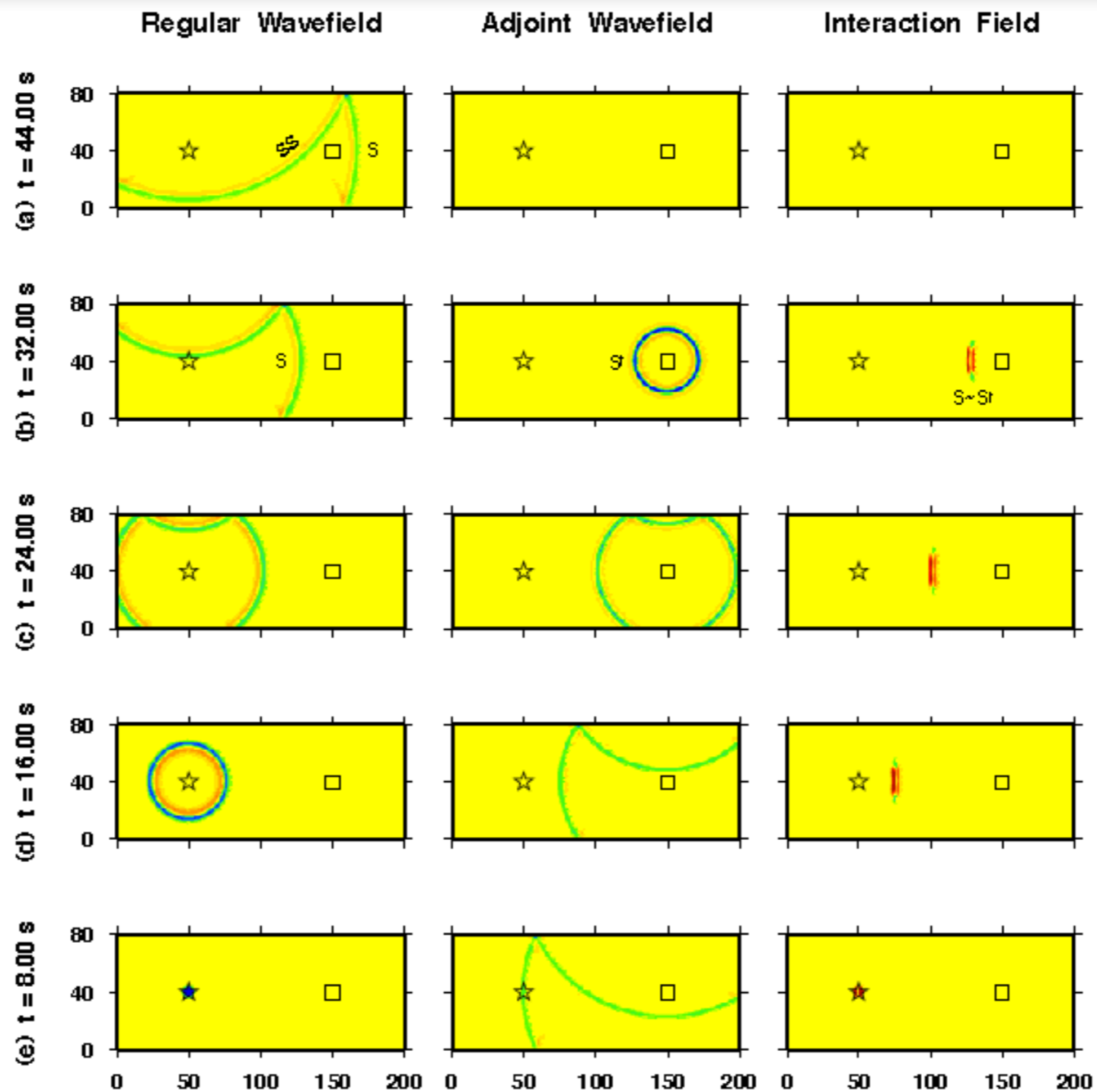
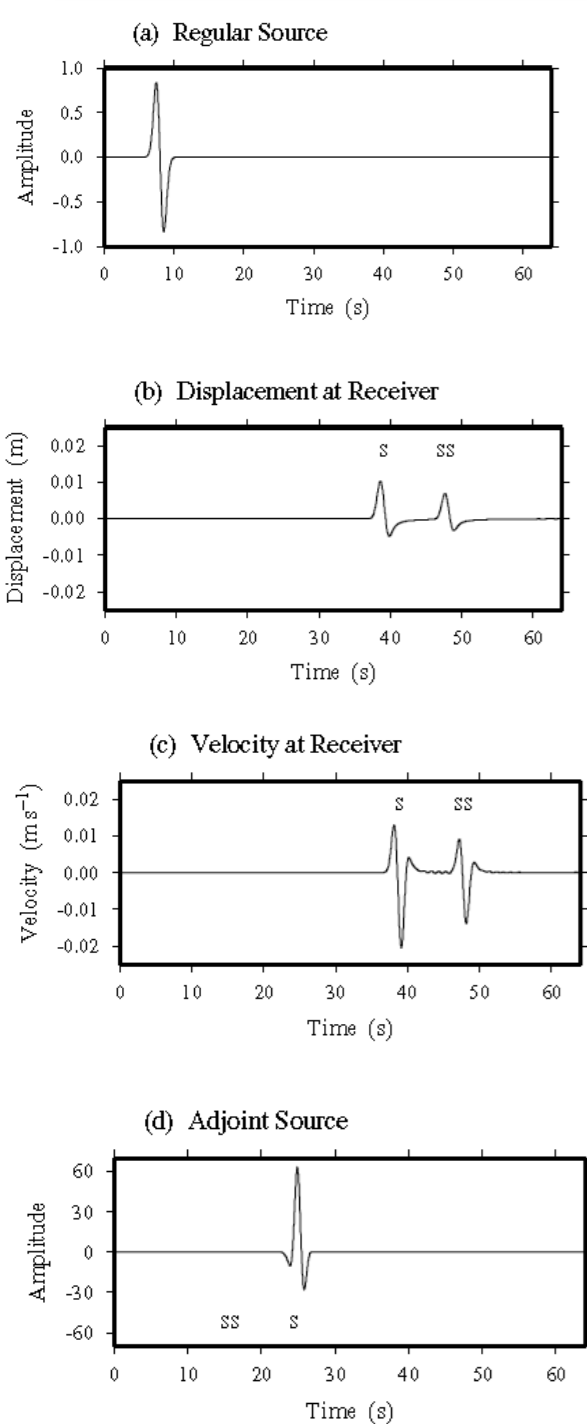


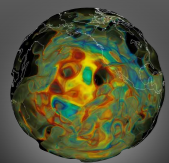
How to compute adjoint kernels?



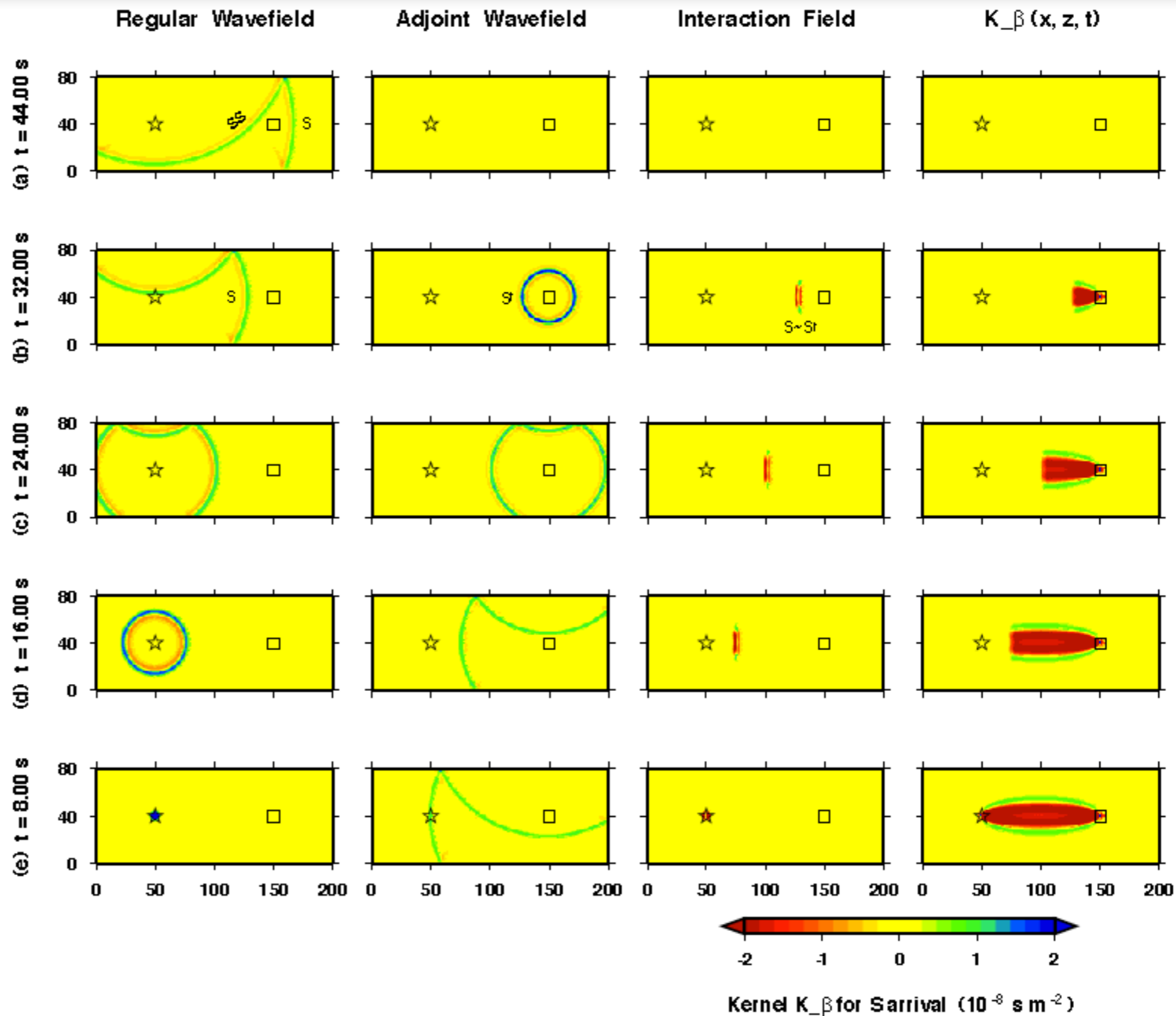
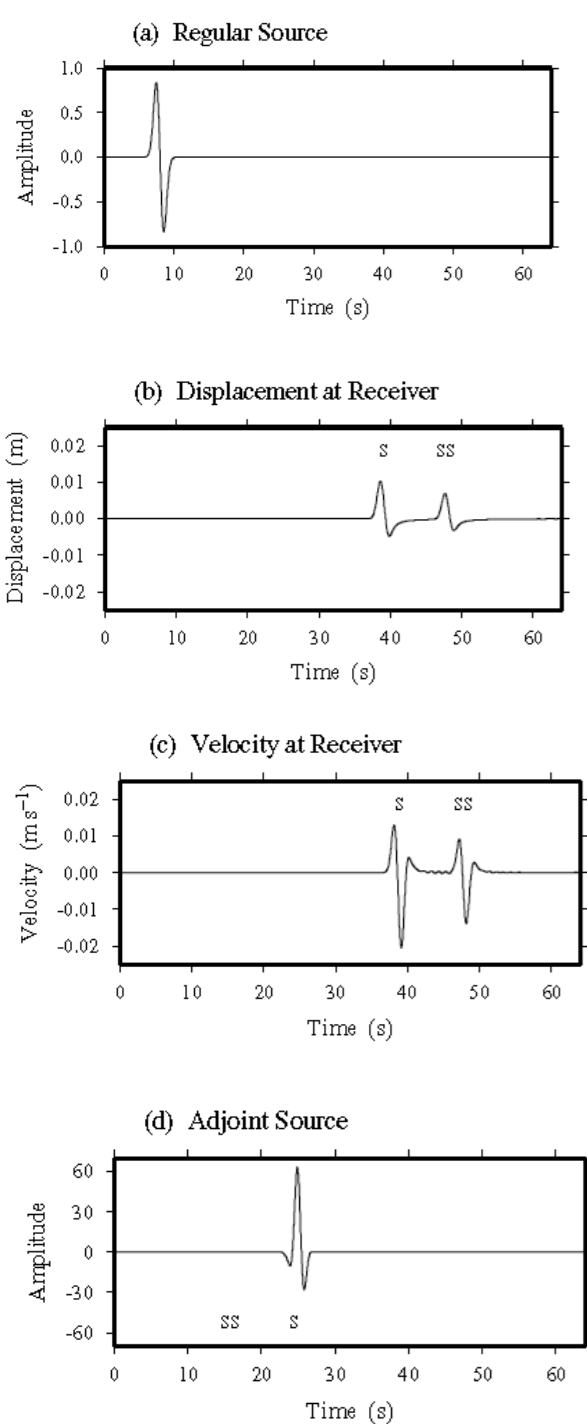


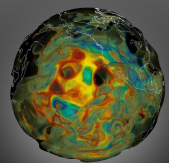
How to compute adjoint kernels?



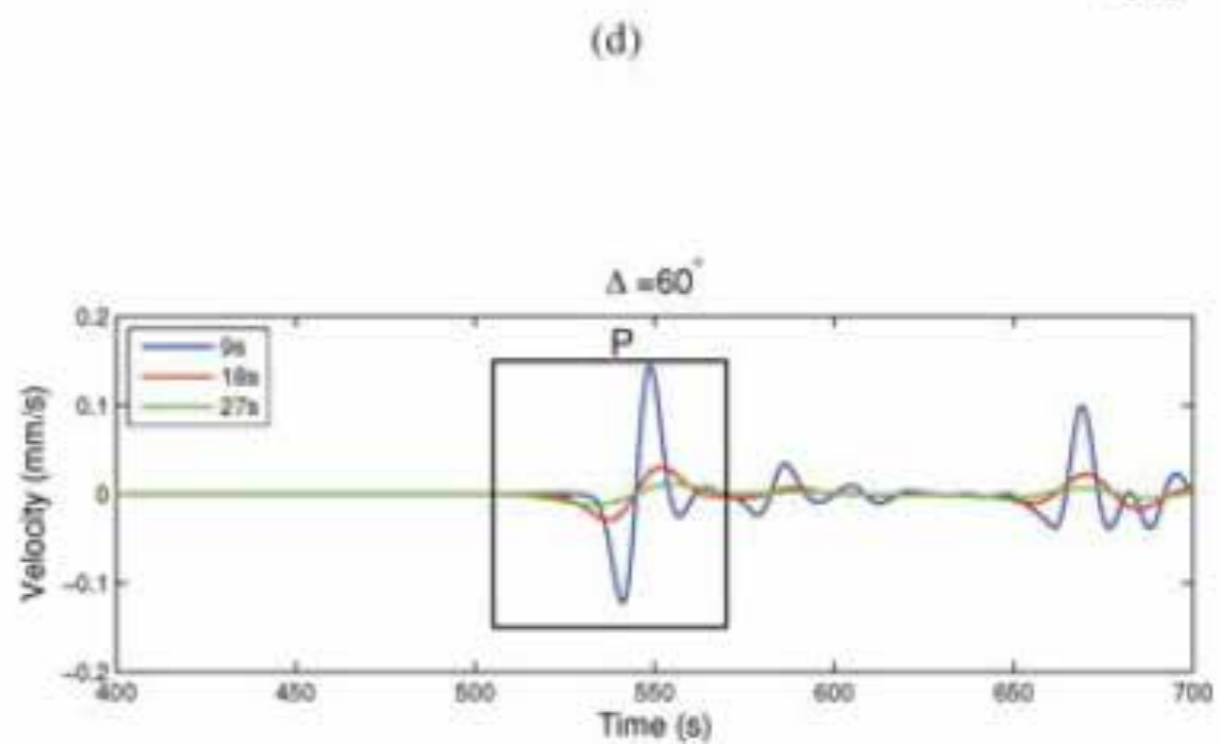
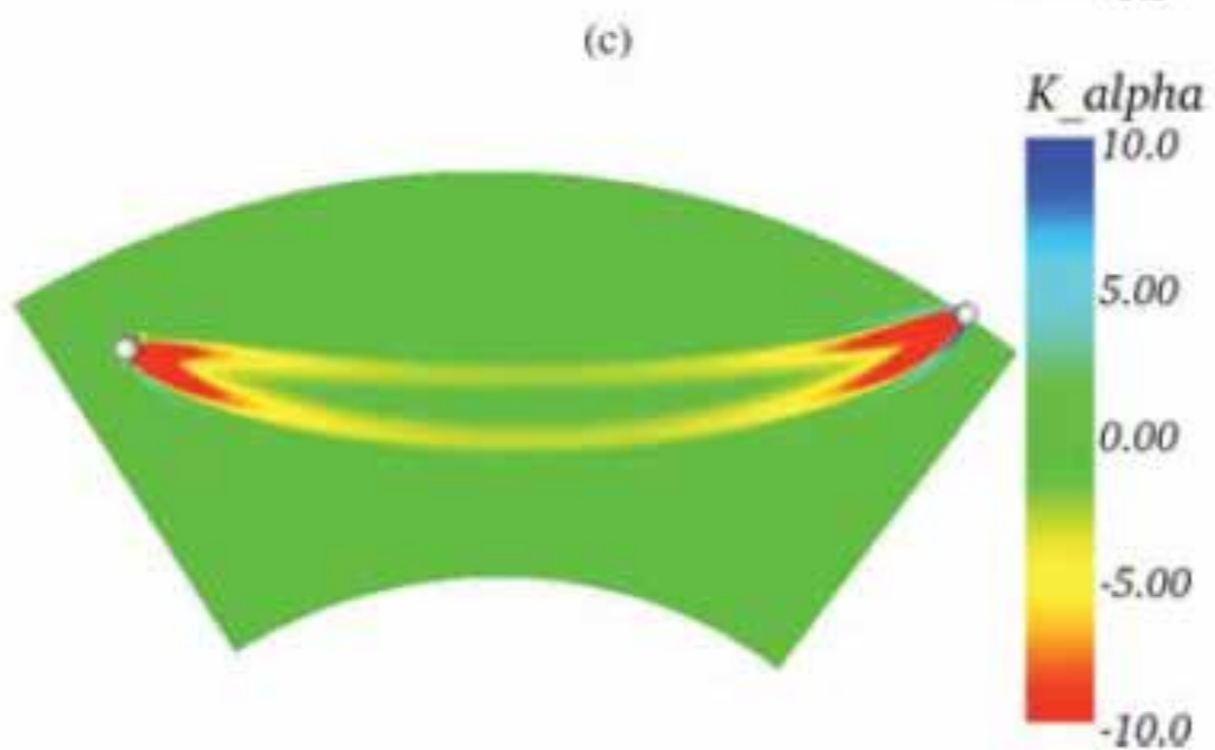
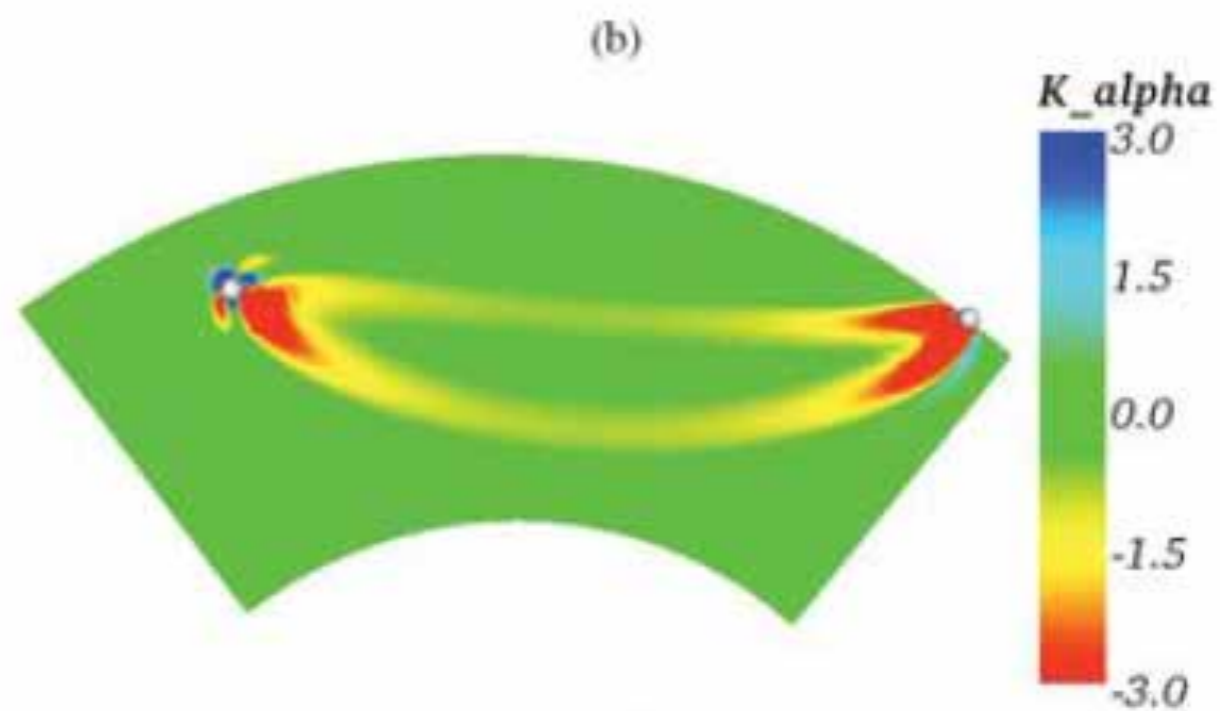
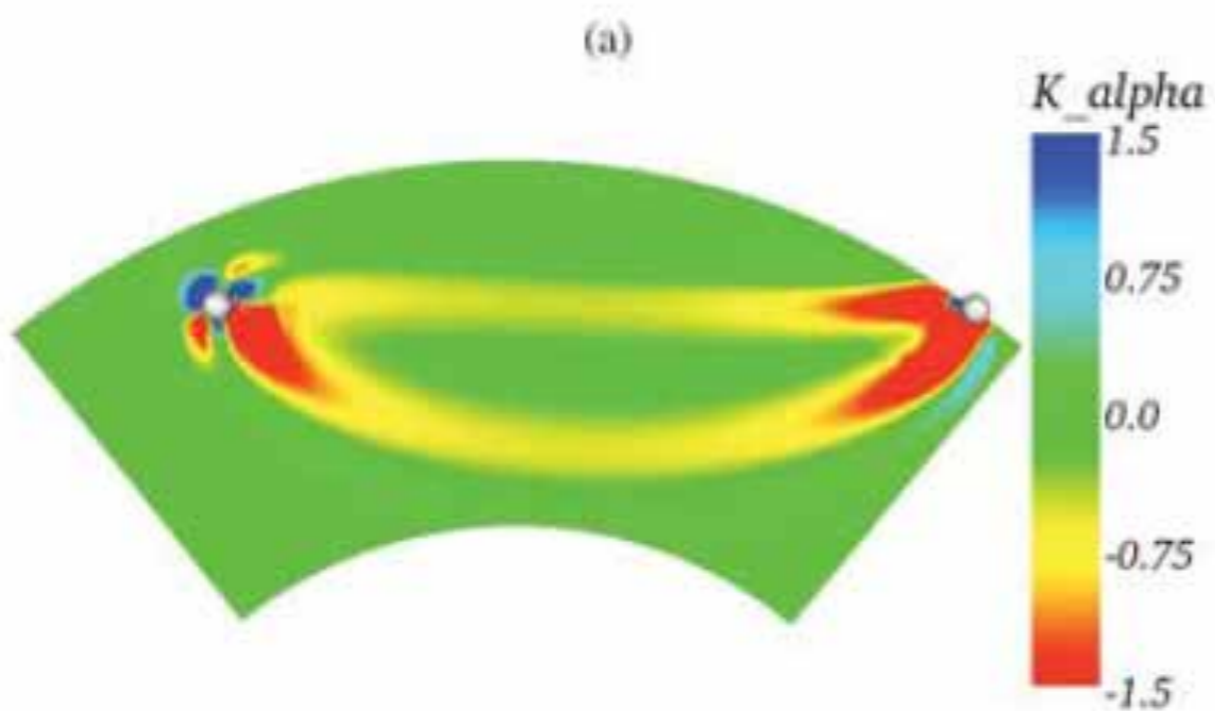


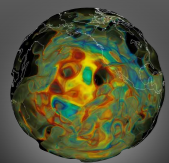
How to compute adjoint kernels?



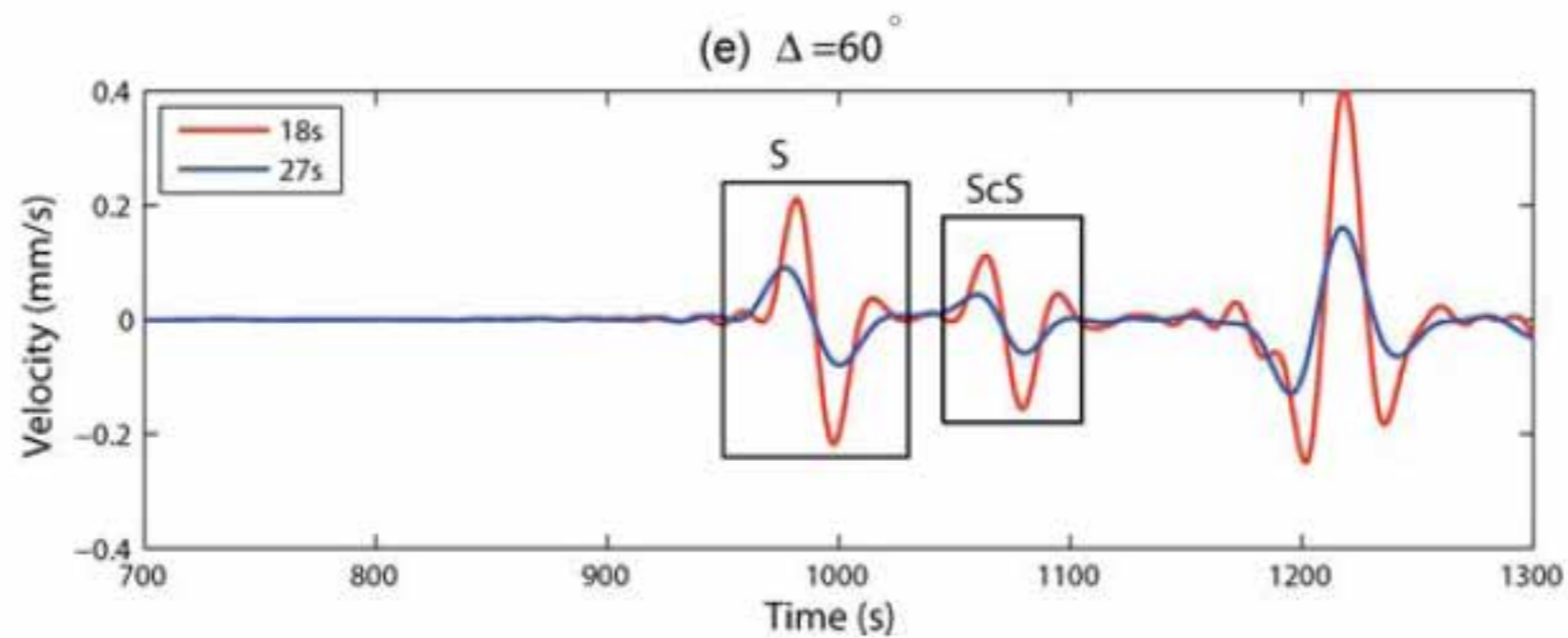
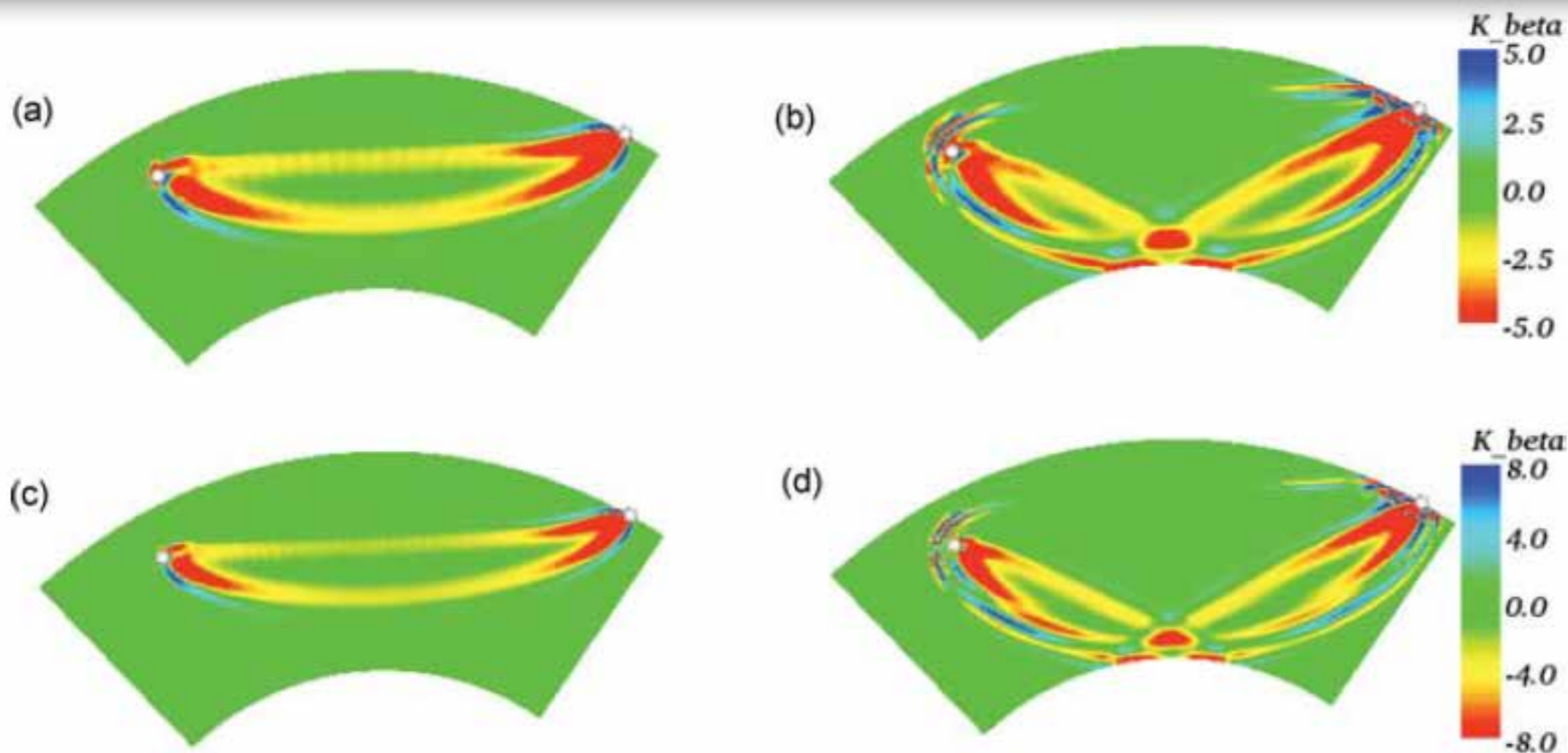


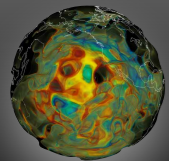
Banana-doughnuts



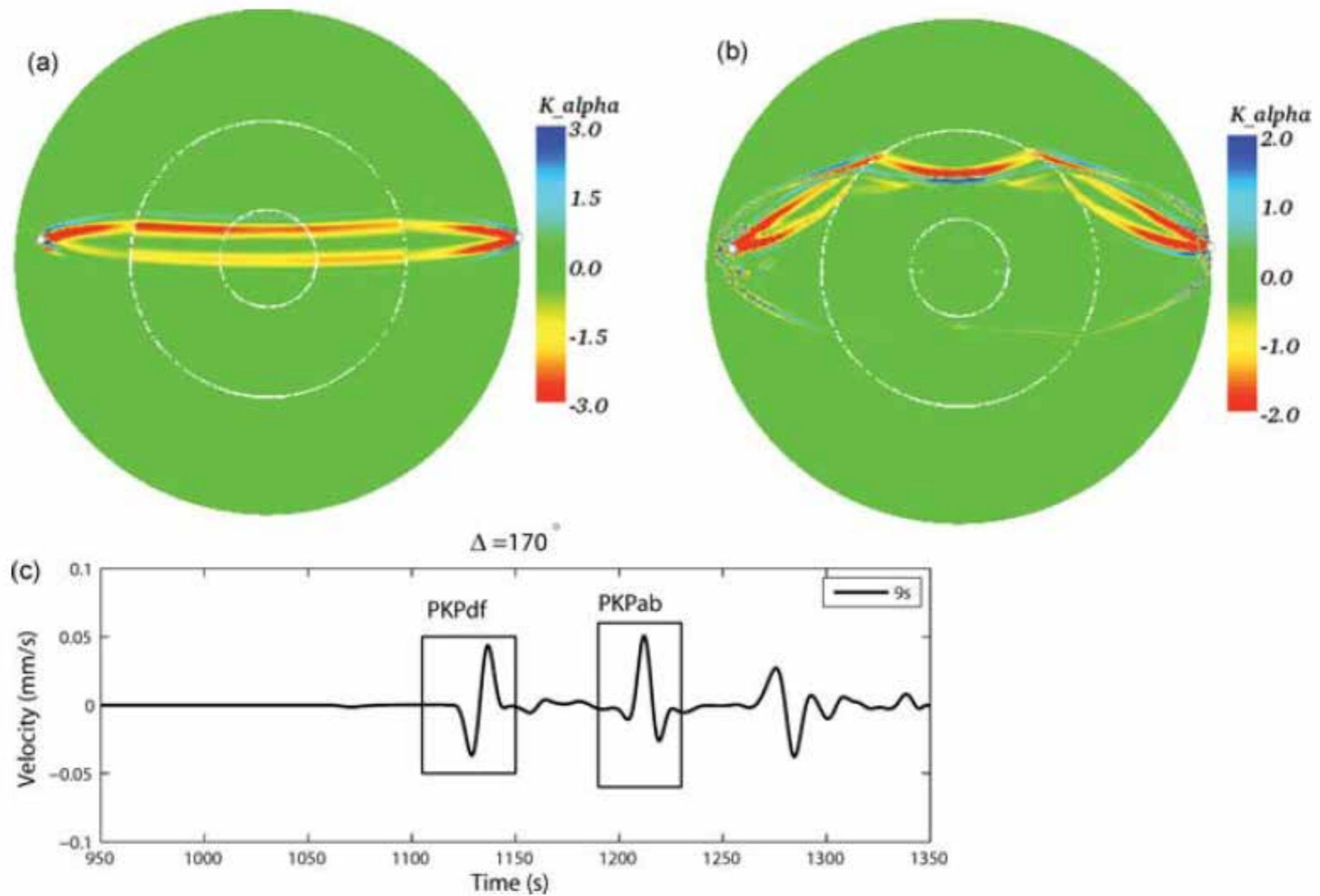


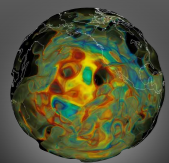
More banana-doughnuts



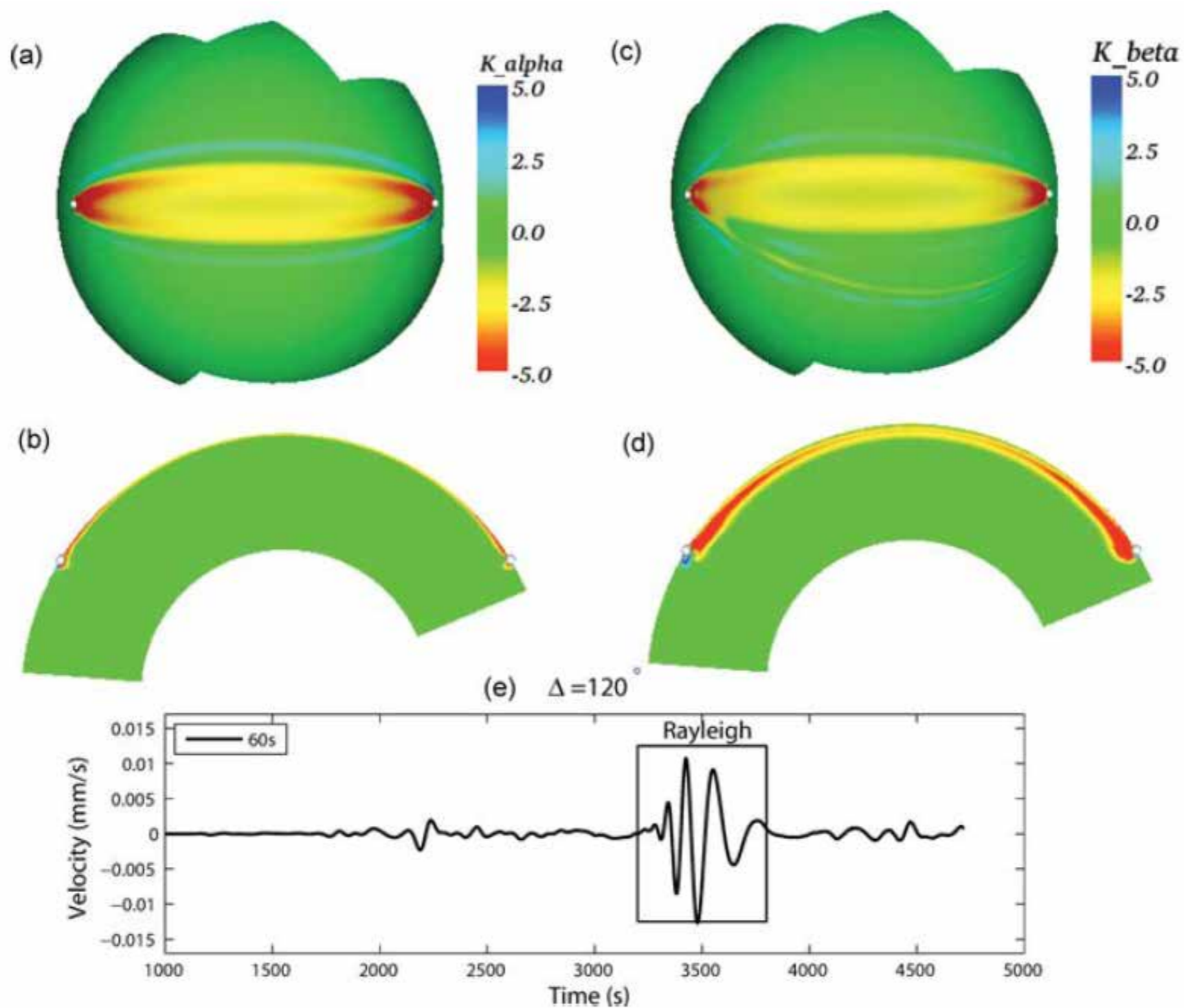


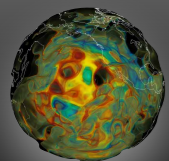
More banana-doughnuts



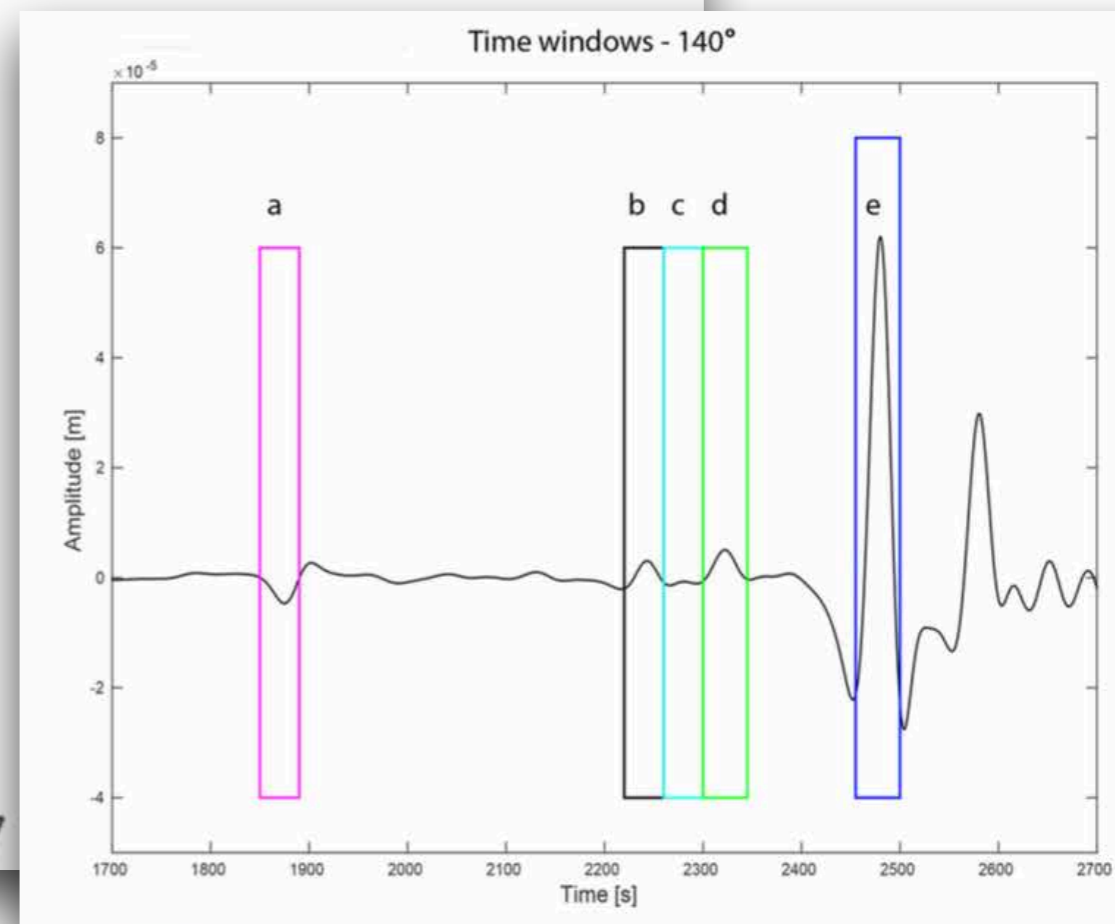
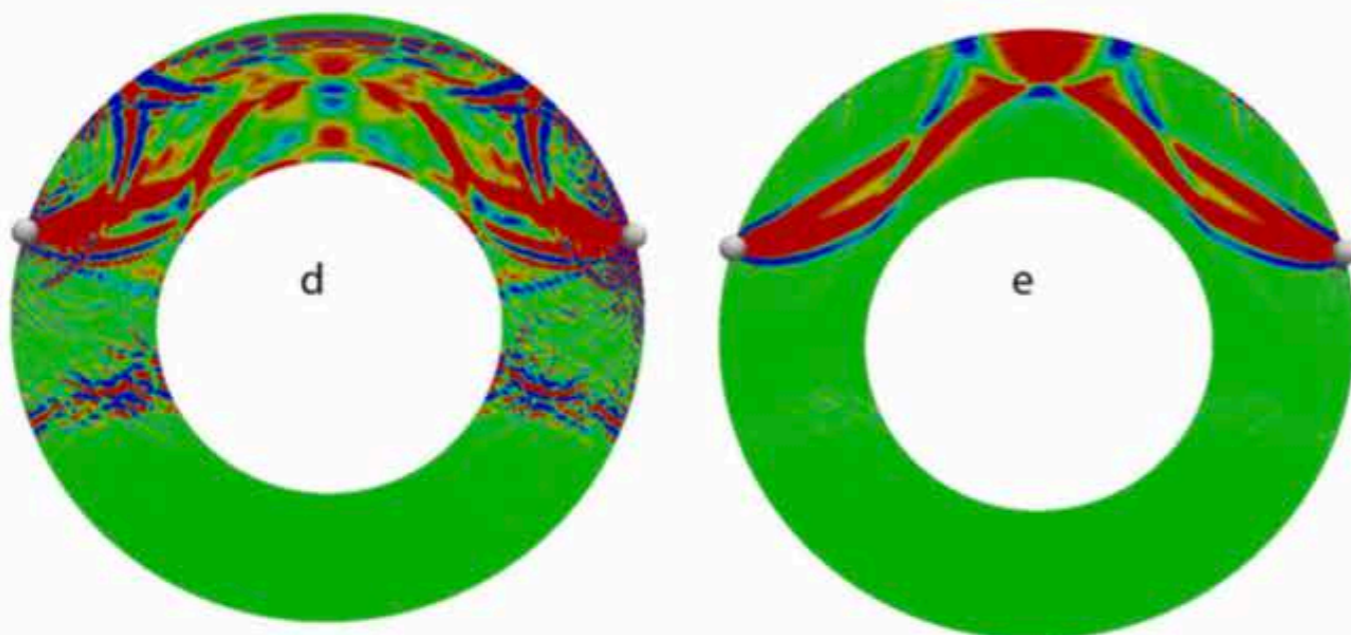
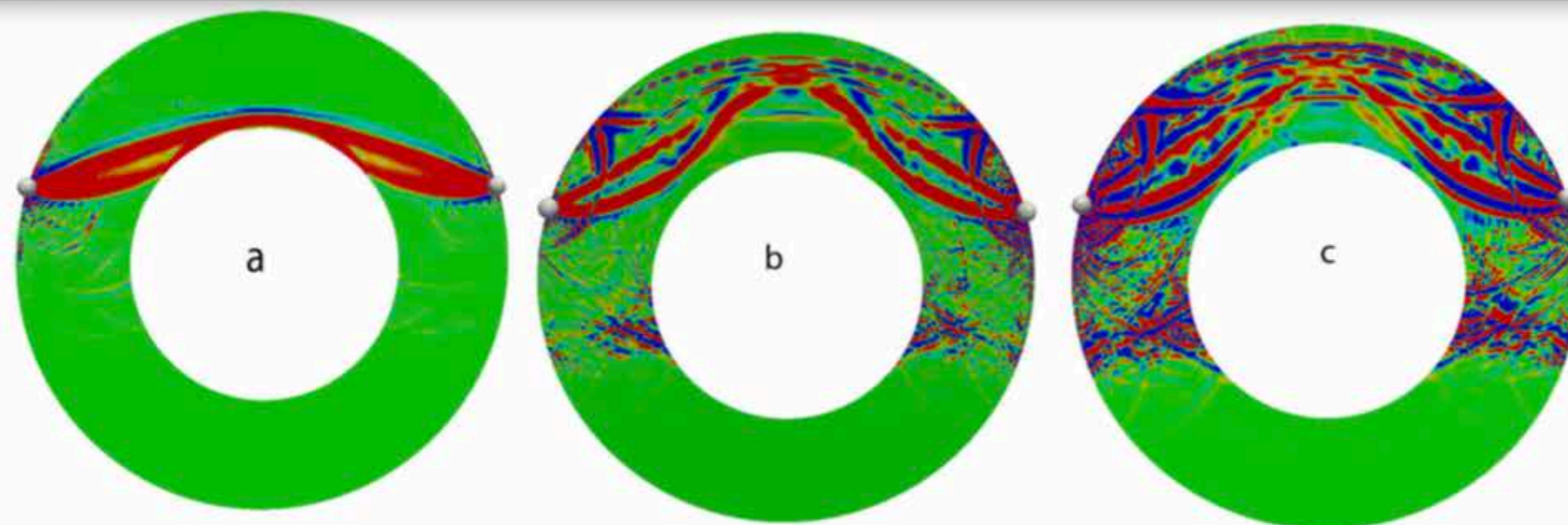


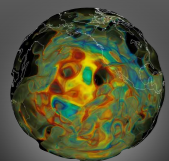
Surface wave sensitivities





Sensitivity to upper-mantle discontinuities





First global seismic models

JOURNAL OF GEOPHYSICAL RESEARCH, VOL. 89, NO. B7, PAGES 5953-5986, JULY 10, 1984

Mapping the Upper Mantle: Three-Dimensional Modeling of Earth Structure by Inversion of Seismic Waveforms

JOHN H. WOODHOUSE AND ADAM M. DZIEWONSKI

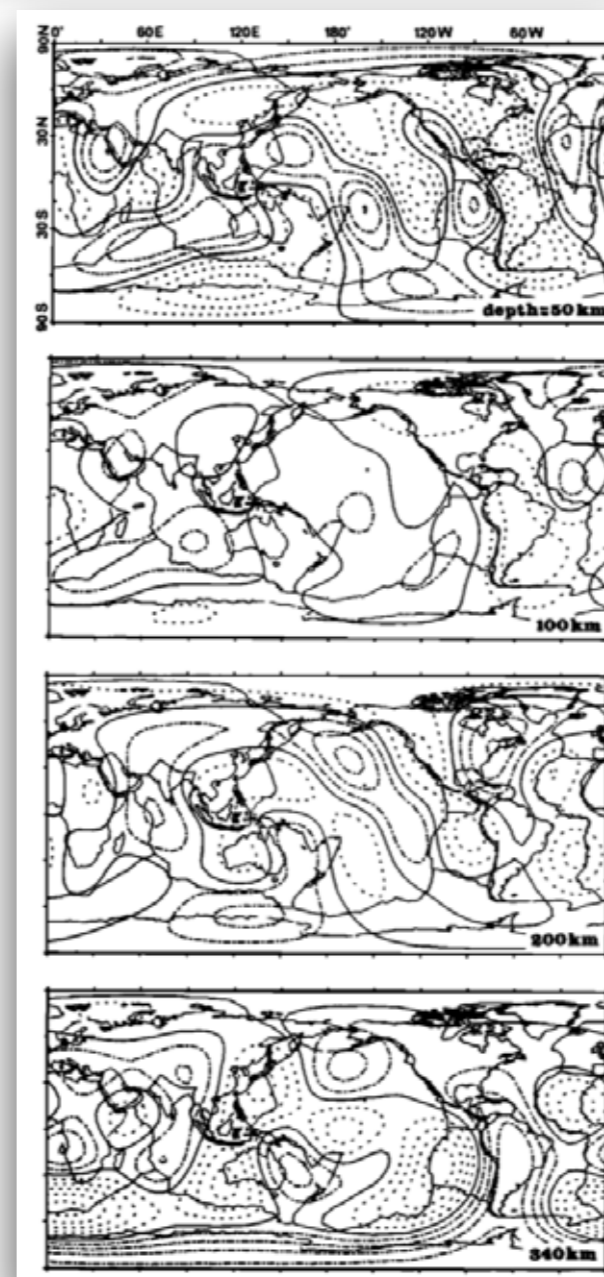
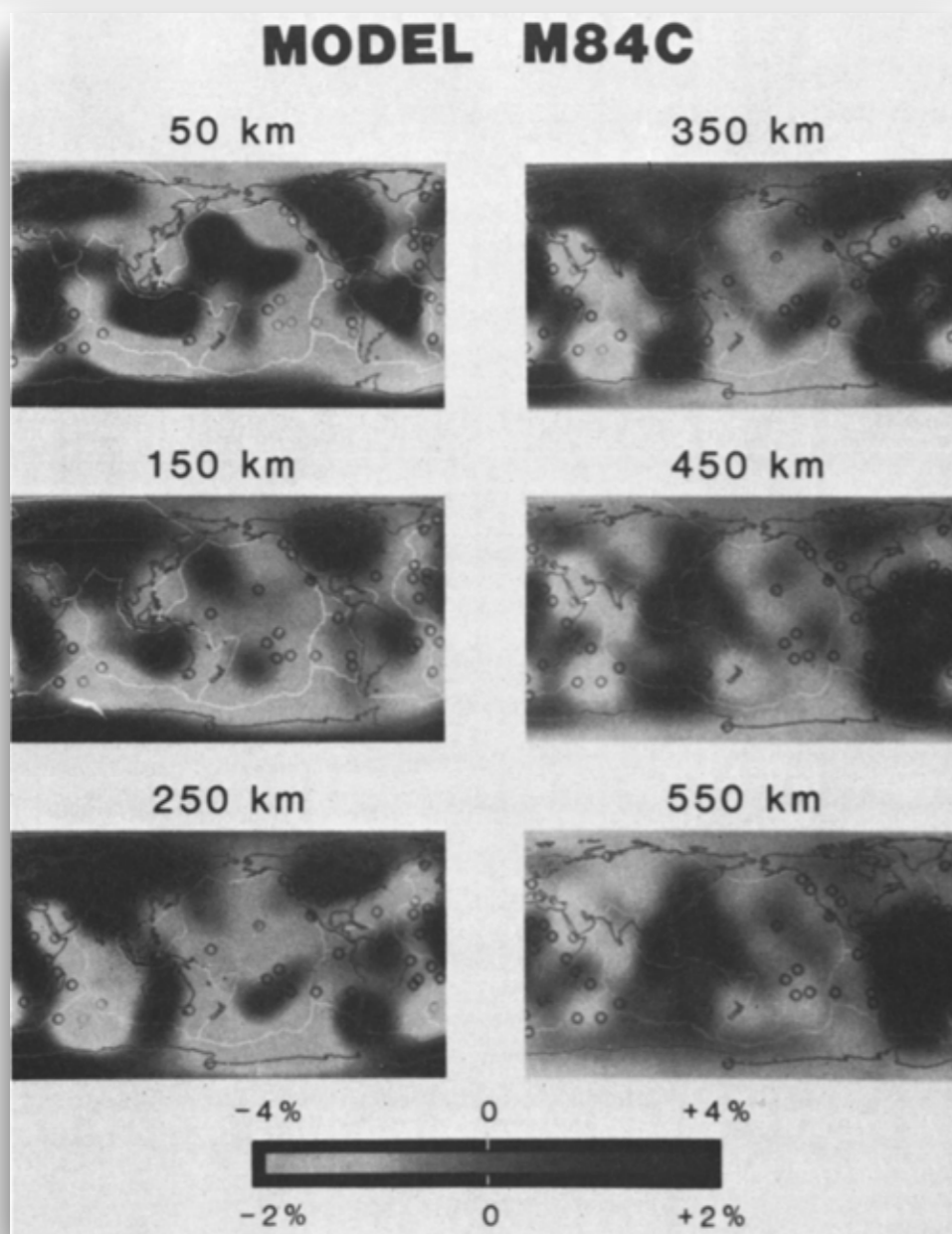
Department of Geological Sciences, Harvard University, Cambridge, Massachusetts

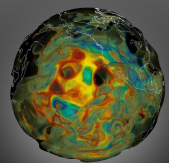
GEOPHYSICAL RESEARCH LETTERS, VOL. 11, NO. 2, PAGES 109-112, FEBRUARY 1984

ANISOTROPY AND SHEAR-VELOCITY HETEROGENEITIES IN THE UPPER MANTLE

H.-C. Nataf¹, I. Nakanishi², and Don L. Anderson

Seismological Laboratory, California Institute of Technology, Pasadena, Ca., 91125, USA.





First global seismic models

Global Images of the Earth's Interior

ADAM M. DZIEWONSKI AND JOHN H. WOODHOUSE

Science, New Series, Vol. 236, No. 4797 (Apr. 3, 1987), pp. 37-48

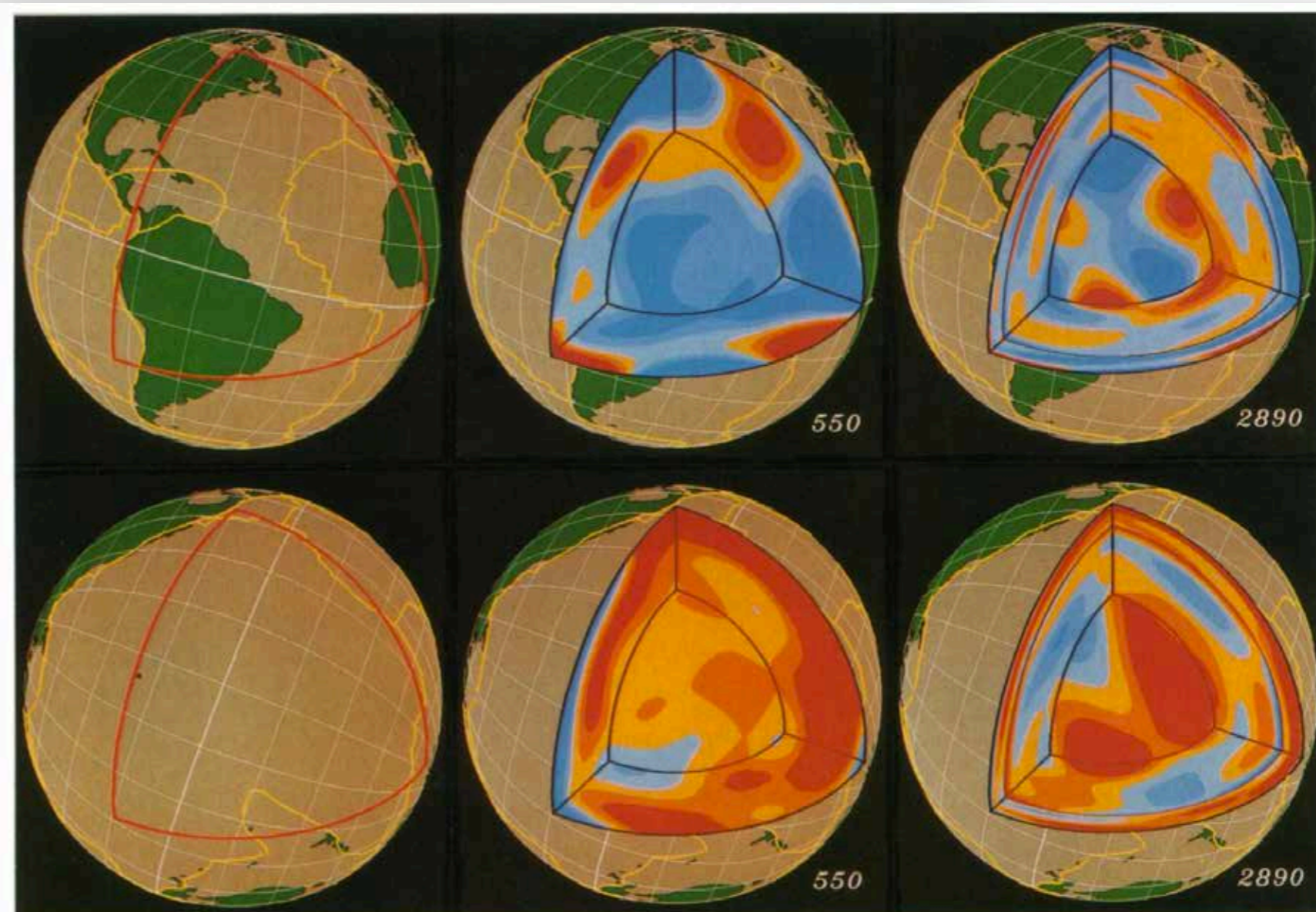
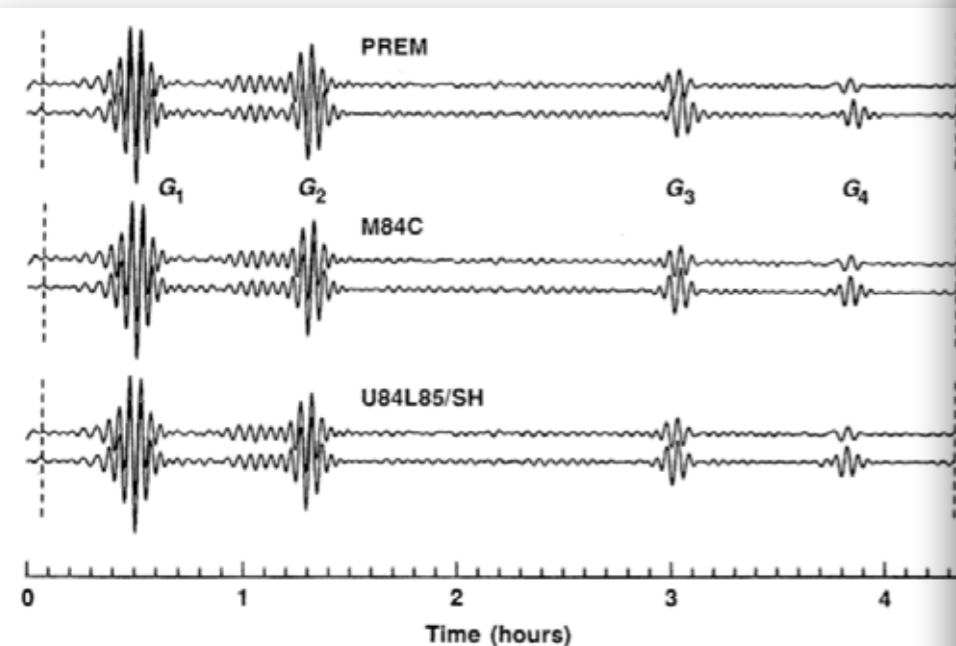
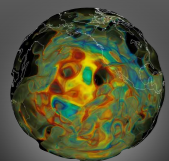


Fig. 12. Windows into the earth. Three-dimensional plots of velocity anomalies under two oceans viewed from an altitude of 35,000 km: the Atlantic (top panels) and the Pacific (bottom panels). Depth to the bottom is indicated in kilometers. In the upper mantle panels (550 km) the depth scale is exaggerated by a factor of 5. Based on models in (3) and (15). Yellow lines are plate boundaries.



Adjoint tomography

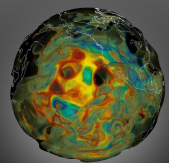
GEOPHYSICS, VOL. 49, NO. 8 (AUGUST 1984); P. 1259–1266.

Inversion of seismic reflection data in the acoustic approximation

Albert Tarantola*

ABSTRACT

The *nonlinear* inverse problem for seismic reflection data is solved in the acoustic approximation. The method is based on the generalized least-squares criterion, and it can handle errors in the data set and a priori information on the model. Multiply reflected energy is naturally taken into account, as well as refracted energy or surface waves. The inverse problem can be solved using an iterative algorithm which gives, at each iteration, updated values of bulk modulus, density, and time source function. Each step of the iterative algorithm essentially consists of a forward propagation of the actual sources in the current model and a forward propagation (backward in time) of the data residuals. The correlation at each point of the space of the two fields thus obtained yields the corrections of the bulk modulus and density models. This shows, in particular, that the general solution of the inverse problem can be attained by methods strongly related to the methods of migration of unstacked data, and commercially competitive with them.

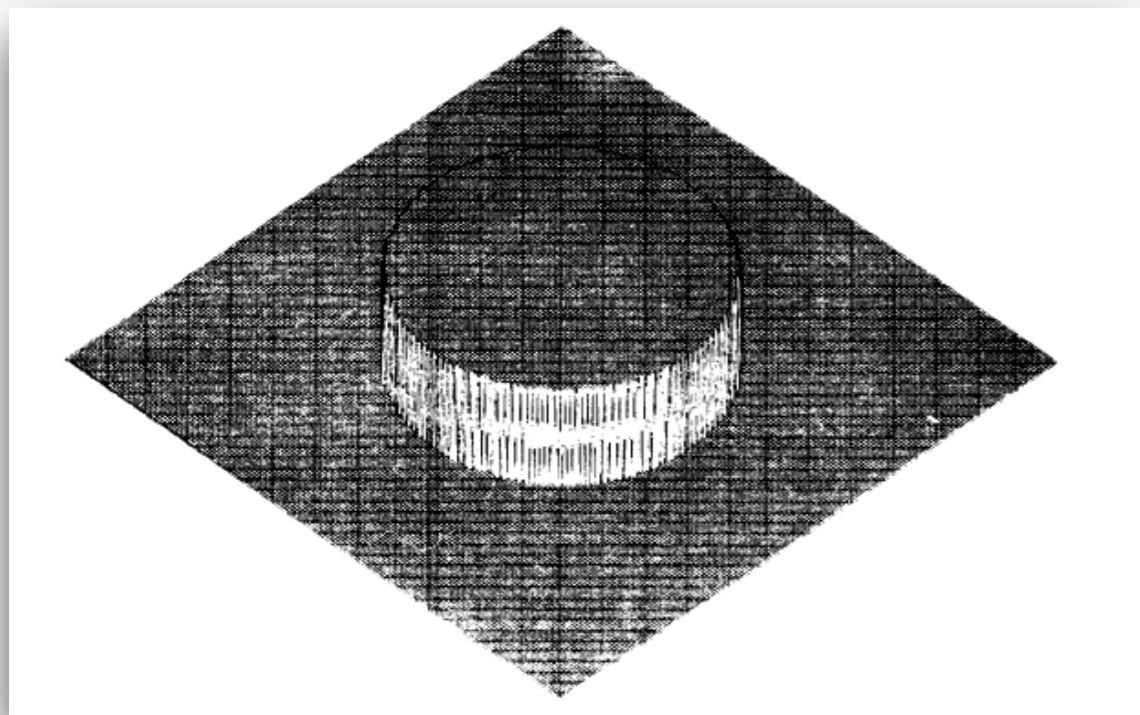


Adjoint tomography

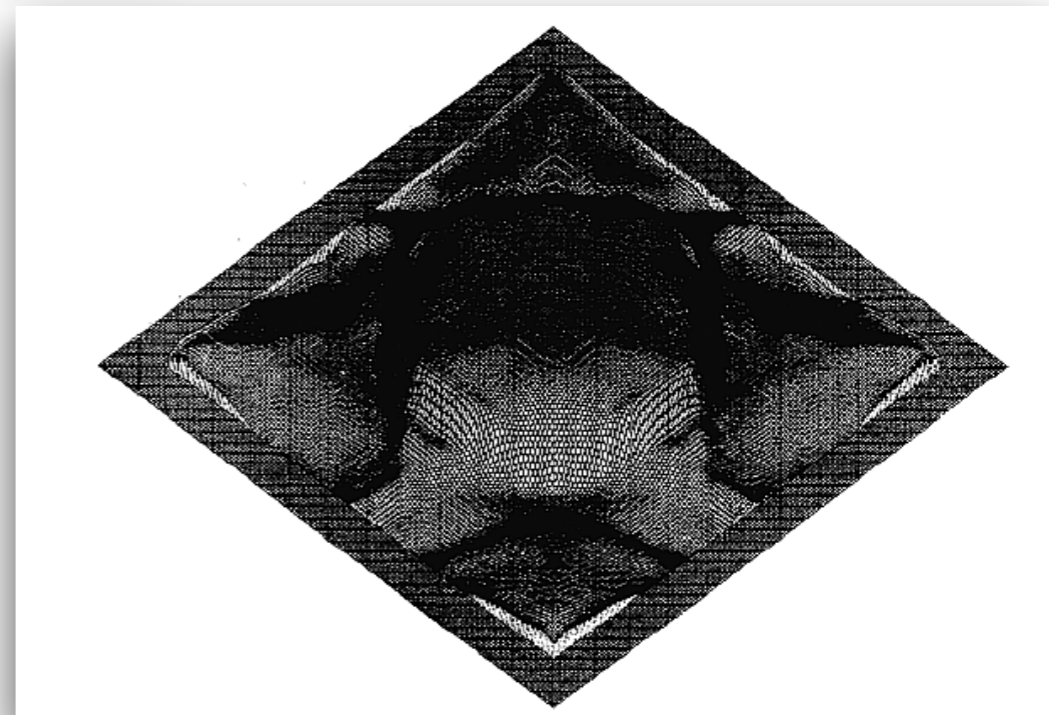
GEOPHYSICS, VOL. 51, NO. 7 (JULY 1986); P. 1387-1403, 16 FIGS.

Two-dimensional nonlinear inversion of seismic waveforms: Numerical results

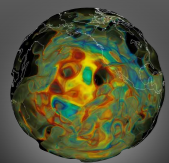
Odile Gauthier*, Jean Virieux*, and Albert Tarantola*



target model



after 5 iterations



Adjoint tomography

GEOPHYSICS, VOL. 51, NO. 7 (JULY 1986); P. 1387-1403, 16 FIGS.

Two-dimensional nonlinear inversion of seismic waveforms: Numerical results

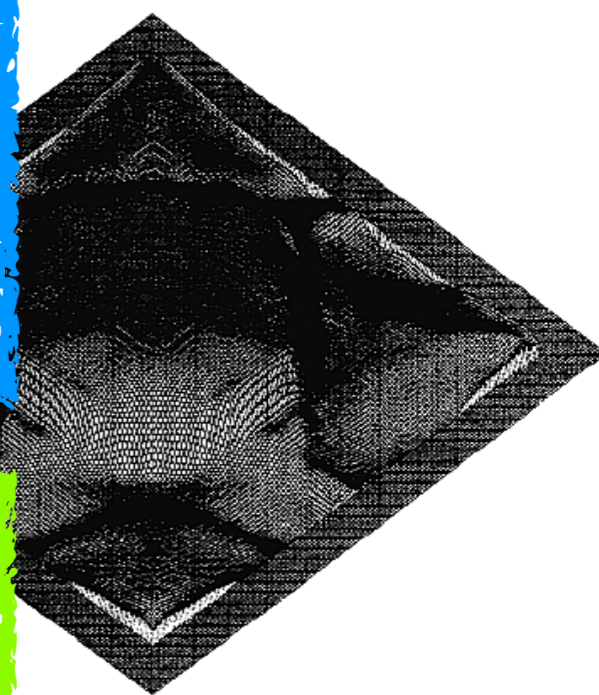
Odile Gauthier*, Jean Virieux* and Albert Tarantola*

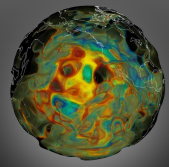
On the other hand, numerical methods based on discretization of the wave equation synthesize the different waves arriving at the station as a whole. This can make physical interpretation of a forward calculation difficult, but it does not prevent inversion. The main limitation of numerical methods is their computational cost, in both time and memory requirements.

Future interpretation of seismic records will certainly be performed through very accurate forward modeling. The method described here shows that accurate forward modeling codes can be used for waveform matching.

target model

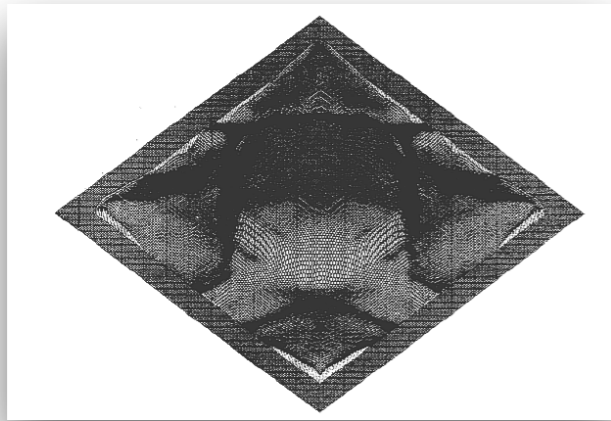
after 5 iterations





Global adjoint tomography

Gauthier et al. 1987



● $dx=dz=5$ m

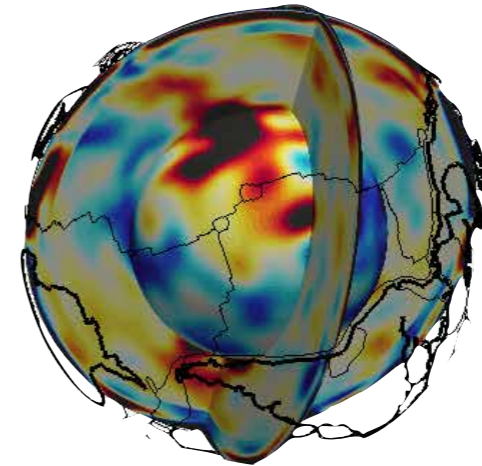
● 200x200 grid points

● 8 shots, 400 receivers

● 2D acoustic

● ~8 s seismograms

Bozdağ et al. 2016



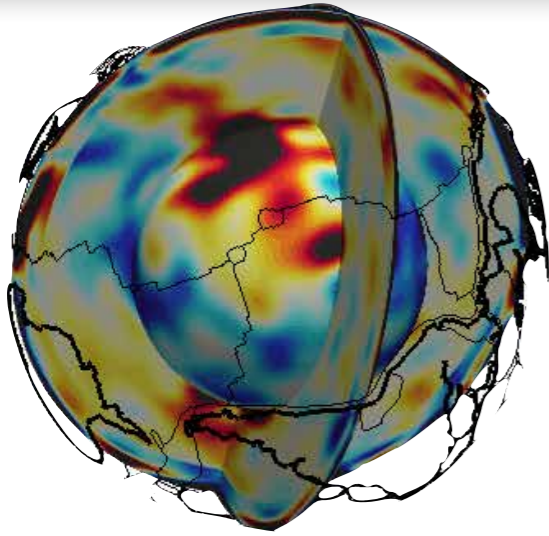
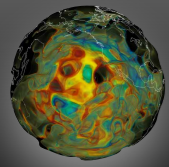
● entire globe

● ~300 millions grid points

● 253 equakes, ~1000 stations/event
(>6000 earthquakes)

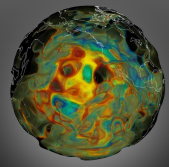
● 3D anelastic

● 180 m seismograms

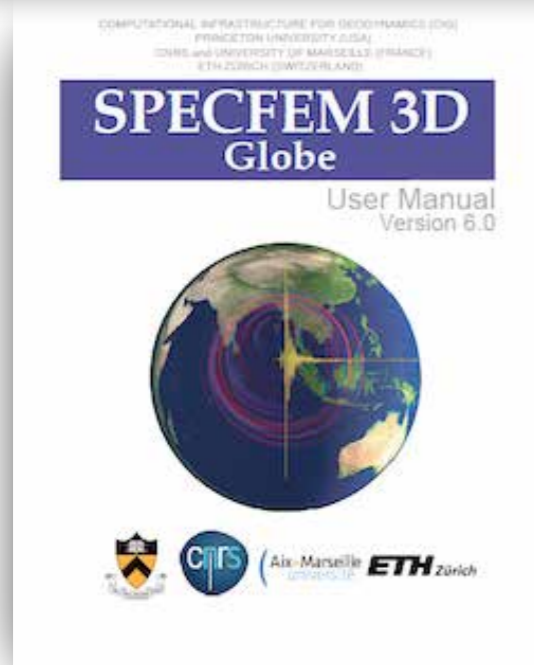


E. Bozdag, D. Peter, M. Lefebvre, D. Komatitsch, J. Tromp, J. Hill, N. Podhorszki, D. Pugmire, 2016, GJI

- **First-generation global adjoint model: TI in upper mantle, *P* & *S* wavespeeds**
- *Forward simulations & Fréchet sensitivity kernels in realistic 3D Earth models using spectral-element & adjoint methods.*
- **253 earthquakes and 15 conjugate-gradient iterations**
- **Crust and mantle are inverted simultaneously: *no crustal corrections!***



Numerical simulations

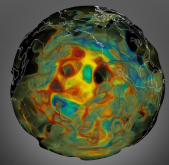


- **SPECIFEM3D_GLOBE**
(Komatitsch & Tromp 2002)
GPU version by Daniel Peter

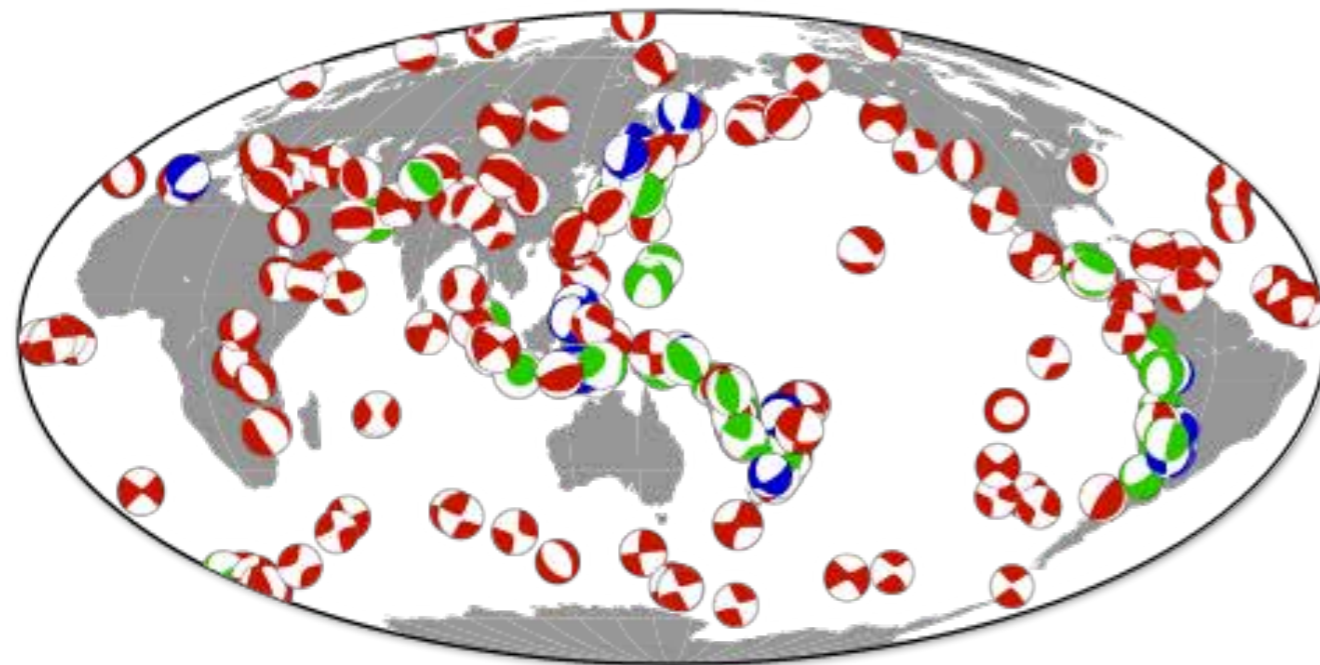
- **3D reference model:**

S362ANI (Kustowski et al. 2008) + **Crust2.0** (Bassin et al. 2000)

- **Topography/bathymetry/attenuation/ellipticity/rotation/gravity (Cowling approx.)/ocean load**
- **Length of seismograms = 100 m-200 m,**
Tmin = ~27 s, now Tmin=~17 s, soon Tmin=~9 s



Earthquakes & seismic stations

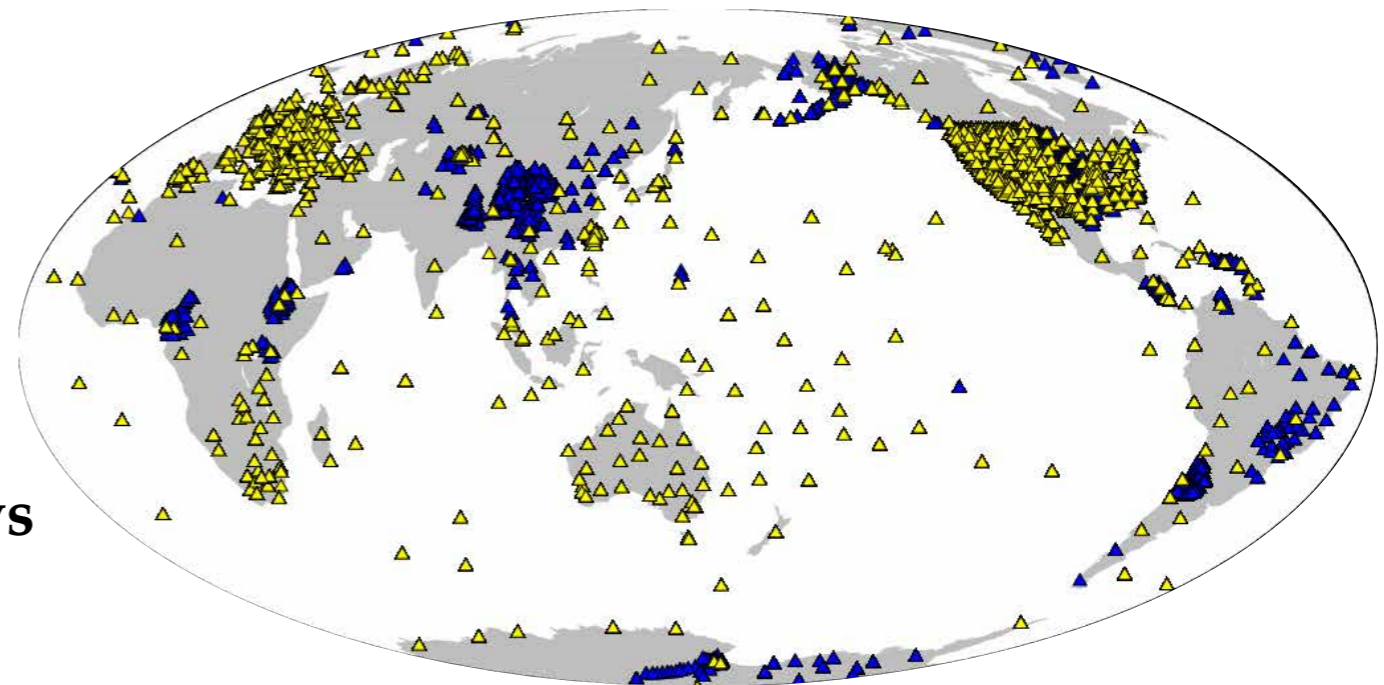


Initially selected 253 CMT events
($5.8 \leq M_w \leq 7.0$)

shallow: $d \leq 50$ km

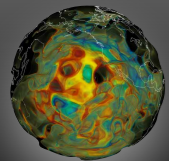
intermediate: $50 \text{ km} < d \leq 300$ km

deep: $d > 300$ km



data from IRIS & ORFEUS

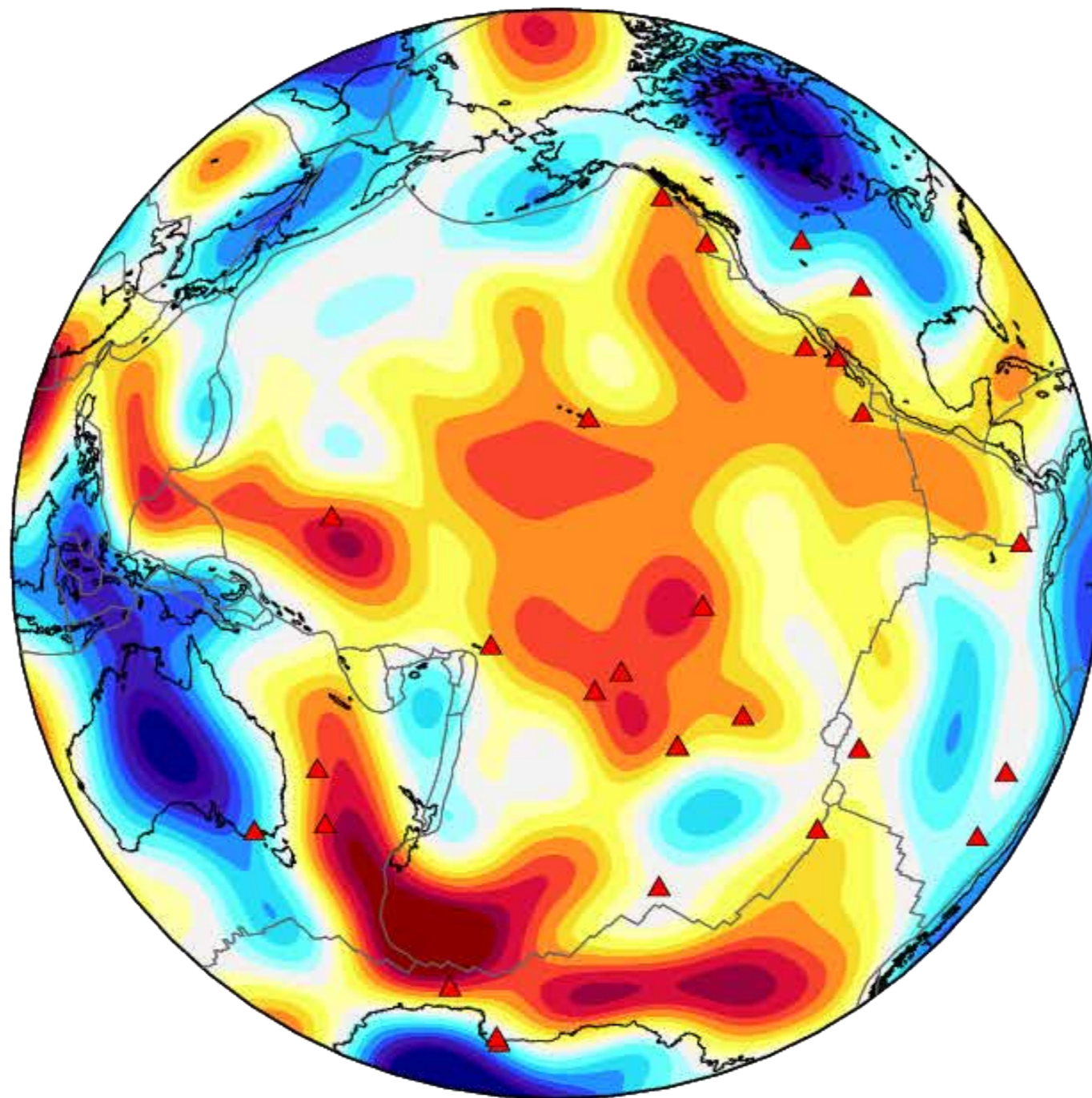
blue: additional local + PASSCAL arrays



Cross-sections - 250 km

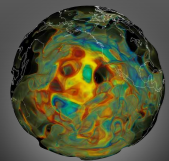
M00: S362ANI

(Kustowski et al. 2007)



$d\ln V_{sv}$ (3%)

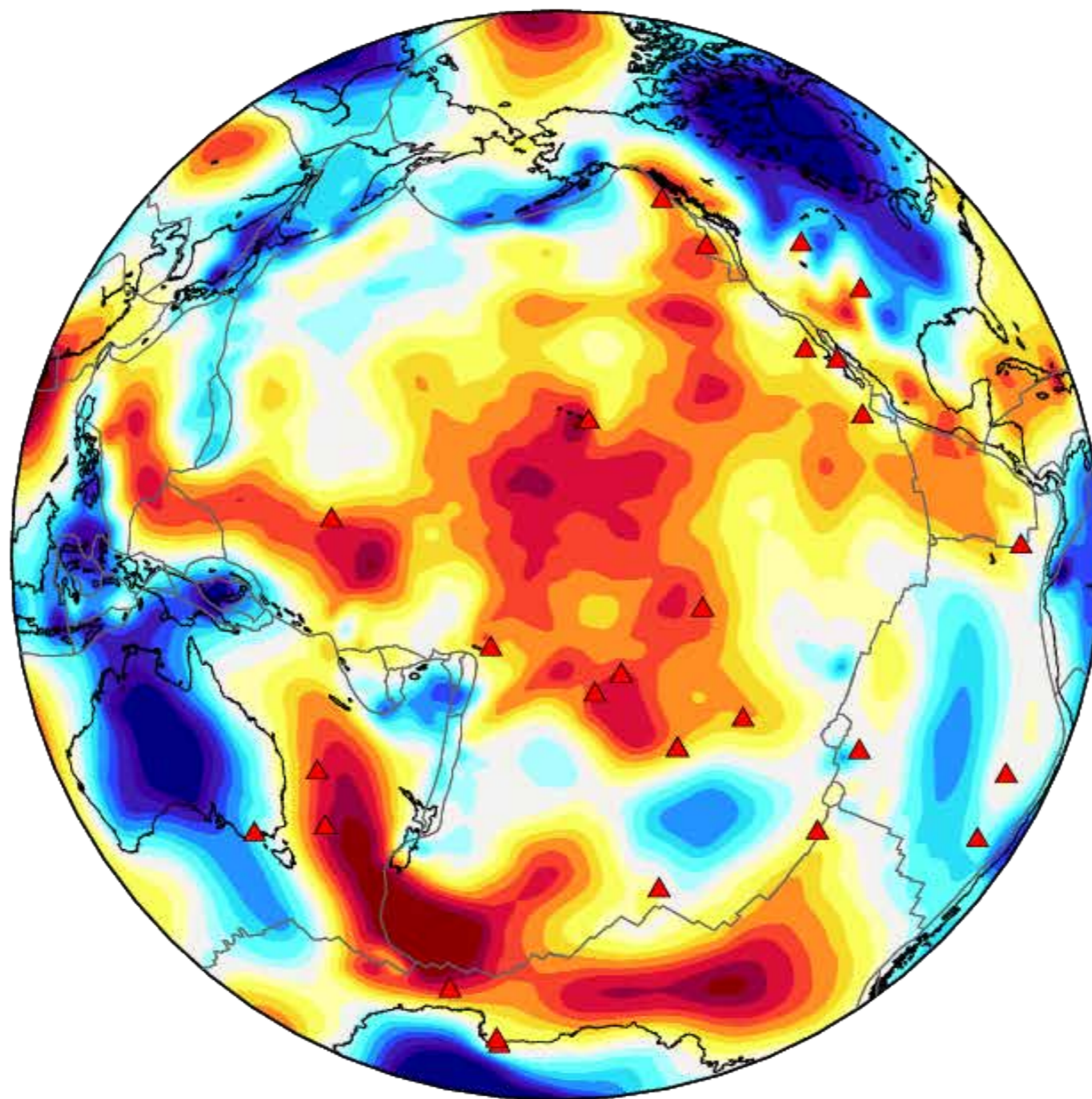
min  max



Cross-sections - 250 km

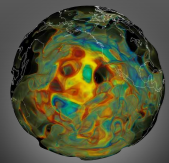
M15: GLAD-M15

(Bozdag et al. 2016)



$d \ln V_{sv}$ (3%)

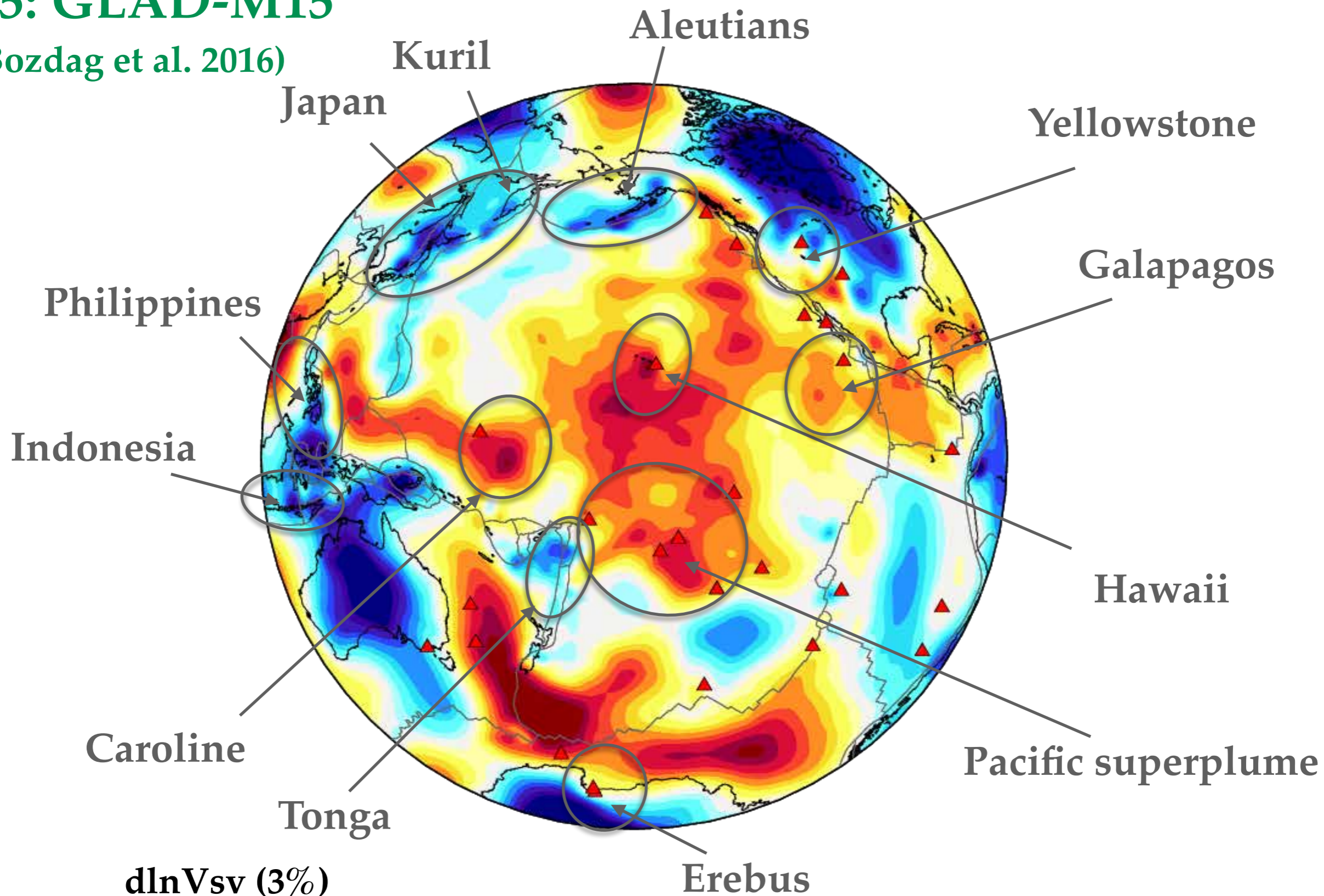
min  max



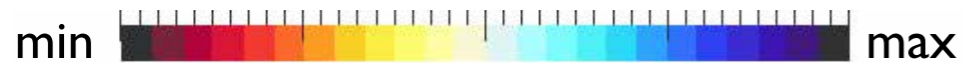
Cross-sections - 250 km

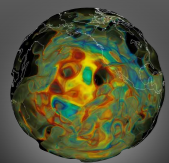
M15: GLAD-M15

(Bozdag et al. 2016)



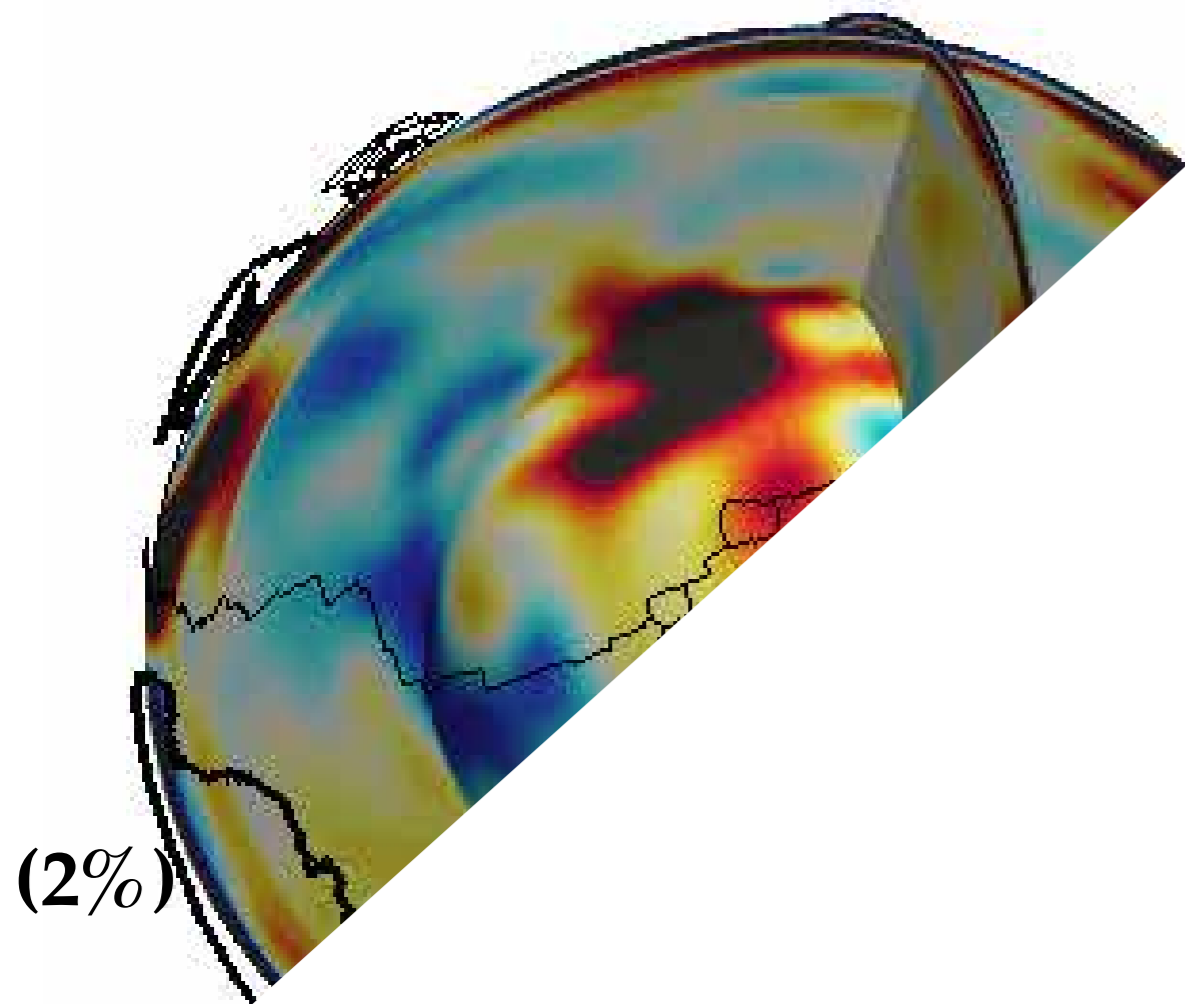
$d \ln V_{sv}$ (3%)





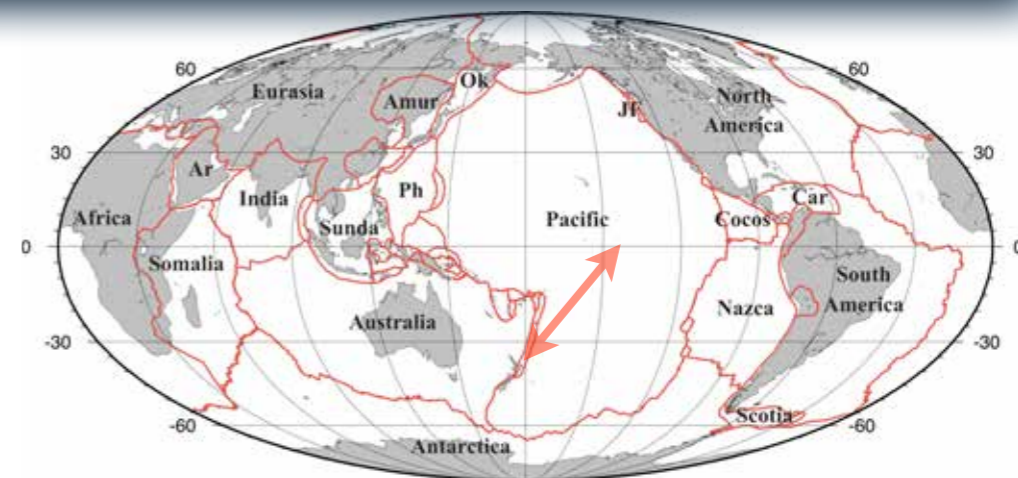
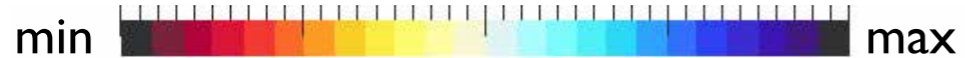
More sections

M00 - 1DREF

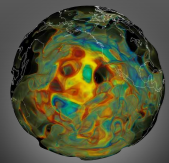


(2%)

$d \ln V_{sv}$



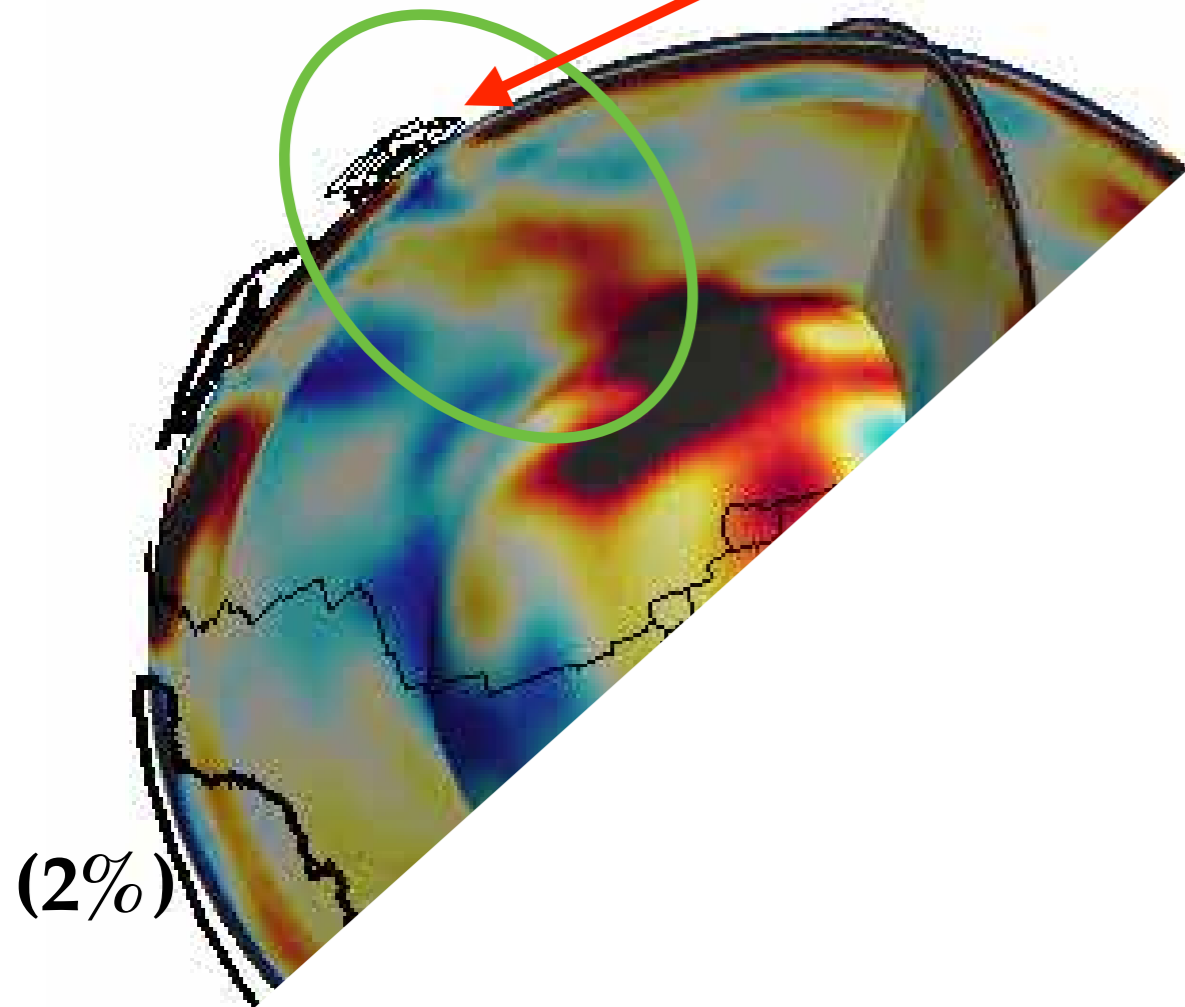
GMT 2005 Apr 29 16:24:38 Plate Boundaries from PS2002 (Peter Bird) Dataset



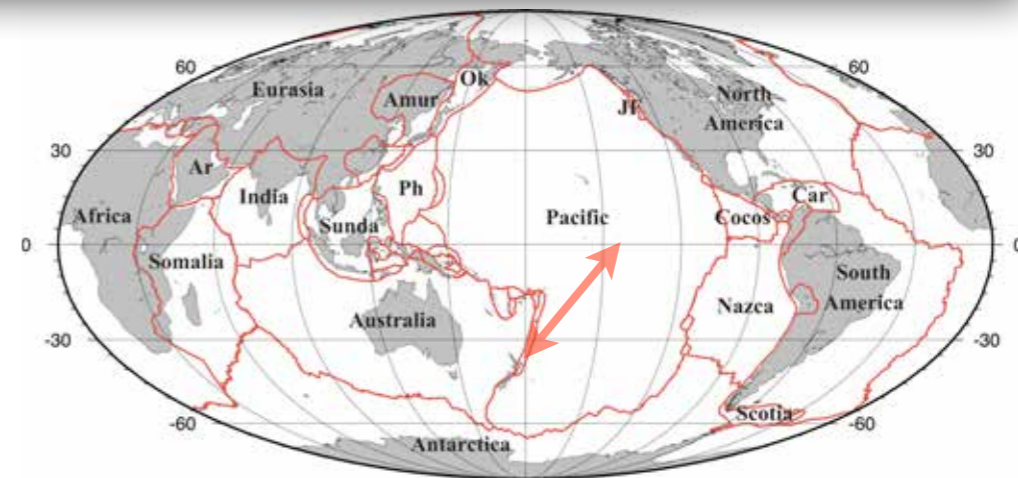
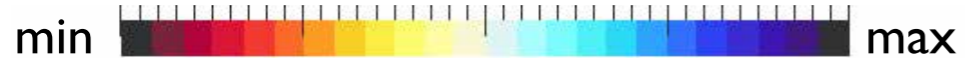
More sections

M15 - 1DREF

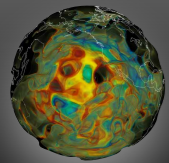
Samoa



$d \ln V_{sv}$



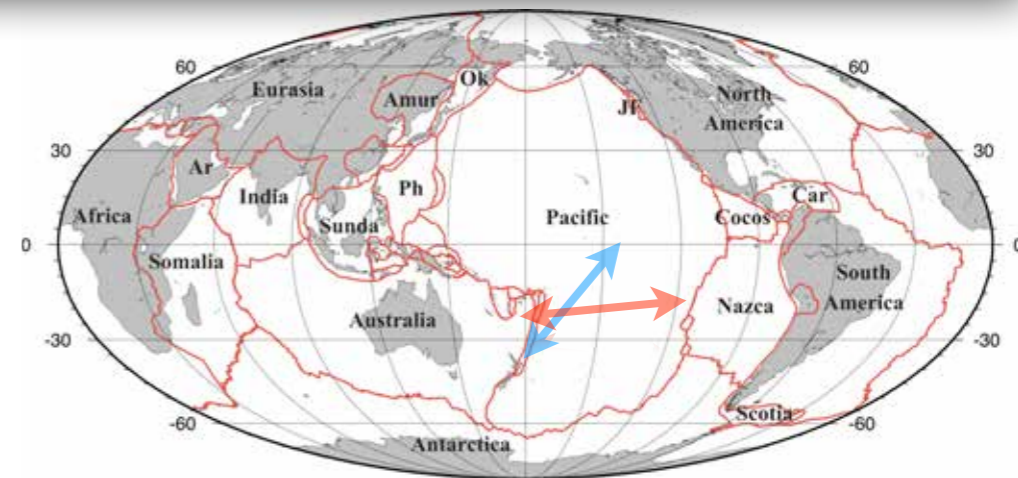
GMT 2005 Apr 29 16:24:38 Plate Boundaries from PB2002 (Peter Bird) Dataset



More sections

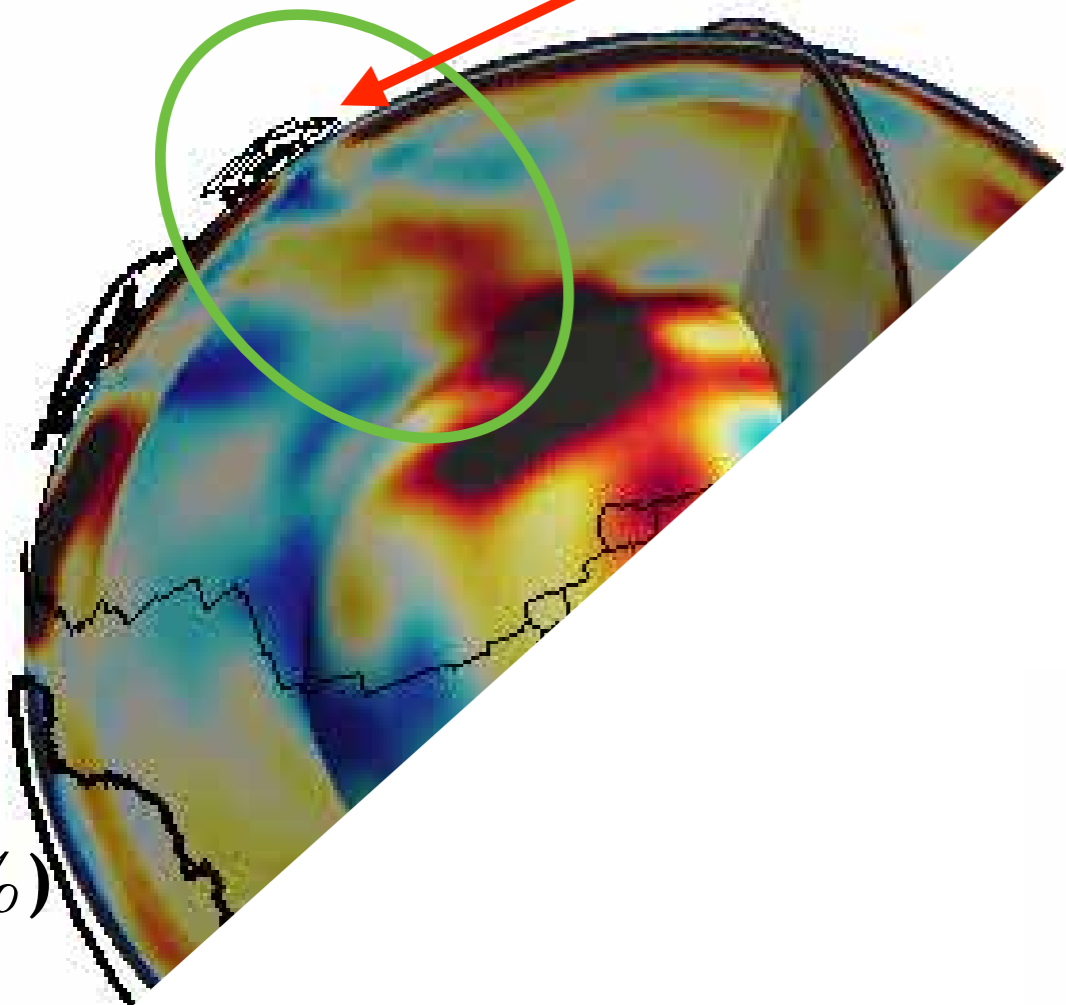
M15 - 1DREF

Samoa

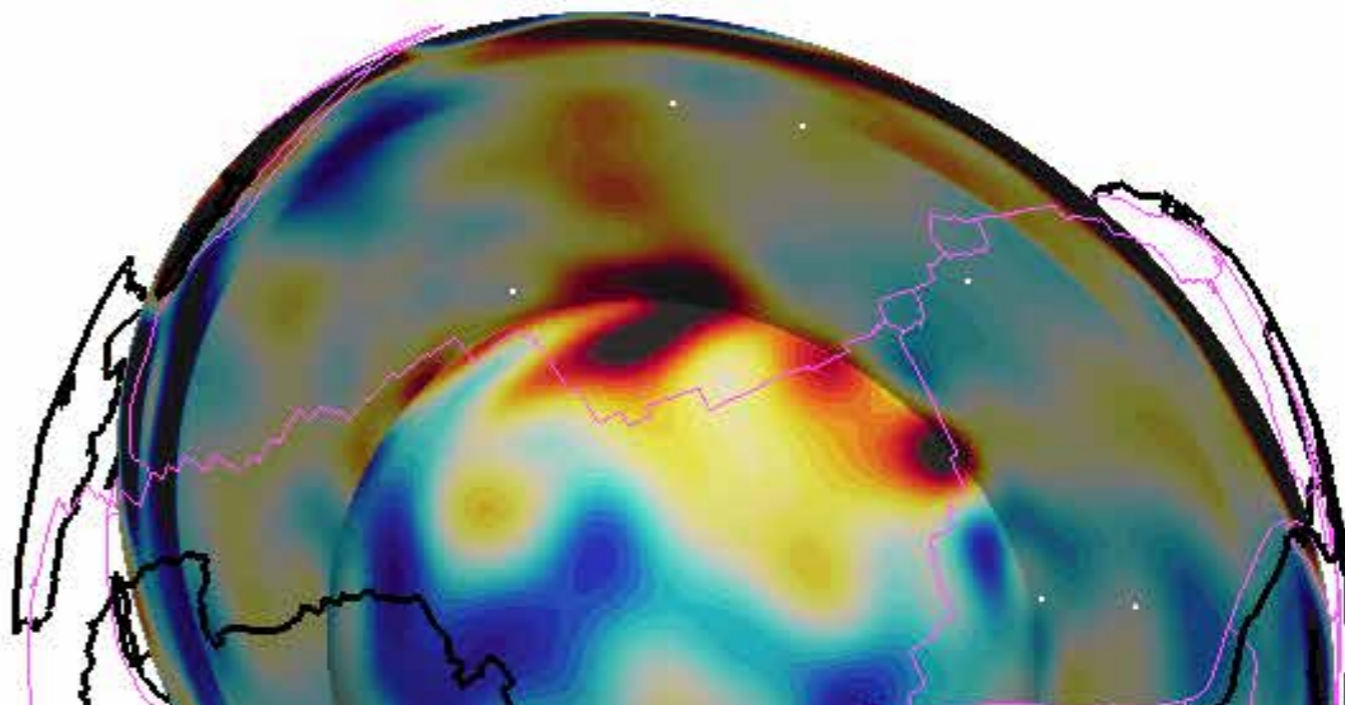


GMT 2005 Apr 29 16:24:38 Plate Boundaries from PB2002 (Peter Bird) Dataset

(2%)

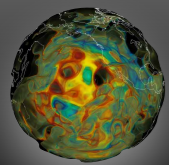


(2%)



$d \ln V_{sv}$

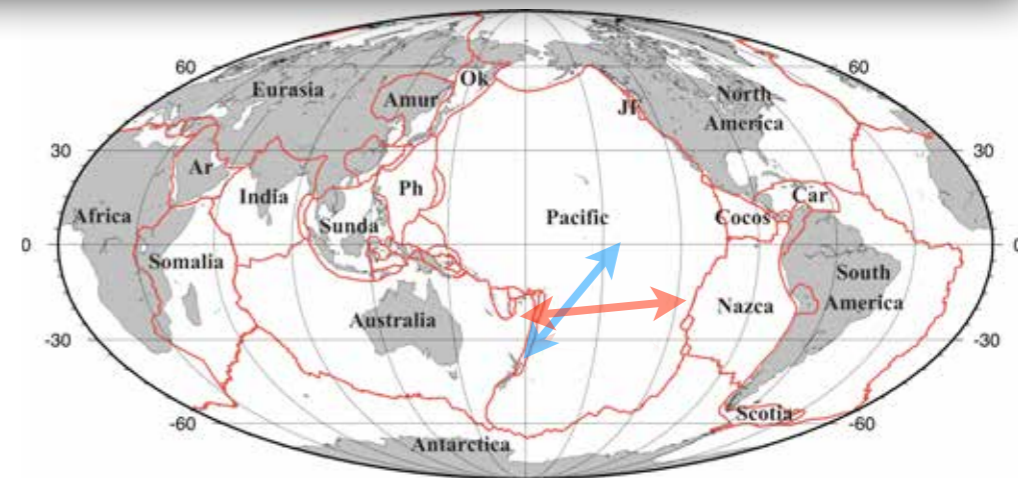




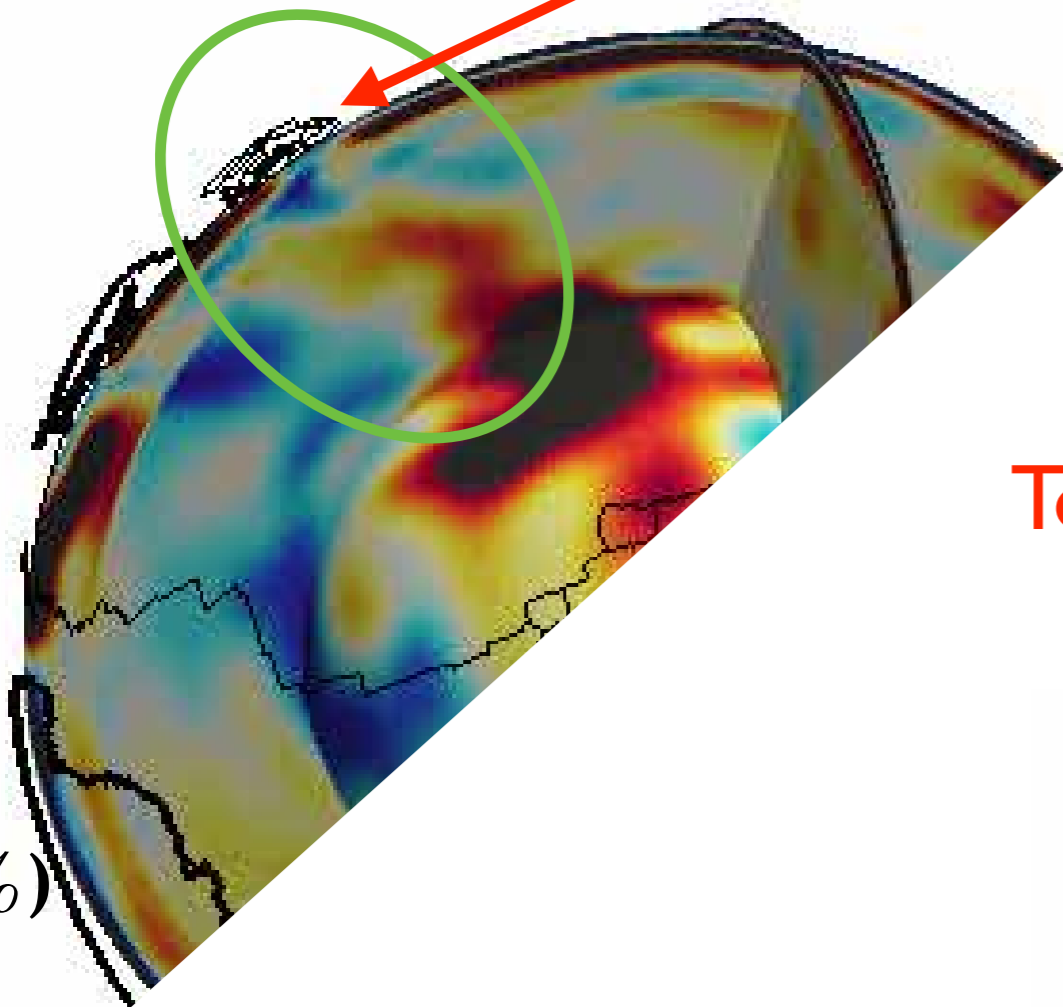
More sections

M15 - 1DREF

Samoa



(2%)

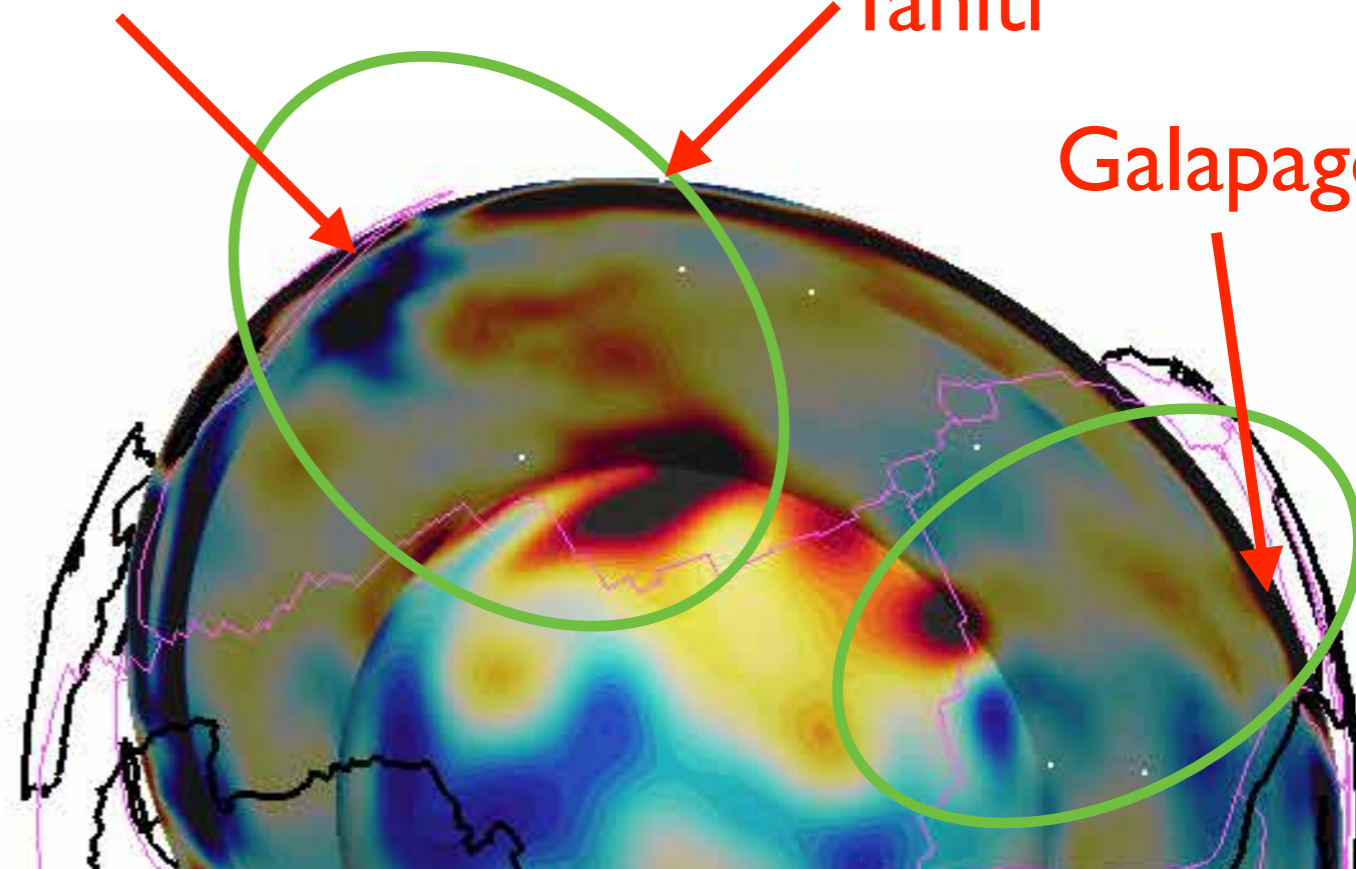


Tonga

Tahiti

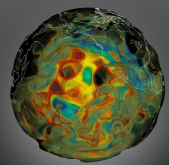
Galapagos

(2%)



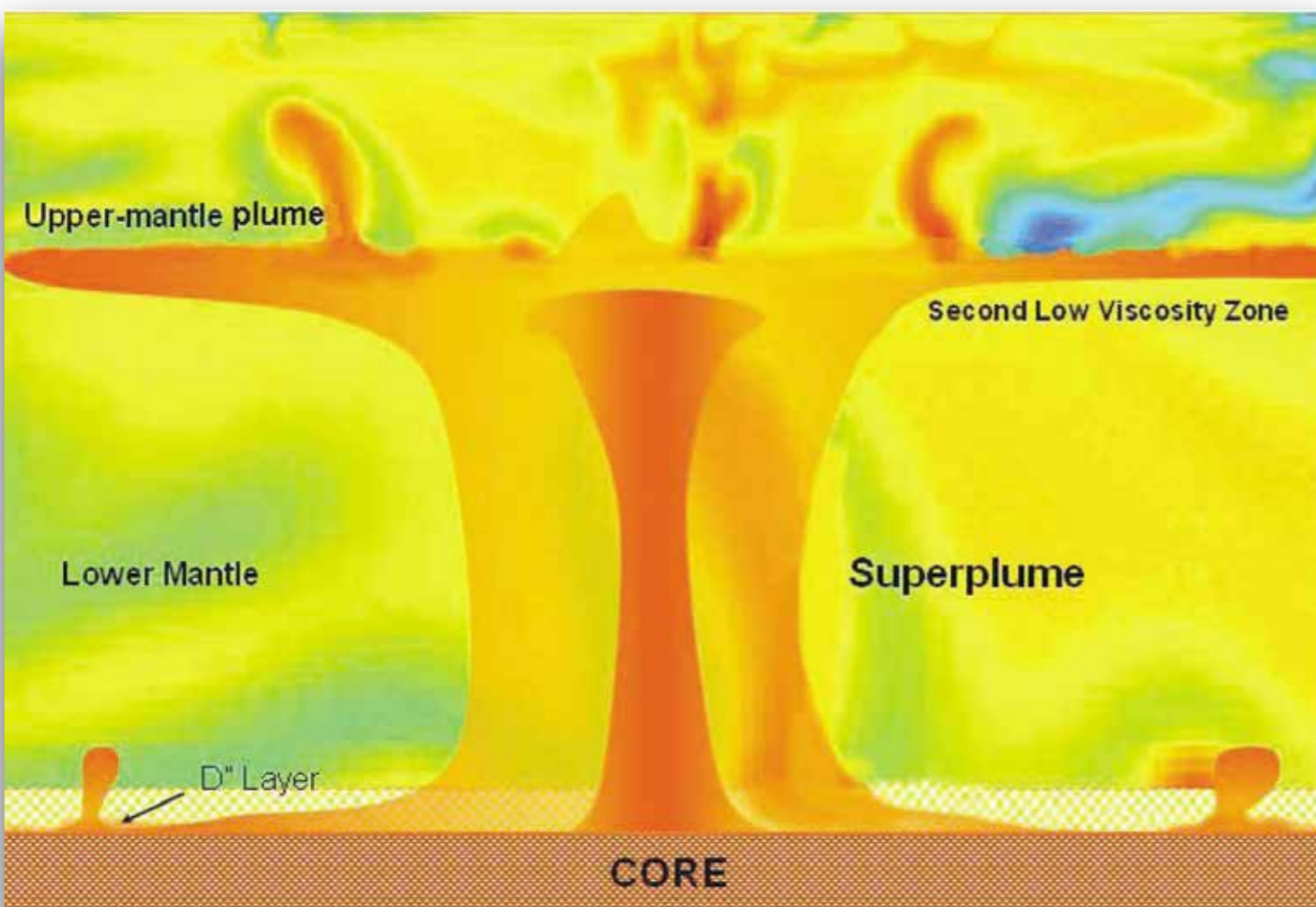
$d \ln V_{sv}$





Multi-scale mantle plumes

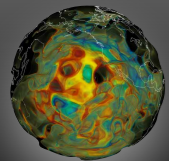
Stable lower-mantle plumes followed by small upper-mantle plumes ==> primary & secondary plumes



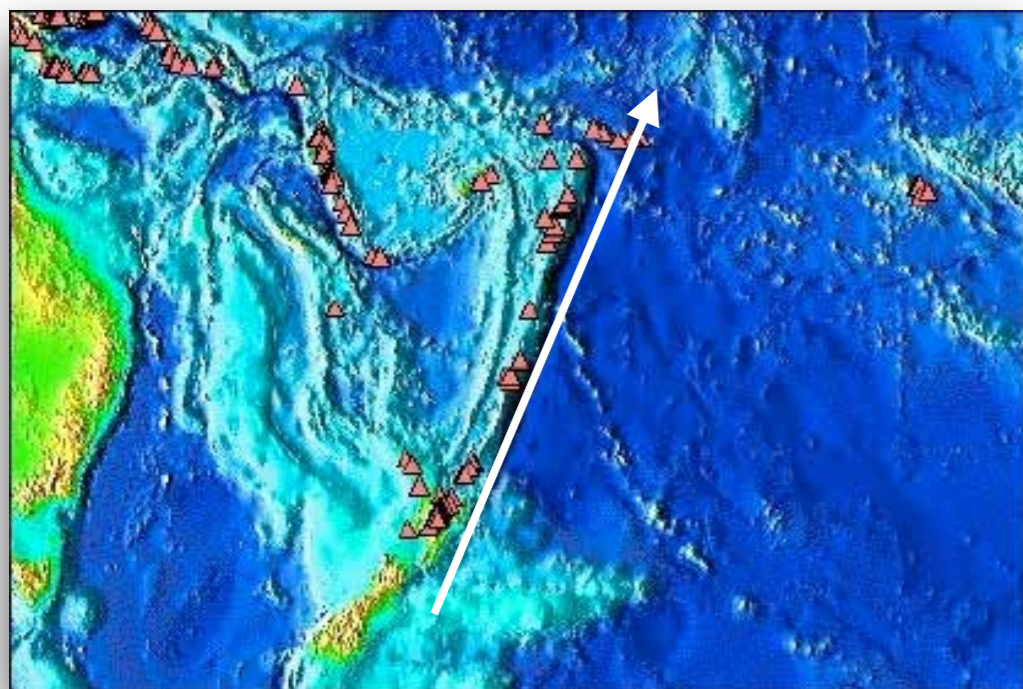
- mantle viscosity
- thermal conductivity
- thermal expansivity
- phase transitions
 - spinel to perovskite
 - perovskite to post-perovskite

Matyska & Yuen, 2007

Courtesy Magali Billen

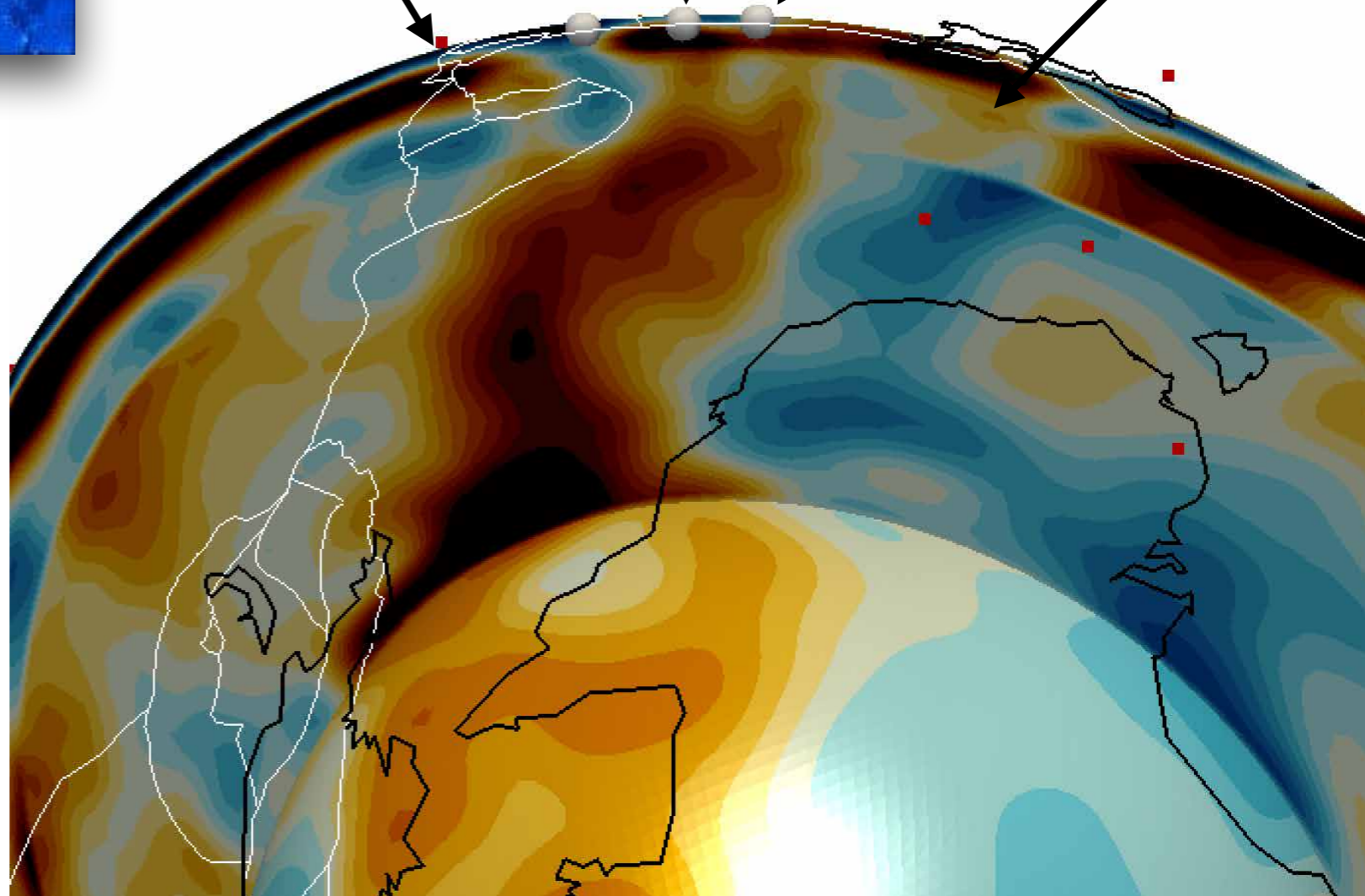
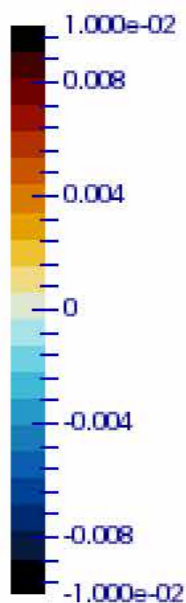


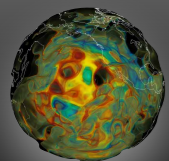
Tonga / Samoa



Raoul Island
Monowai Island
Samoa Island
New Zealand geothermal zone

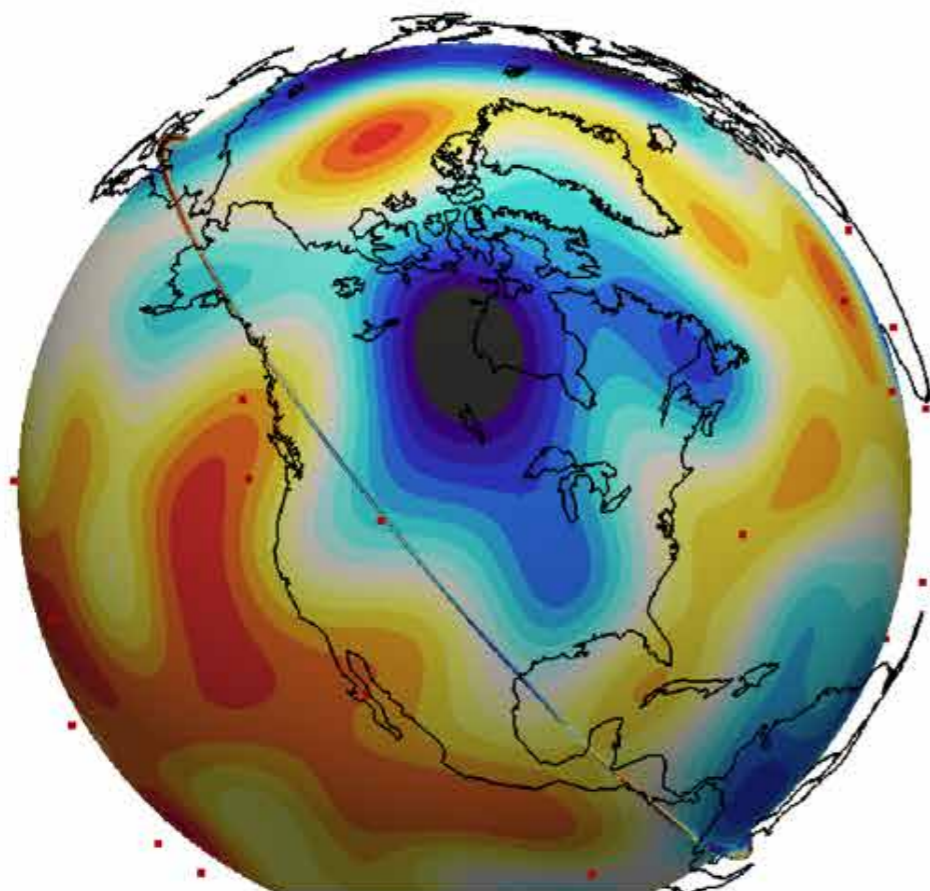
$d\ln(VP/VS)$





North America - 250 km

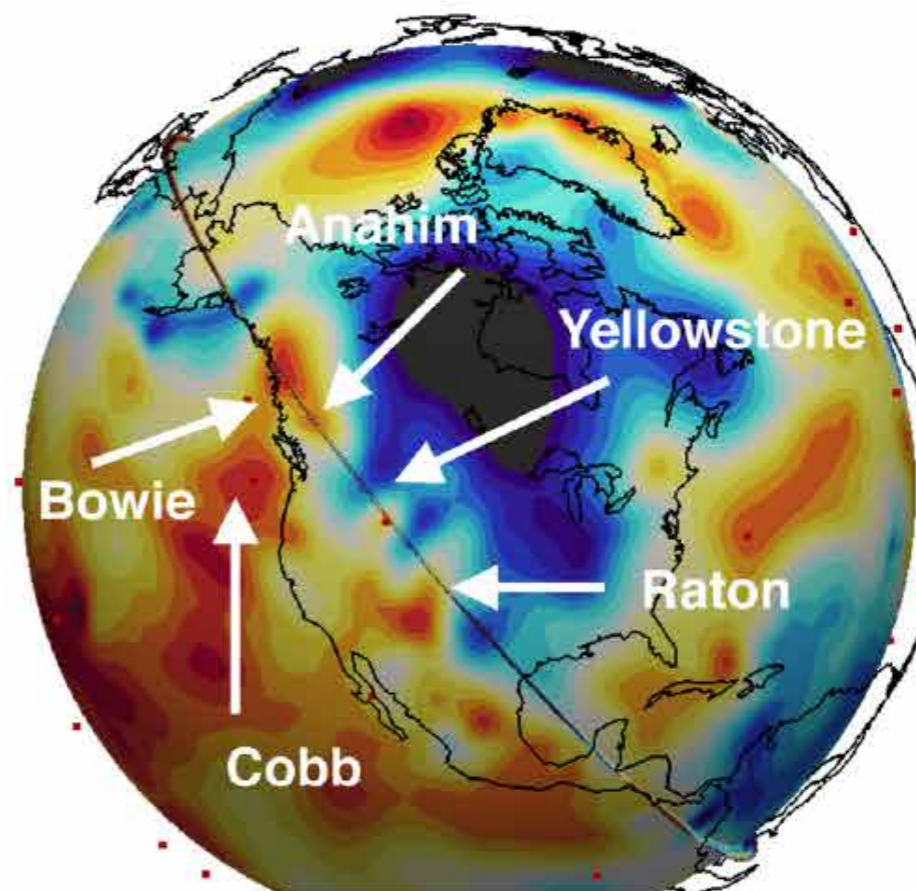
S362ANI



dlnVsv (2%)

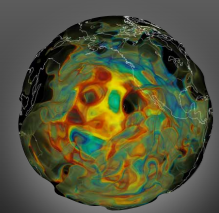
min  max

GLAD-M15



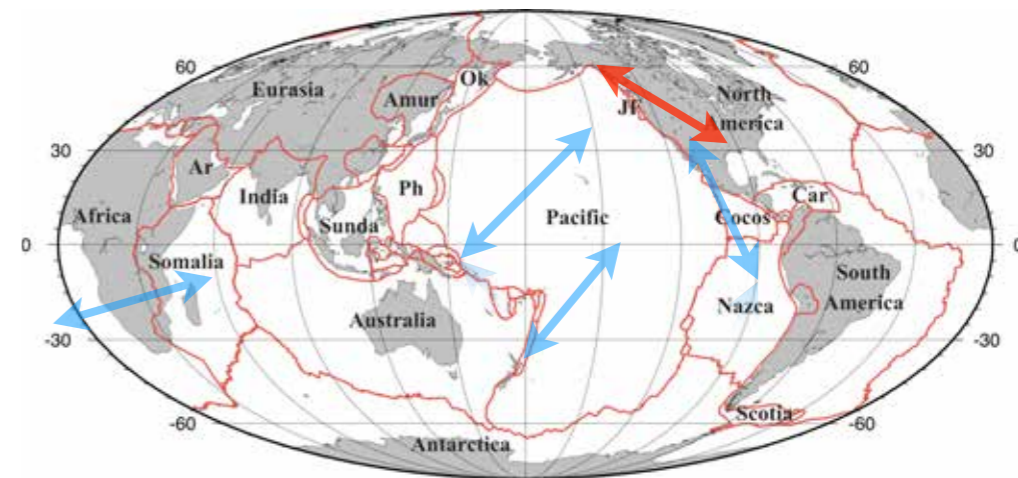
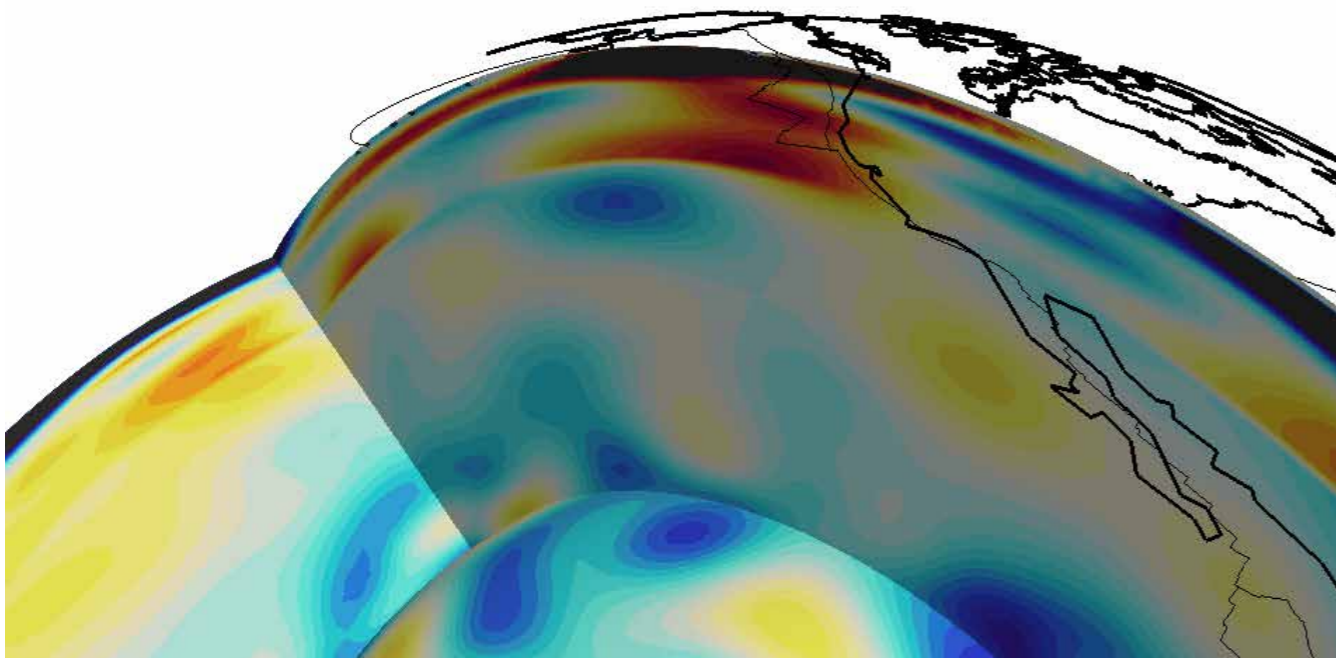
dlnVsv (2%)

min  max



More sections..

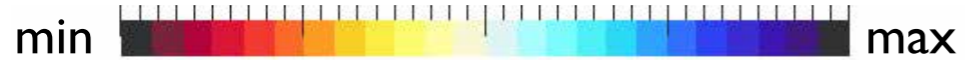
M00 - 1DREF

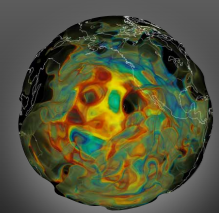


GMT 2005 Apr 29 16:24:38 Plate Boundaries from PB2002 (Peter Bird) Dataset

(2%)

$d \ln V_{sv}$



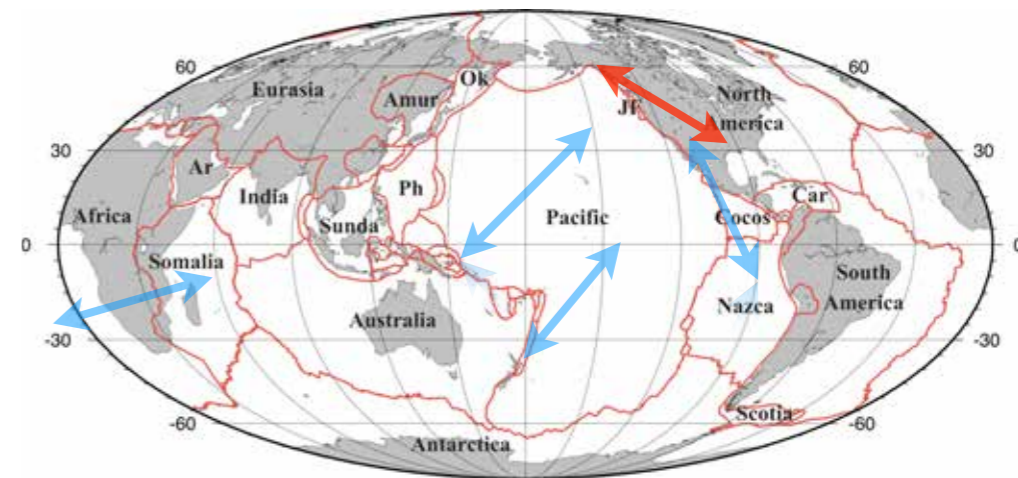
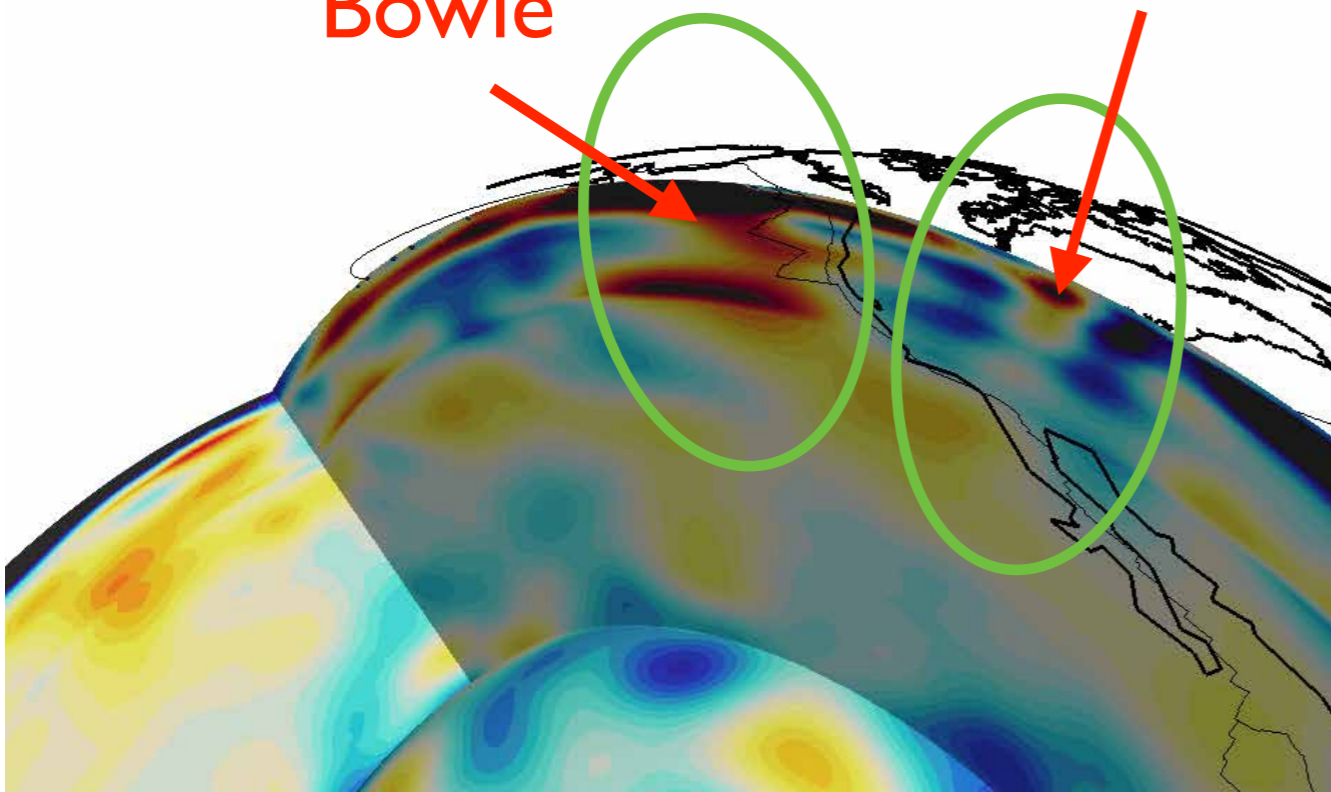


More sections..

M15 - 1DREF

Bowie

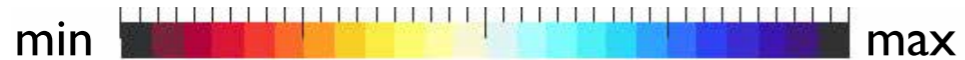
Yellowstone

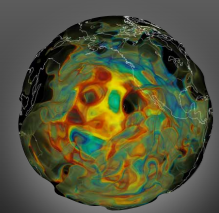


GMT 2005 Apr 29 16:24:38 Plate Boundaries from PB2002 (Peter Bird) Dataset

(2%)

$d \ln V_{sv}$



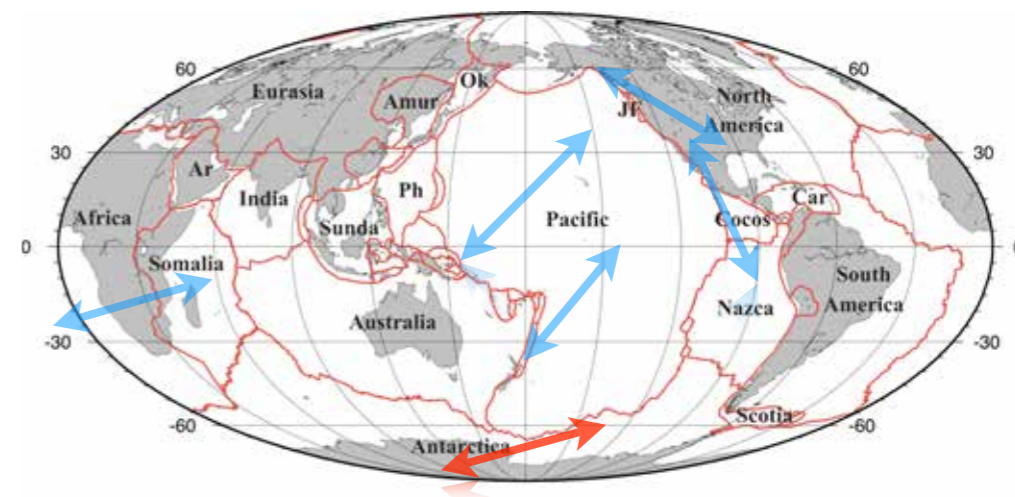
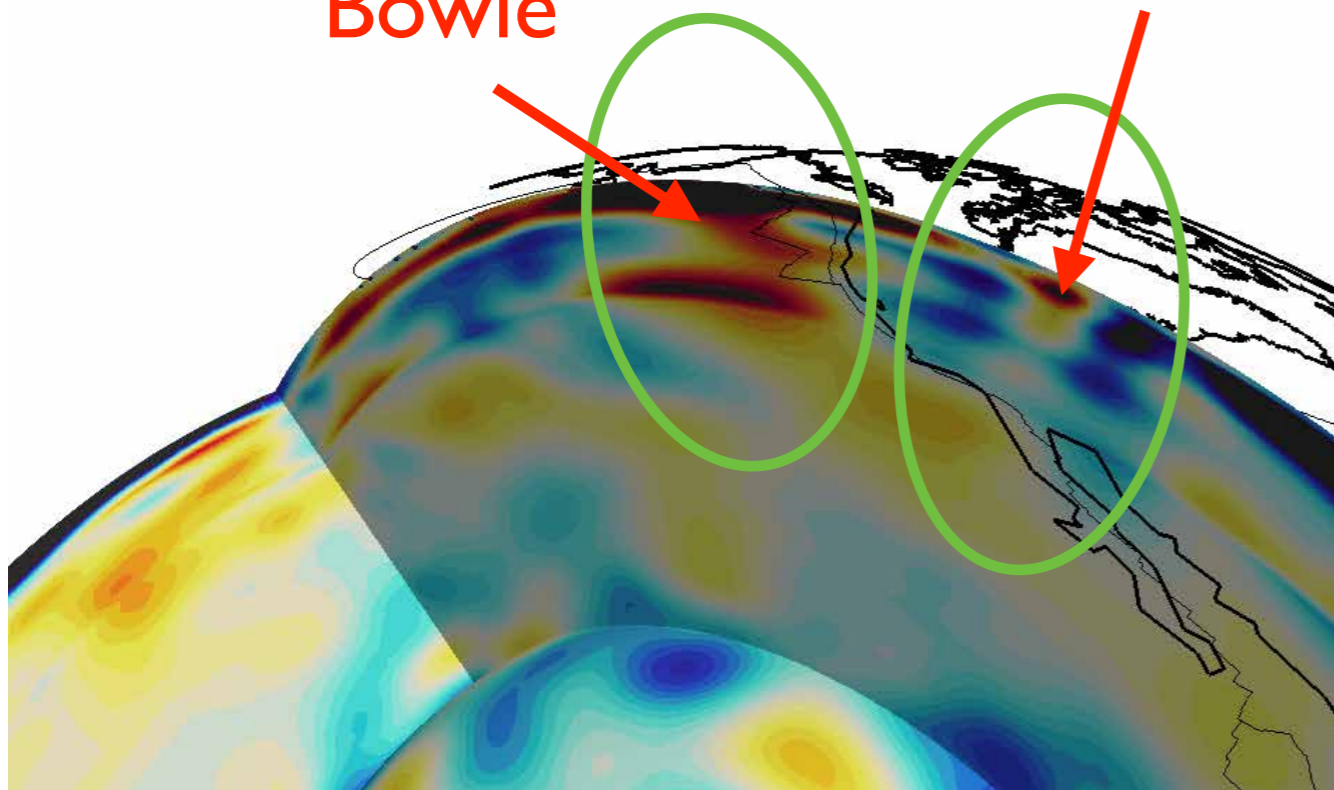


More sections..

M15 - 1DREF

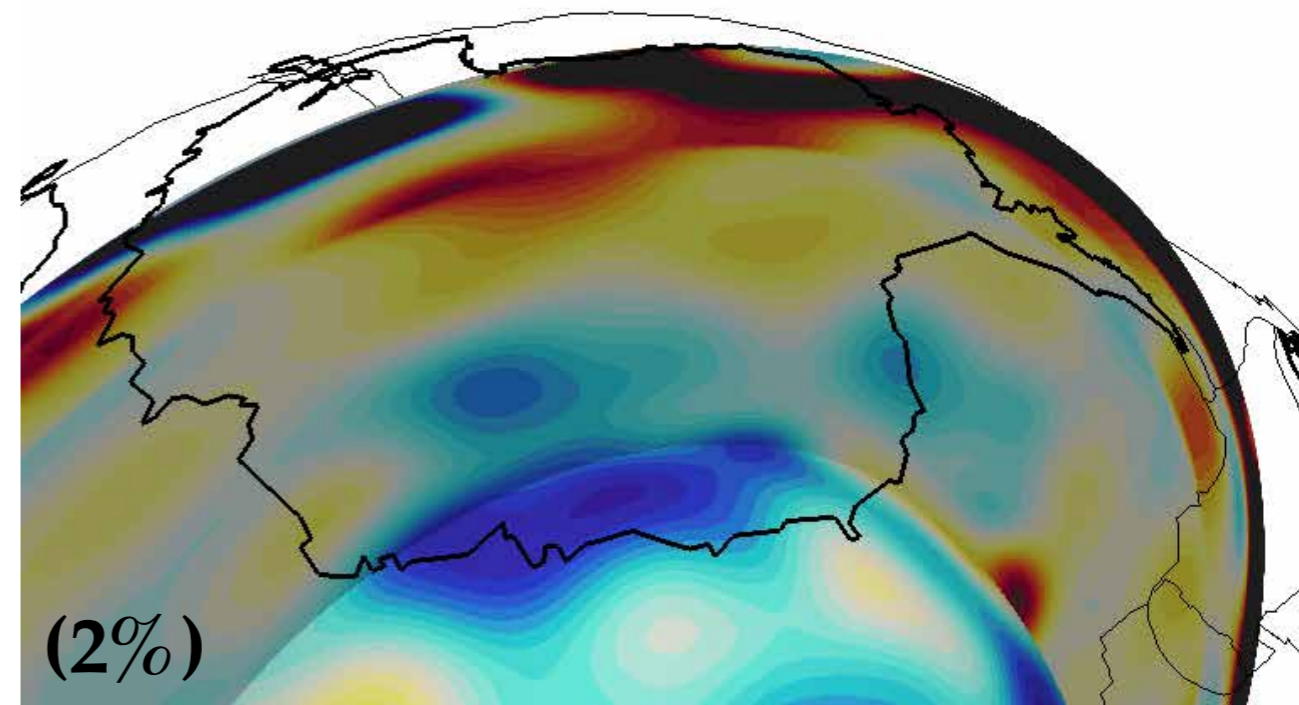
Bowie

Yellowstone



GMT 2005 Apr 29 16:24:38 Plate Boundaries from PB2002 (Peter Bird) Dataset

M00 - 1DREF

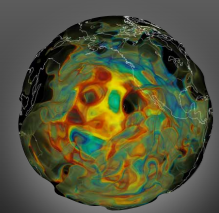


(2%)

(2%)

$d \ln V_{sv}$



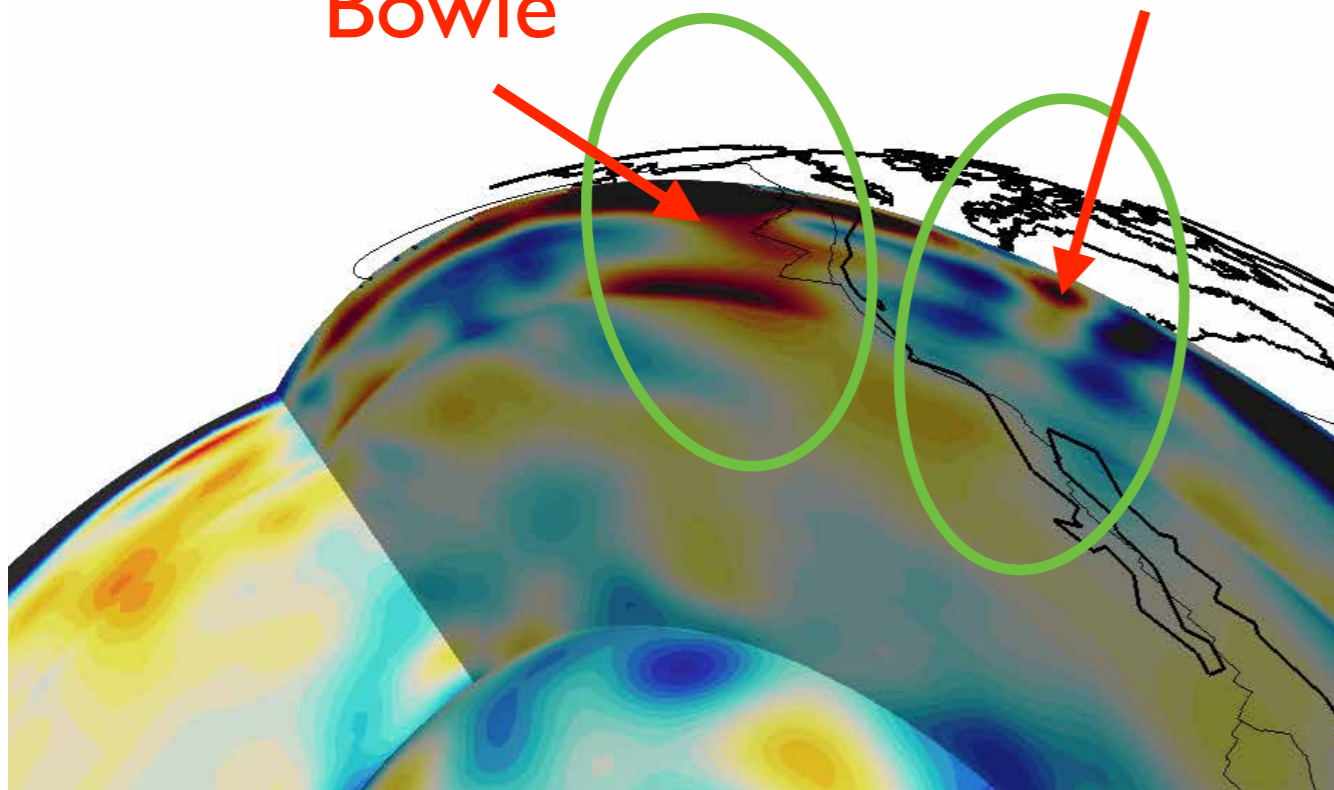


More sections..

M15 - 1DREF

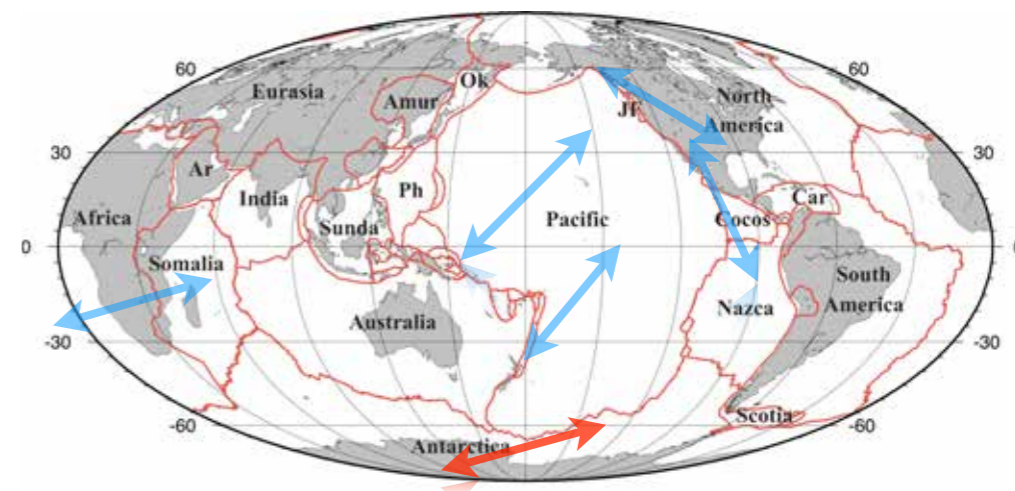
Bowie

Yellowstone



(2%)

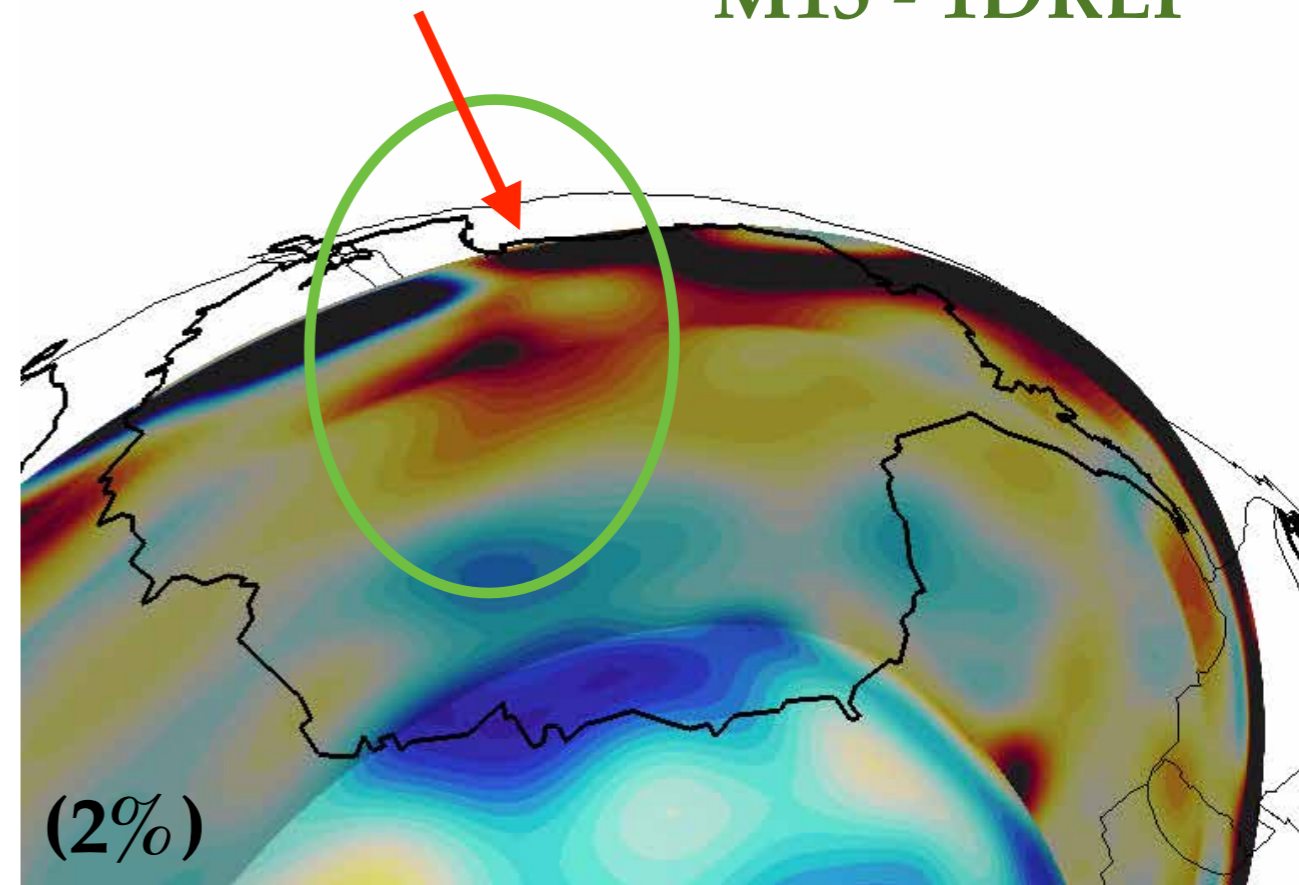
$d \ln V_{sv}$



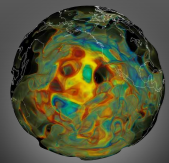
GMT 2005 Apr 29 16:24:38 Plate Boundaries from PB2002 (Peter Bird) Dataset

Erebus

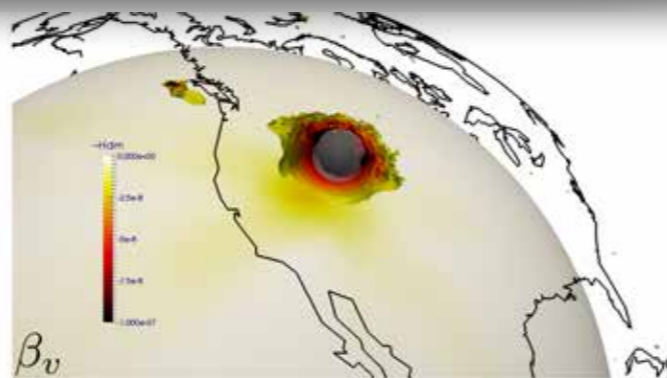
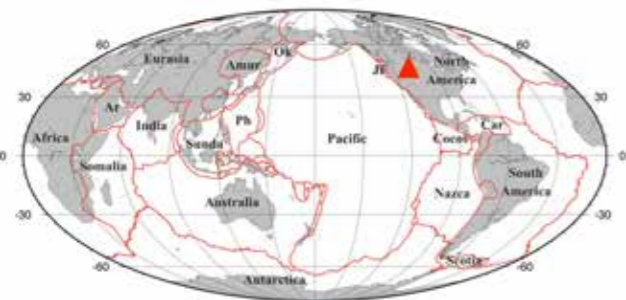
M15 - 1DREF



(2%)

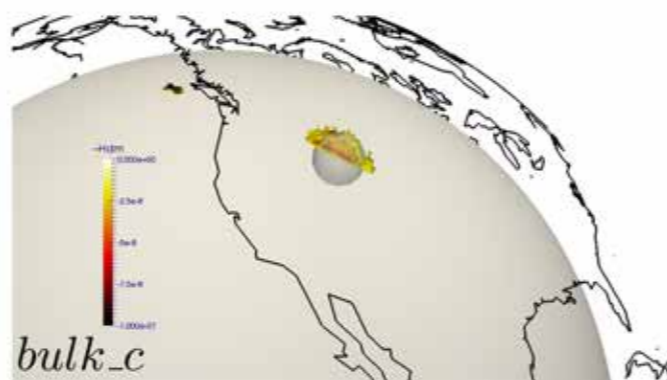
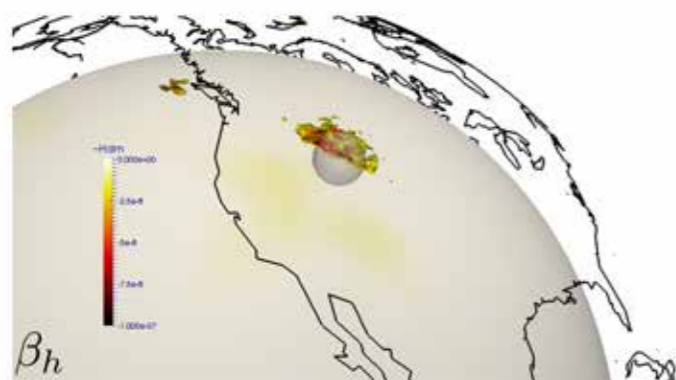


Resolution tests

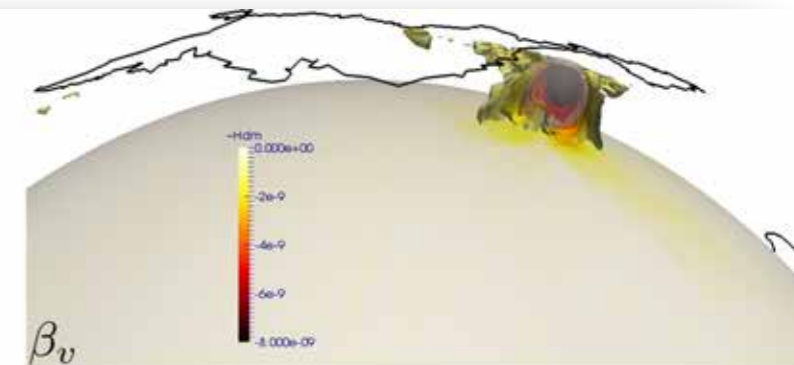


**PSF
(Point-Spread Function)
tests
(Fichtner & Trampert 2011, GJI)**

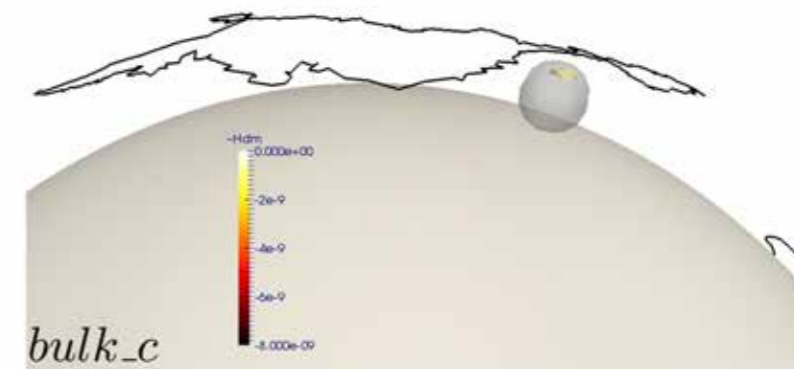
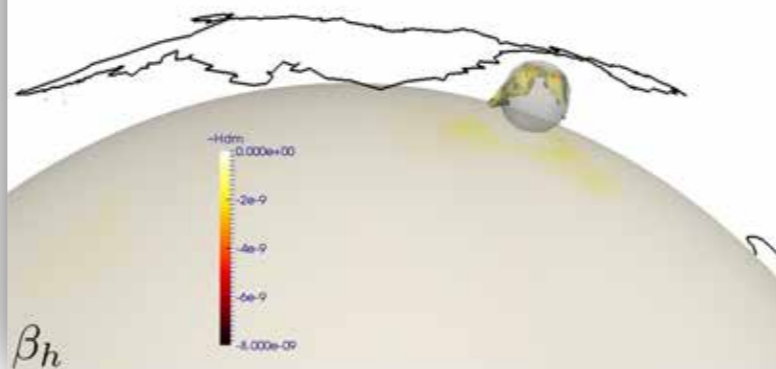
$$H(m)\delta m \approx g(m + \delta m) - g(m)$$

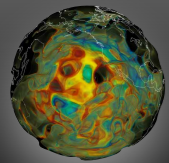


Yellowstone

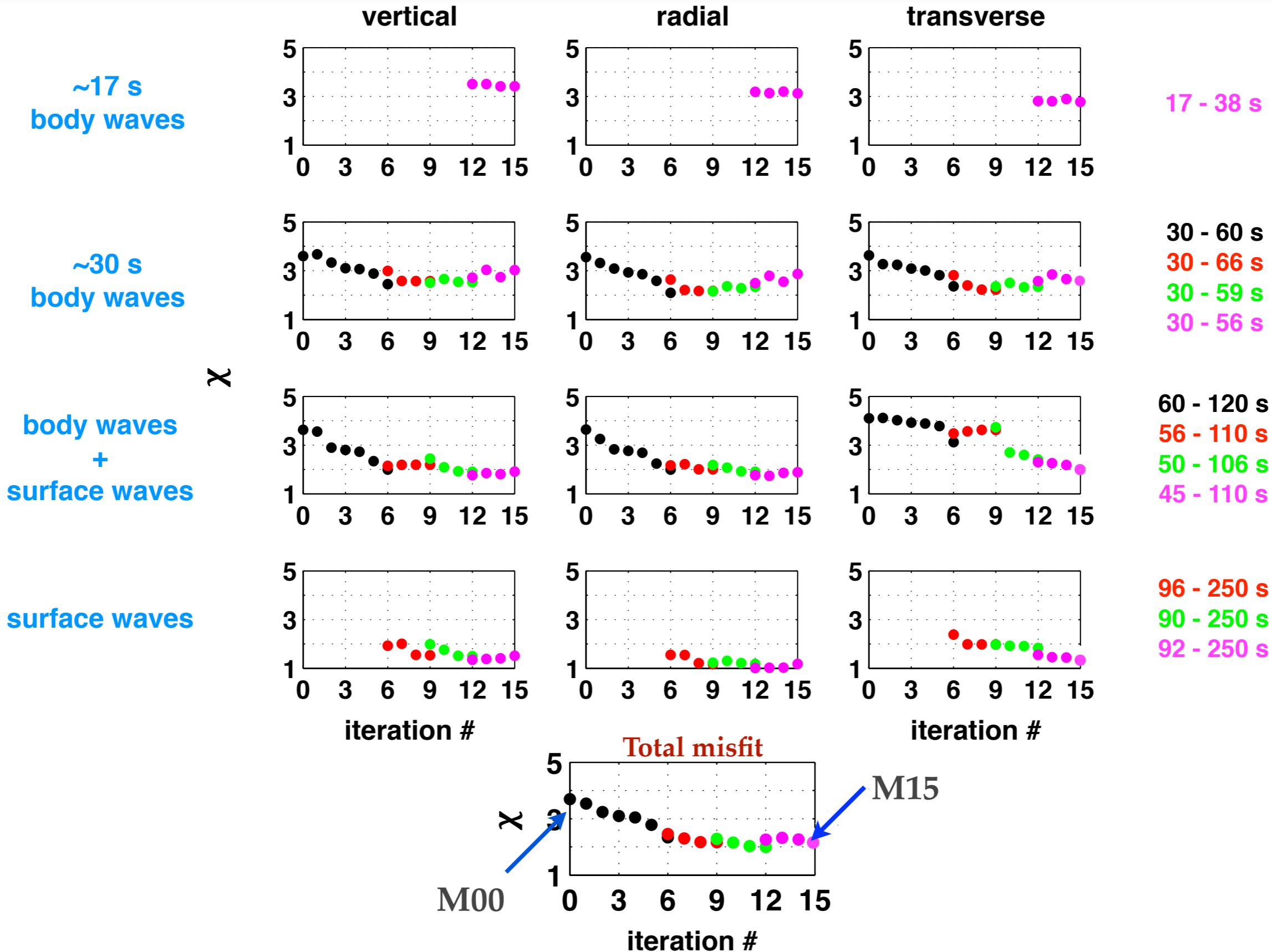


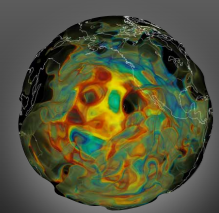
Erebus





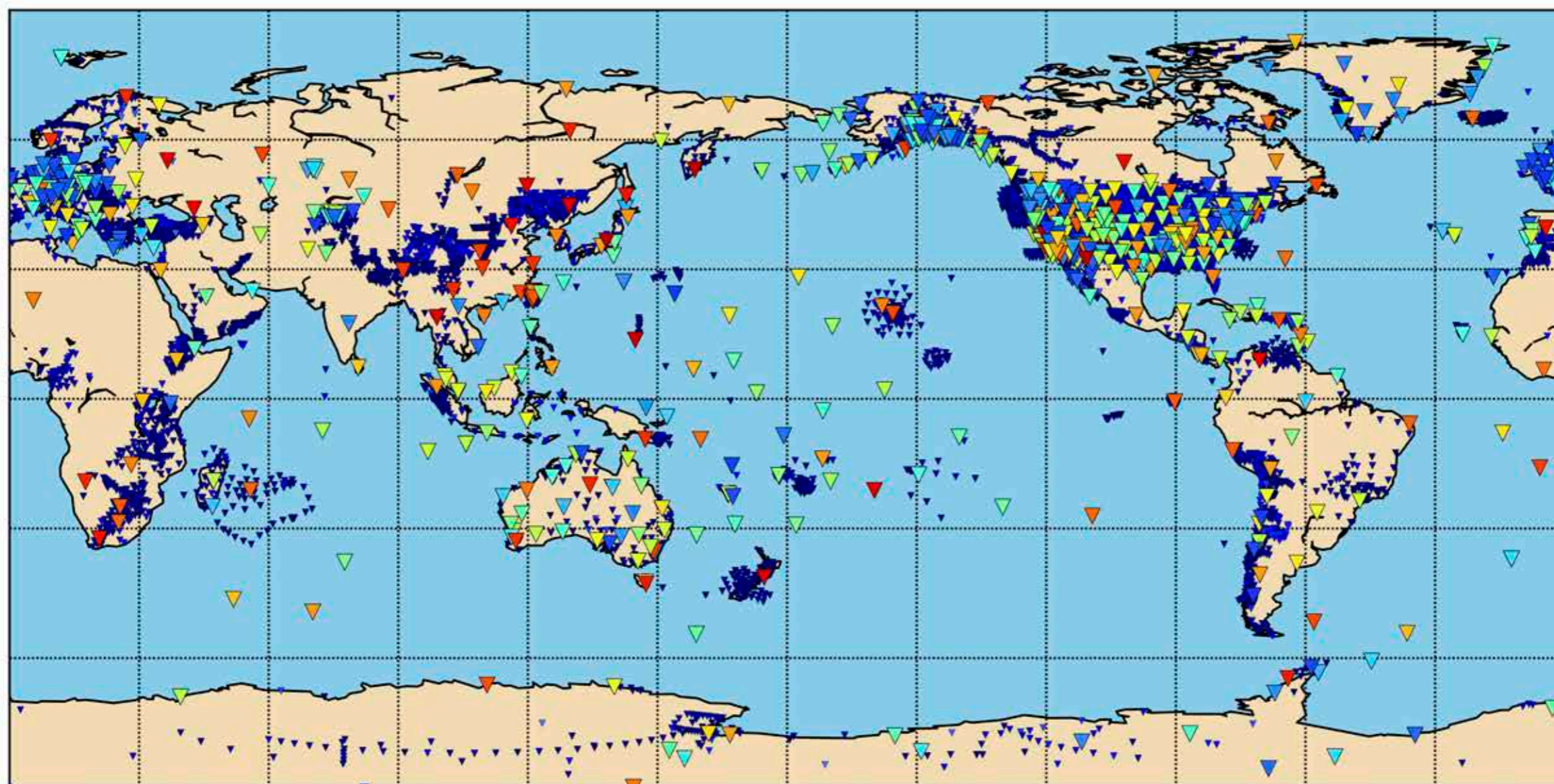
Misfit reduction



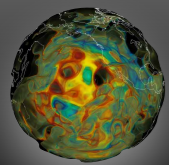


GLAD-M24*

Wenjie Lei, Youyi Ruan

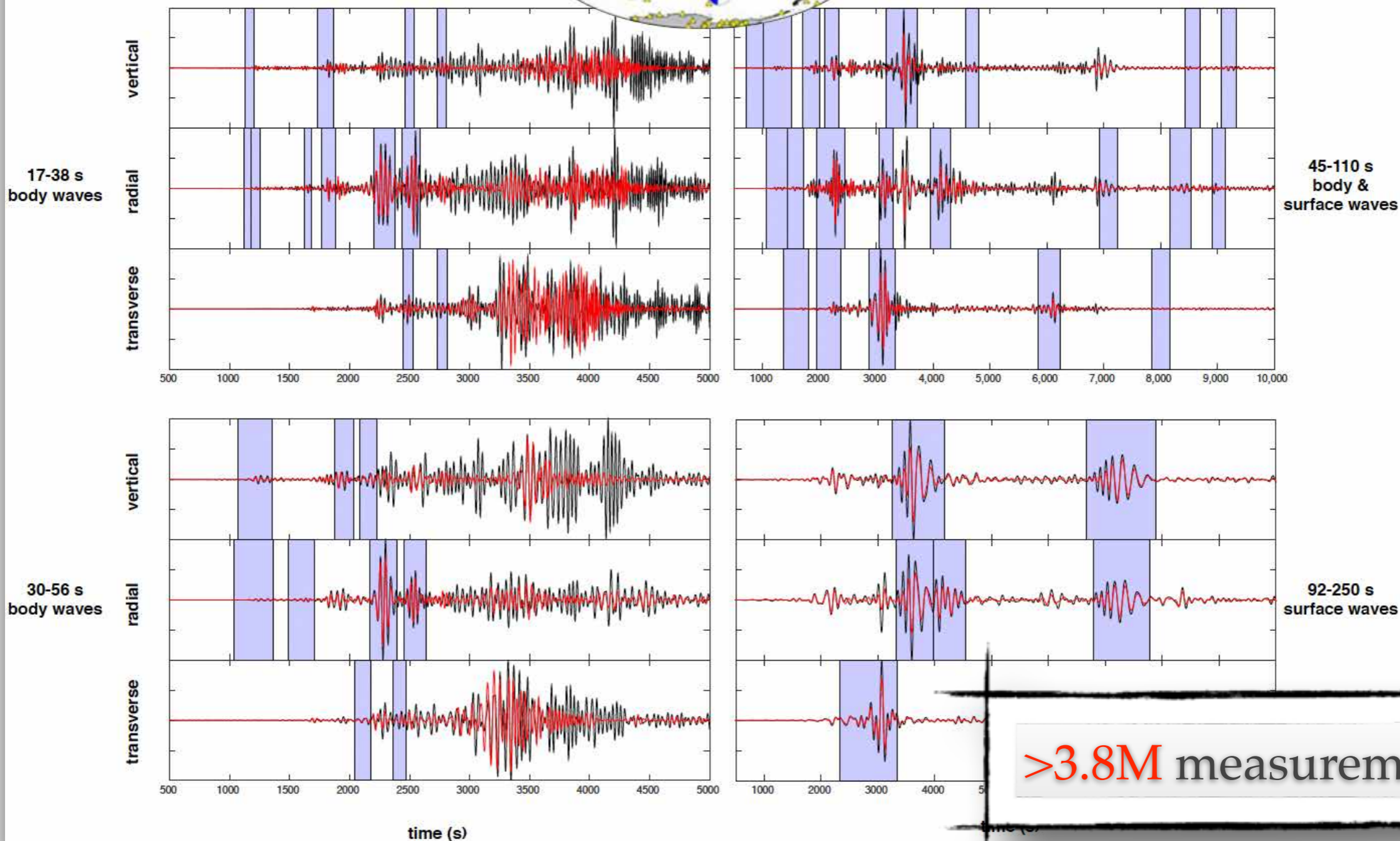
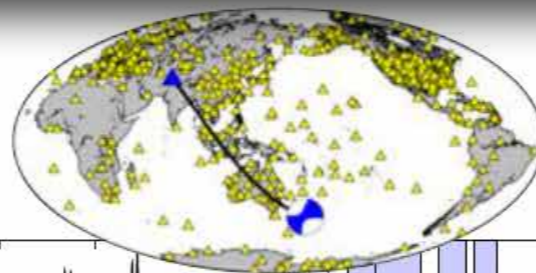


11,800 seismic stations, 1480 earthquakes



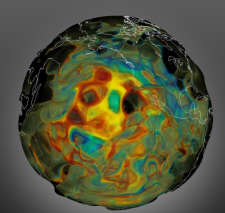
Data selection / measurements

window selections:
FLEXWIN
Maggi et al. 2009

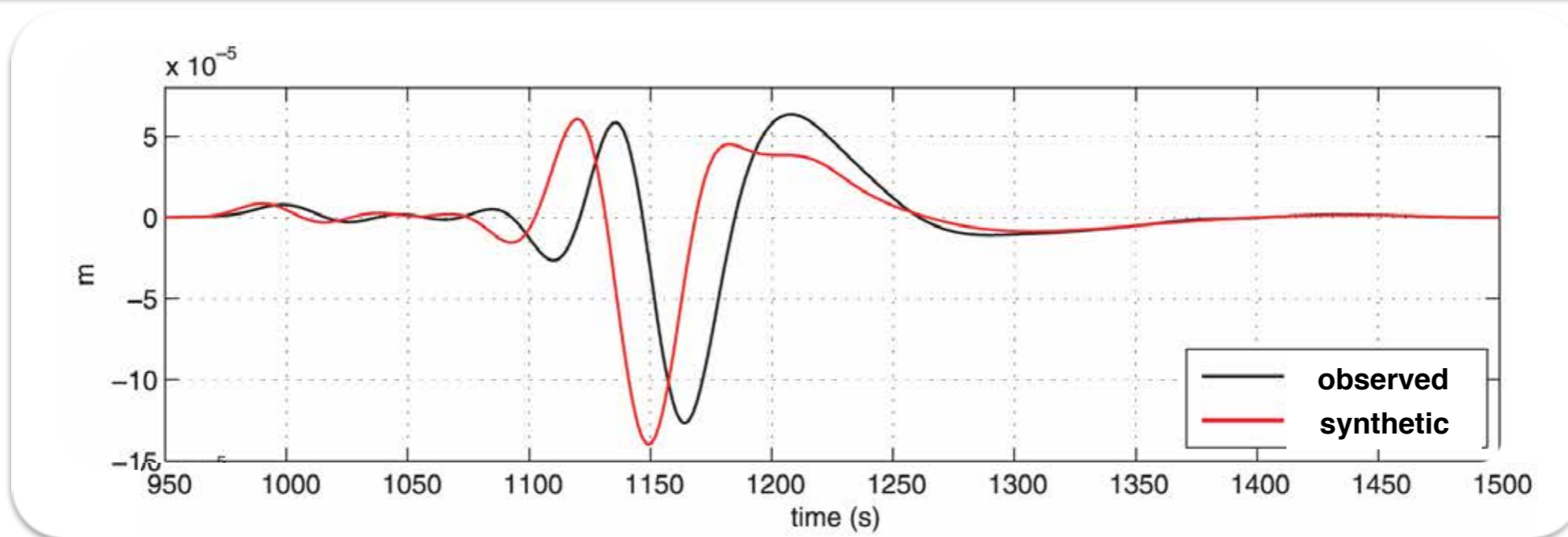


>3.8M measurements

$$\chi_c = \frac{1}{N_c} \sum_{e=1}^E \sum_{i=1}^{N_c^s} \int w_i(\omega) \left[\frac{\Delta\tau_i(\omega)}{\sigma_i} \right]^2 d\omega / \int w_i(\omega) d\omega$$



Data / measurements



$$\chi = \frac{1}{2} \sum_s \sum_r \int ||d^{obs}(t) - d^{syn}(t, m)||^2 dt$$

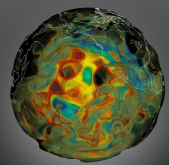
Waveform misfit

$$\chi = \frac{1}{2} \sum_s \sum_r [T^{obs} - T^{syn}]^2$$

Traveltime misfit

$$\chi = \frac{1}{2} \sum_s \sum_r \left[\ln \frac{A^{obs}}{A^{syn}} \right]^2$$

Amplitude misfit

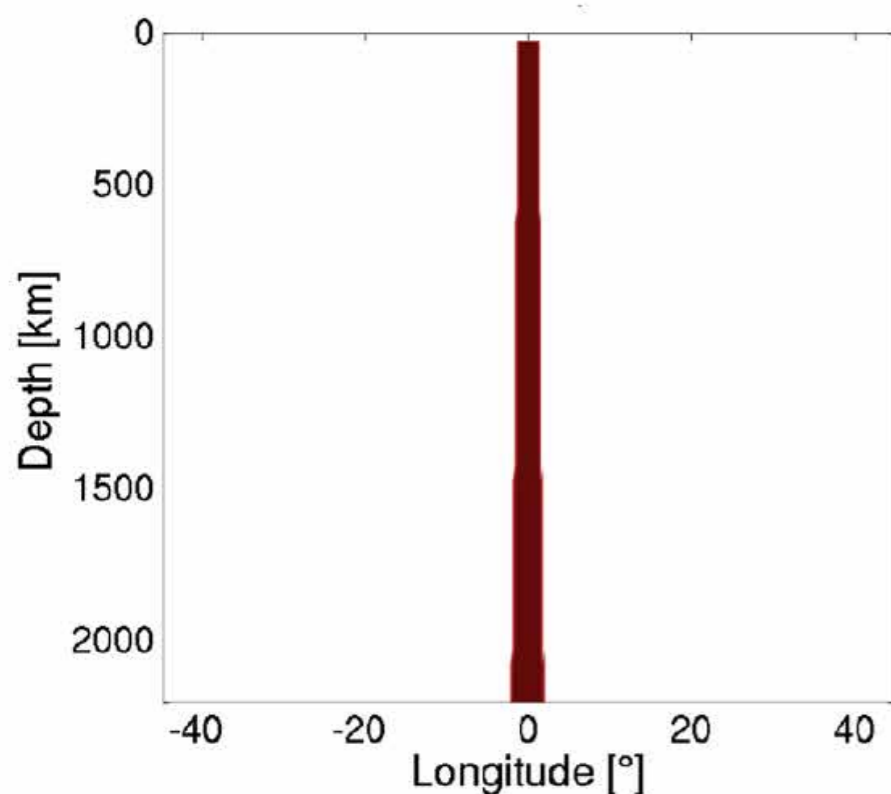


Define new measurements

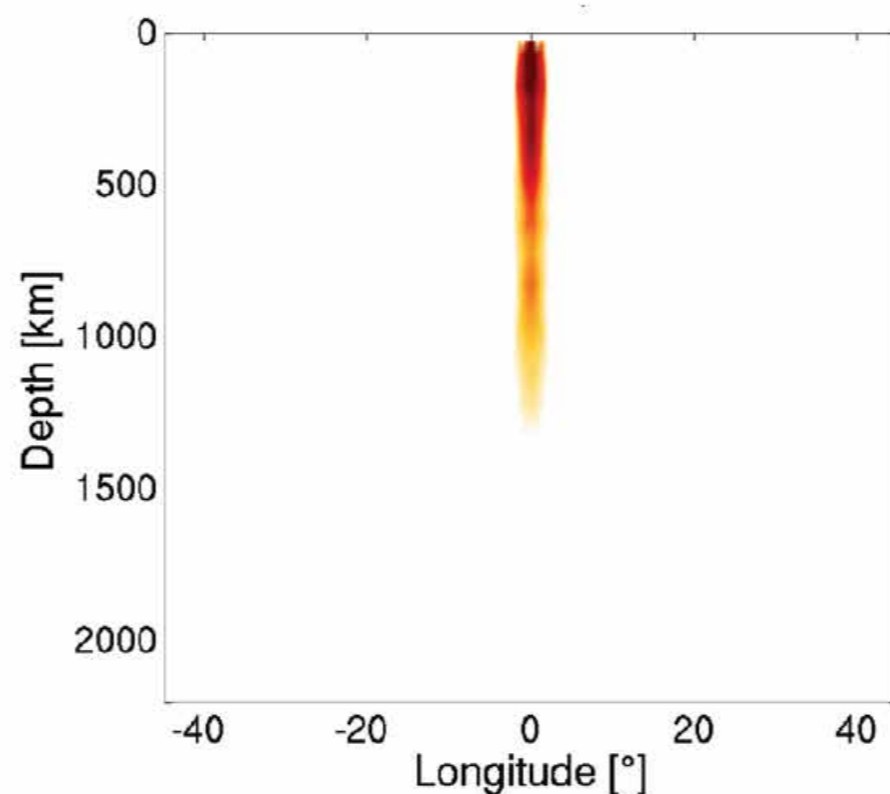
Instantaneous-phase misfit

$$\chi_{ip} = \frac{1}{2} \sum_{s,r} \int_0^T \|\phi_d(t) - \phi_s(t, \mathbf{m})\|^2 dt$$

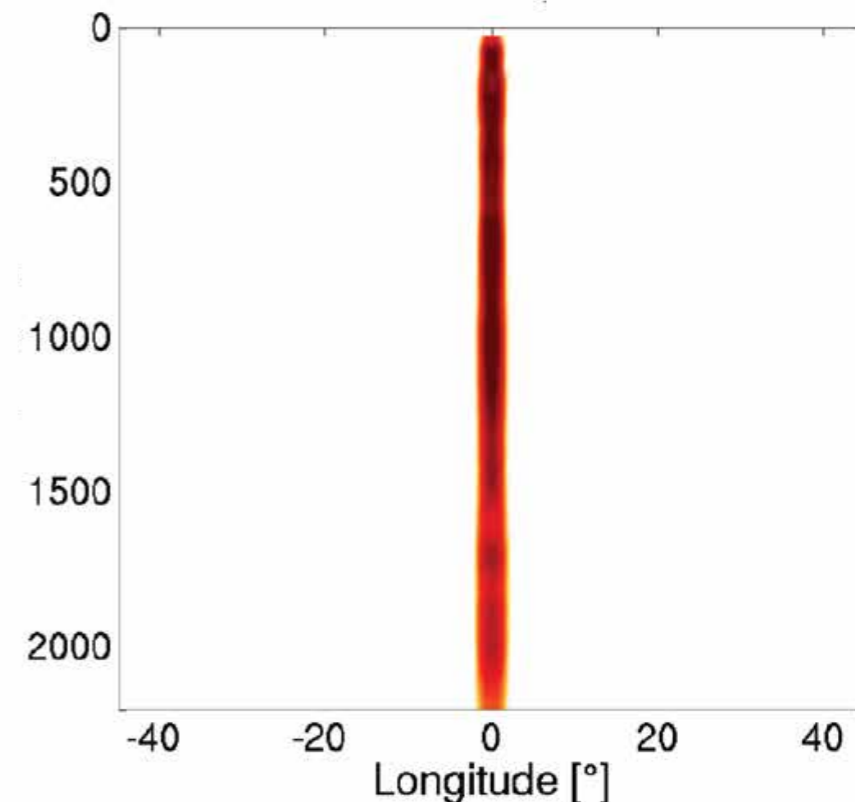
E. Bozdag, J. Trampert &
J. Tromp, 2011, GJI



Target model

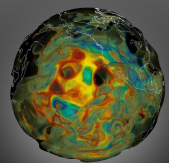


CC traveltimes

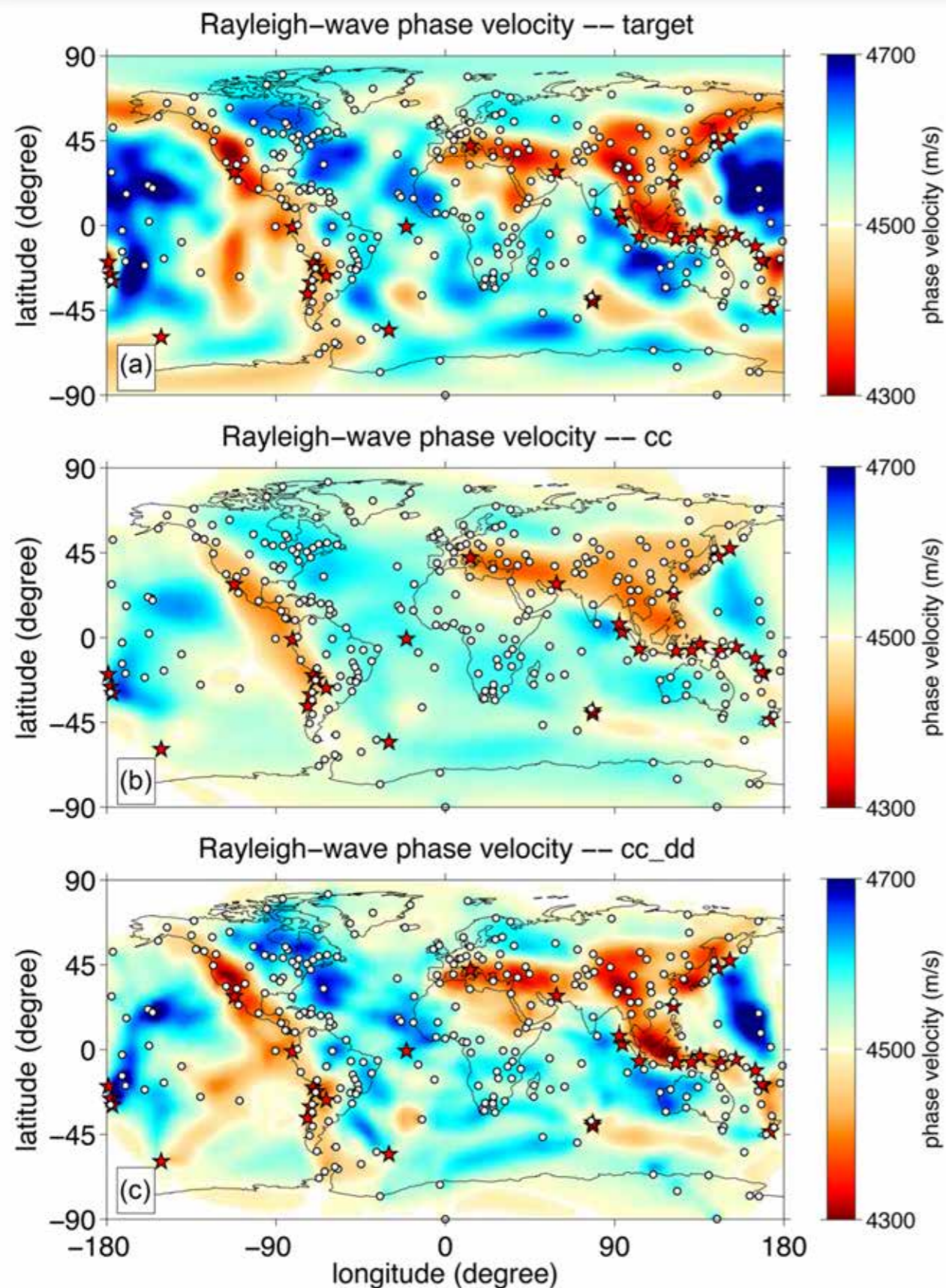


IP

Rickers, Fichtner & Trampert, 2012, GJI

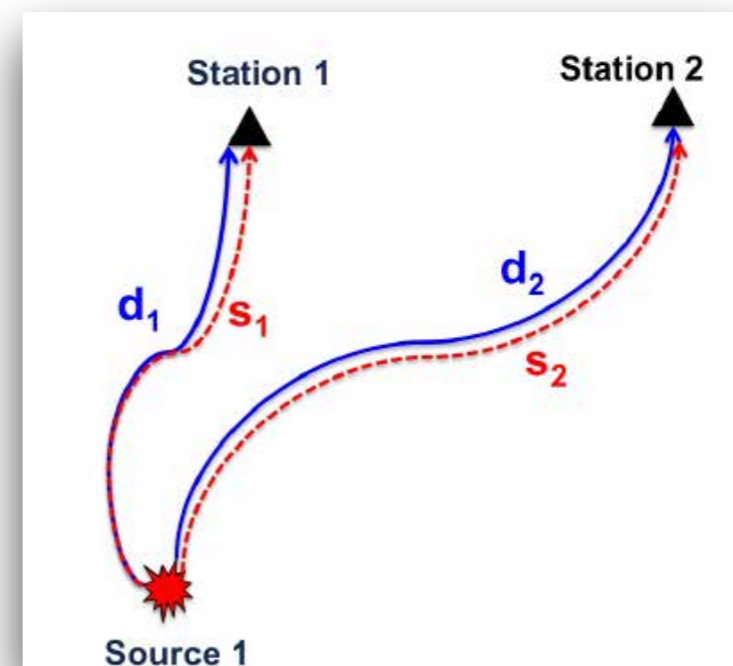


Define new measurements



Double-difference misfit

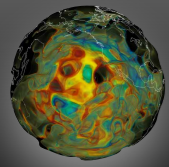
Y. Yuan, F. J. Simons & J. Tromp,
GJI, 2016



**reduce source uncertainties
& bias in sparsely sampled regions**

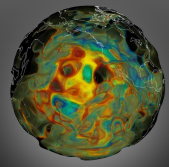
**Exponentiated phase -
Hybrid objective functions
(CC+IP, MT+IP, etc..)**

Yuan, Bozdağ, Simons, Gao, *AGU*, 2016



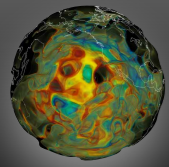
Computational requirements

core hours 180-m seismograms	1 event	1 iteration (255 events)	20 iterations
CPU - ~27 s	~5,400	~1.4M	~28M



Computational requirements

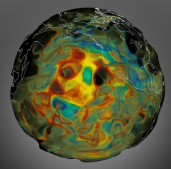
core hours 180-m seismograms	1 event	1 iteration (255 events)	20 iterations
CPU - ~27 s	~5,400	~1.4M	~28M
CPU - ~17 s	~16,890	~4.3M	~86M



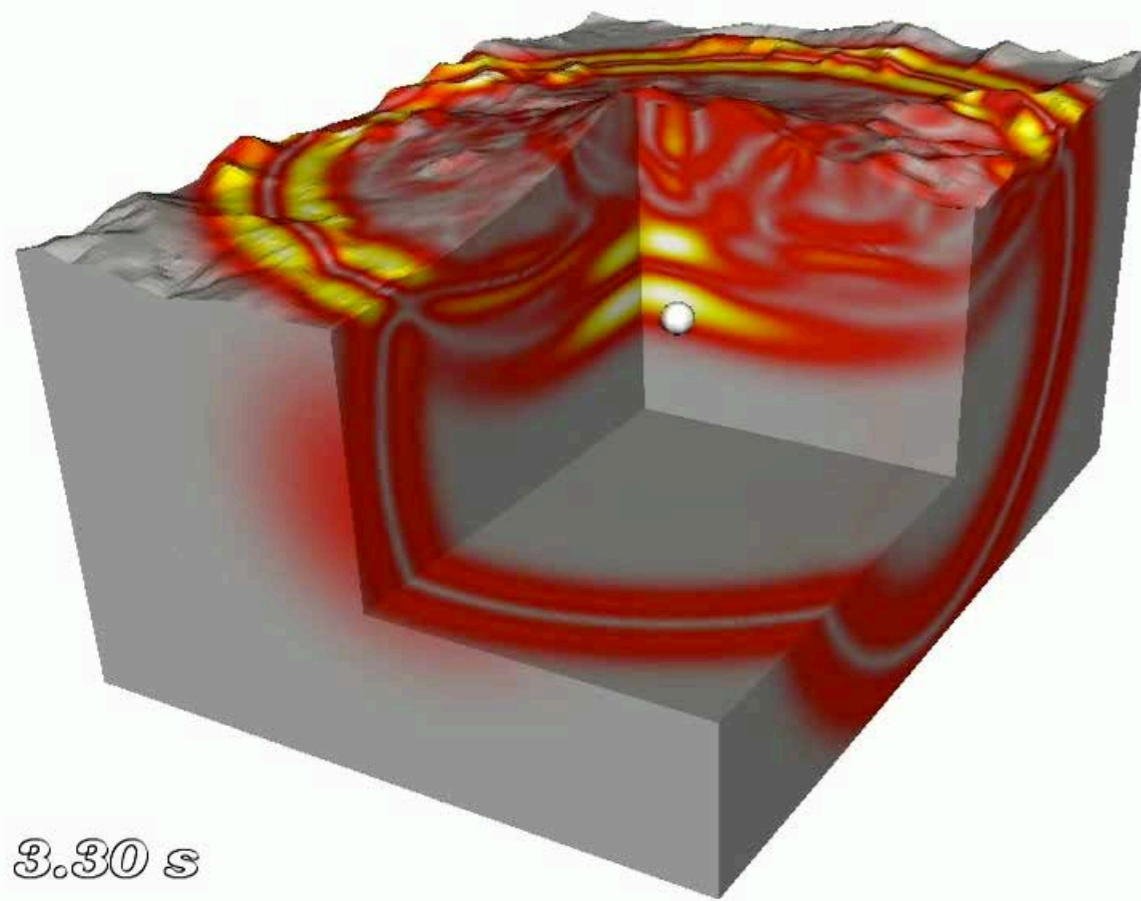
Computational requirements

core hours 180-m seismograms	1 event	1 iteration (255 events)	20 iterations
CPU - ~27 s	~5,400	~1.4M	~28M
CPU - ~17 s	~16,890	~4.3M	~86M

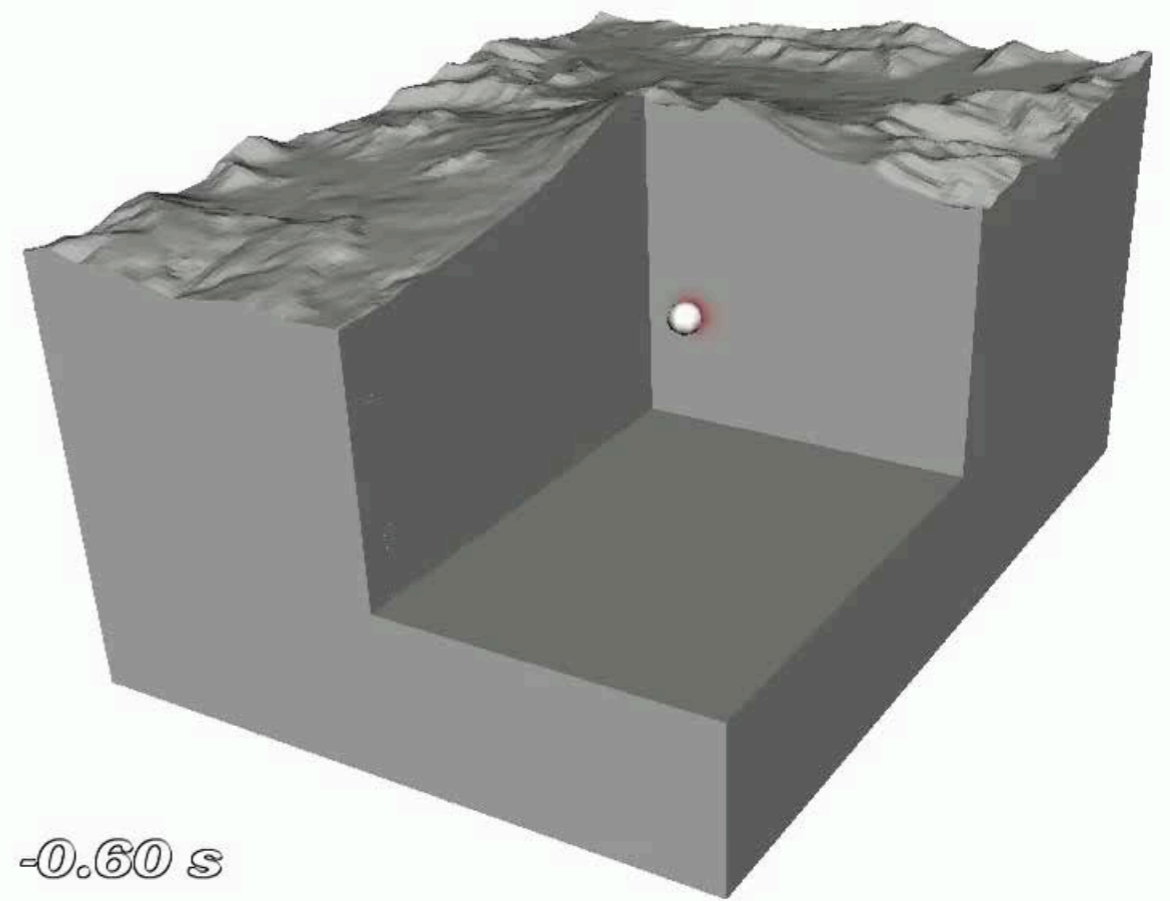
- **Code-method development/optimisation**
- **Current/future computational resources**



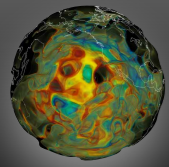
GPU computing



GPU

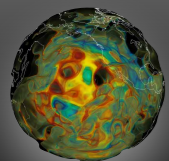


CPU



Computational requirements

core hours 180-m seismograms	1 event	1 iteration (255 events)	20 iterations
CPU - ~27 s	~5,400	~1.4M	~28M
CPU - ~17 s	~16,890	~4.3M	~86M
GPU - ~17 s	~600	~153,000	~3.1M



Computational requirements

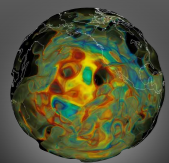
core hours 180-m seismograms	1 event	1 iteration (255 events)	20 iterations
CPU - ~27 s	~5,400	~1.4M	~28M
CPU - ~17 s	~16,890	~4.3M	~86M
GPU - ~17 s	~600	~153,000	~3.1M



ORNL's Cray XK7 TITAN system
16,384 GPUs

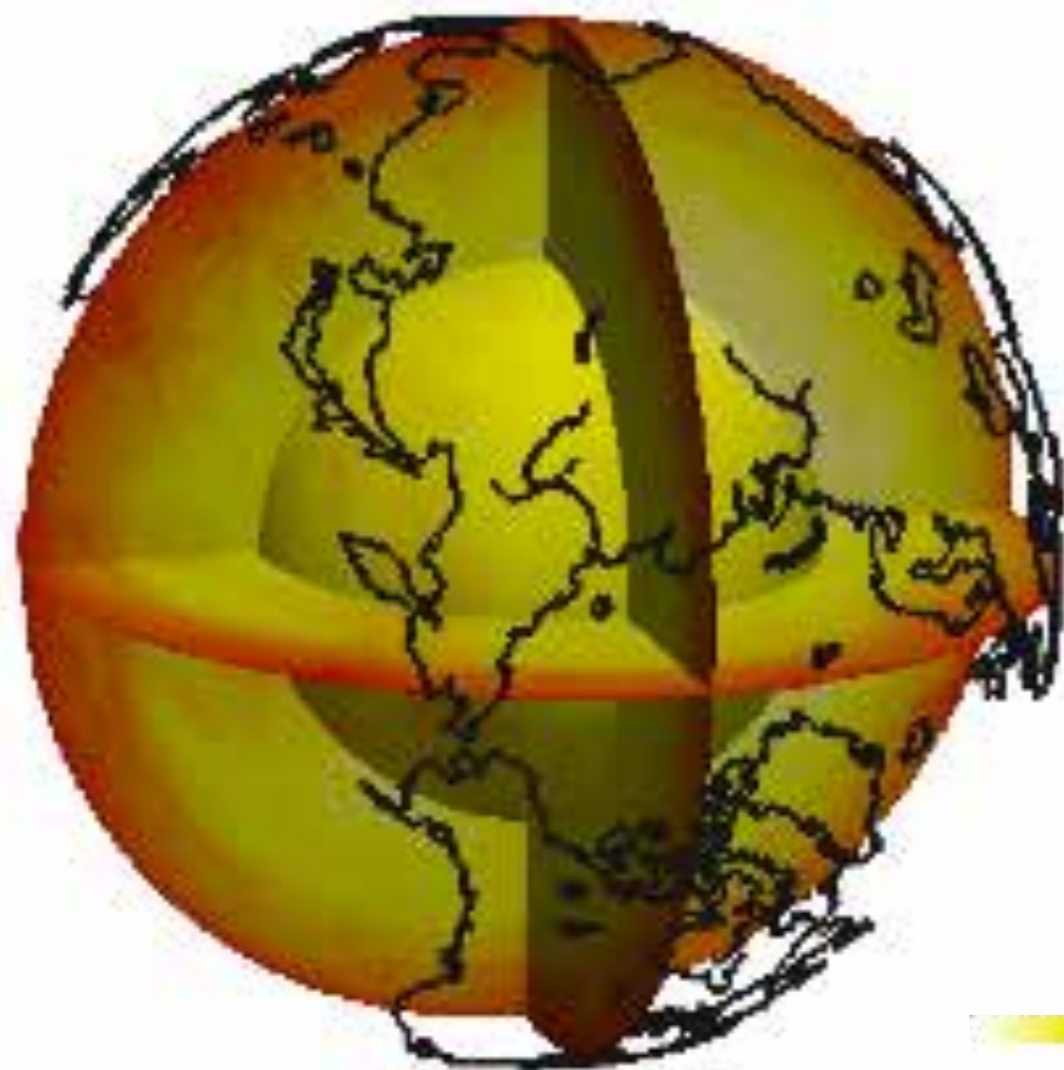


ORNL's next-generation supercomputer
SUMMIT, x5 faster, 4086 nodes
The goal is ~9 s, ultimately 1 s (*exascale era*).

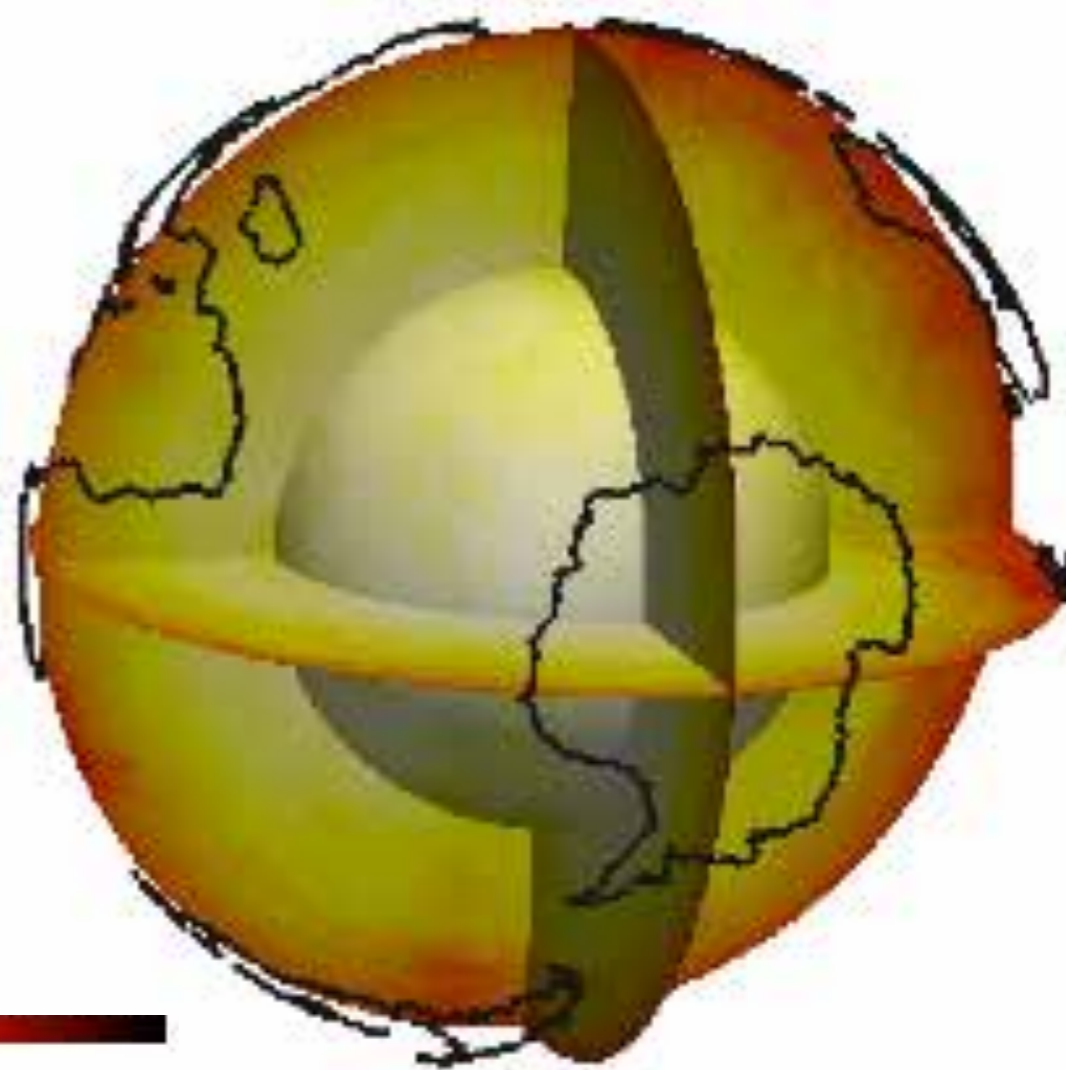


Ray (kernel) density

$$P(\mathbf{x}) = \sum_s \int \partial_t^2 \mathbf{s}^\dagger(\mathbf{x}, -t) \cdot \partial_t^2 \mathbf{s}(\mathbf{x}, t) dt \quad (\text{diagonal terms of Hessian})$$



Northern hemisphere

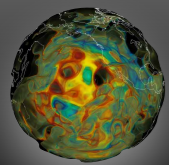


Southern hemisphere

log scale

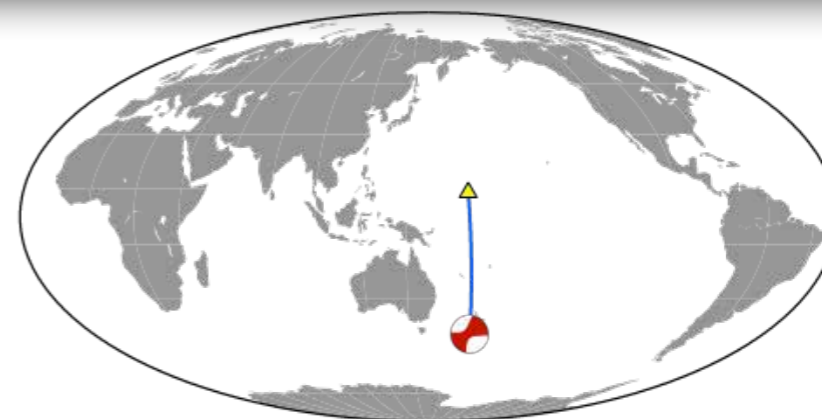
min

max



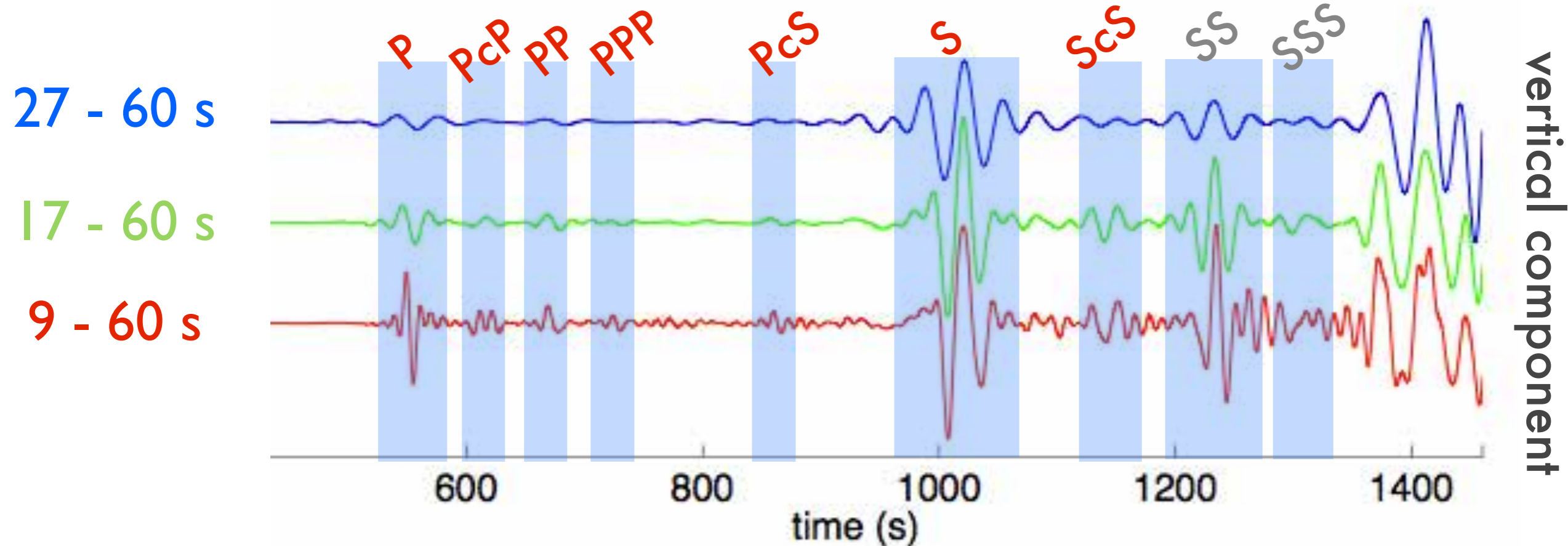
Increase resolution of simulations

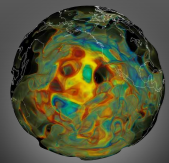
add *high-frequency body waves*



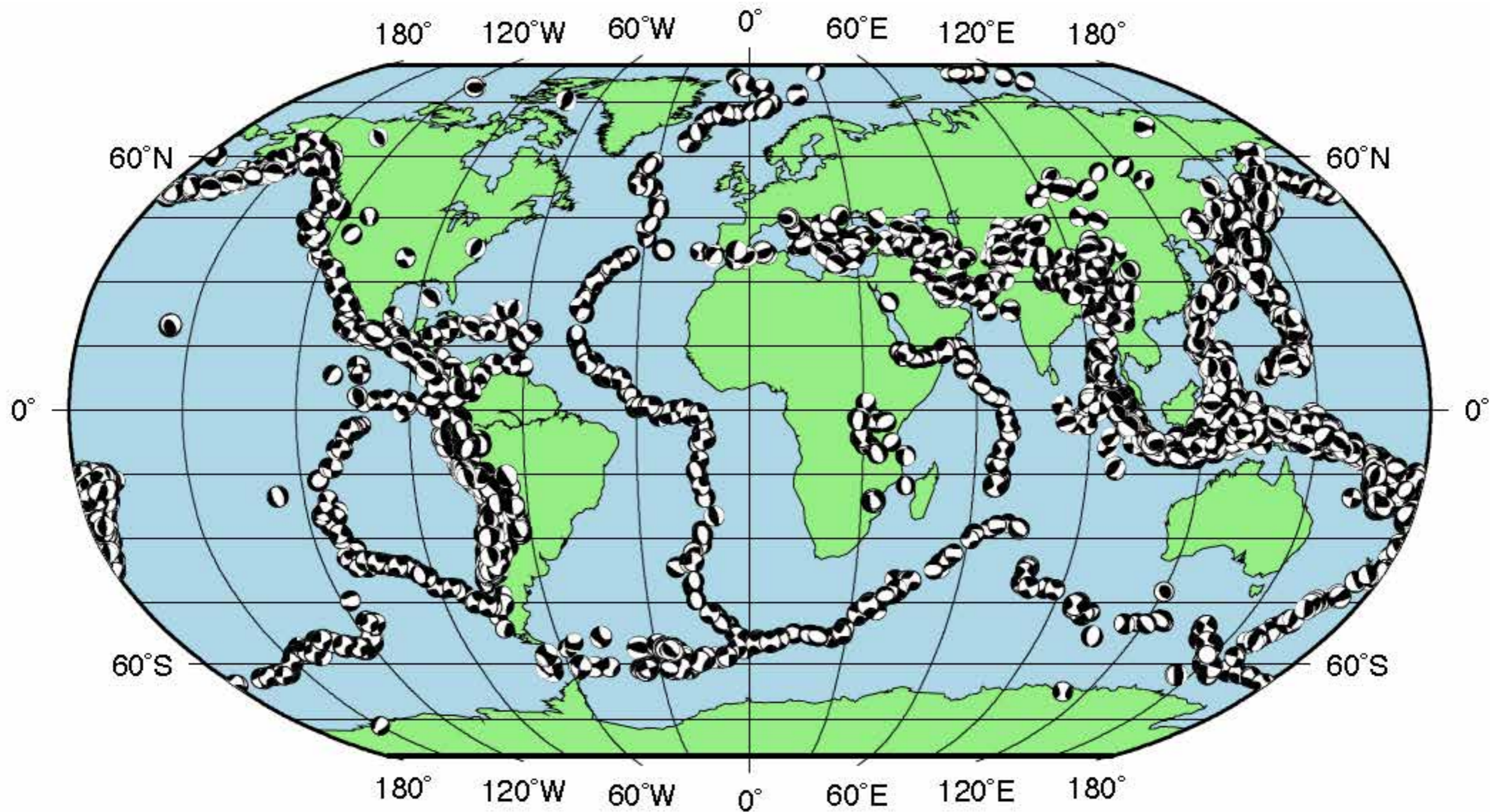
Station: KWAJ

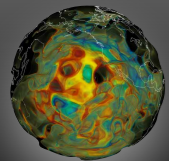
$\Delta = 52^\circ$



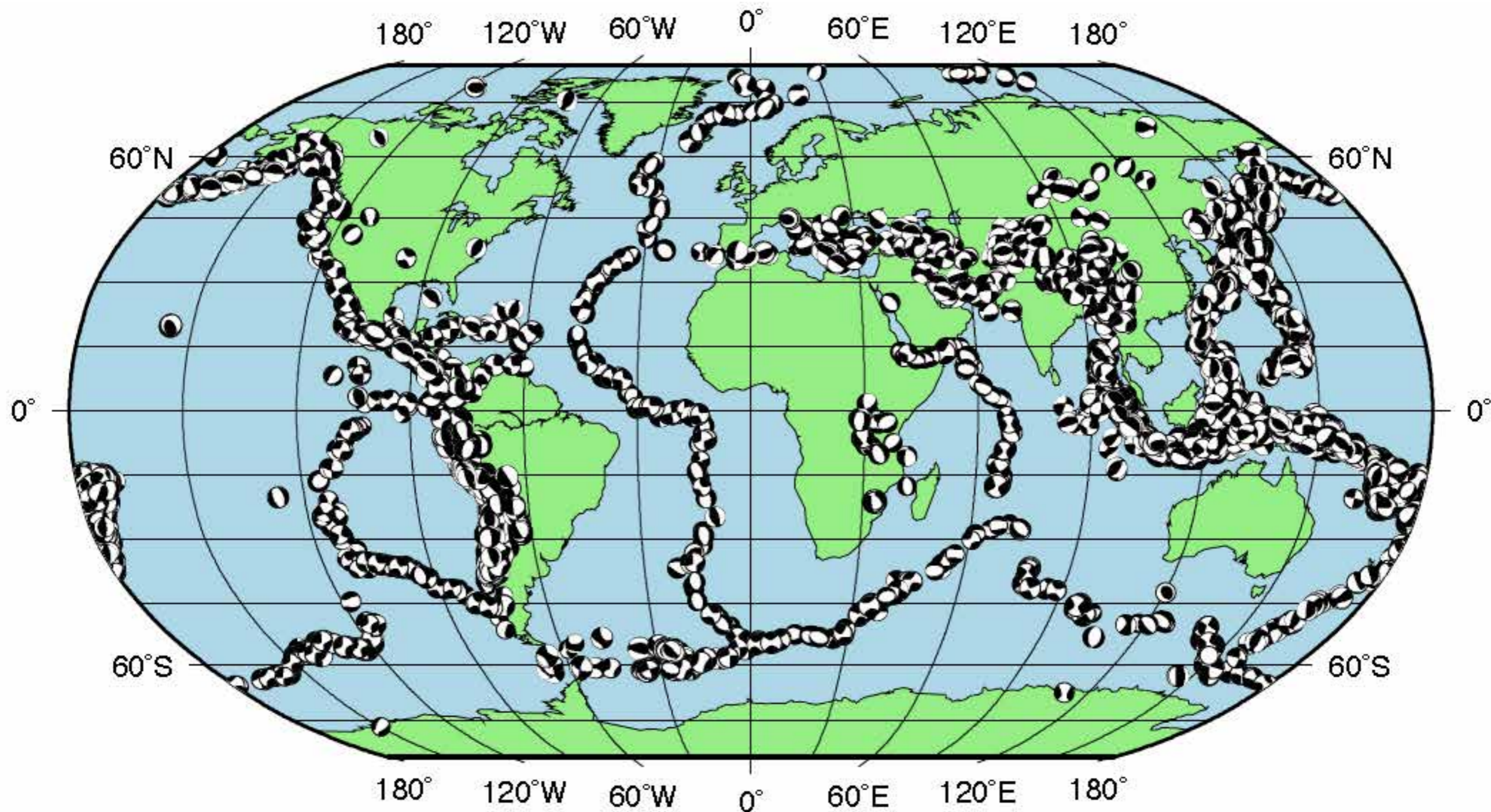


> 6000 global CMT events since 1999 ($5.5 \leq M_w \leq 7.0$)

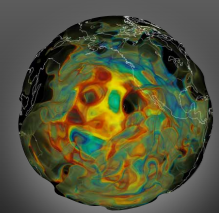




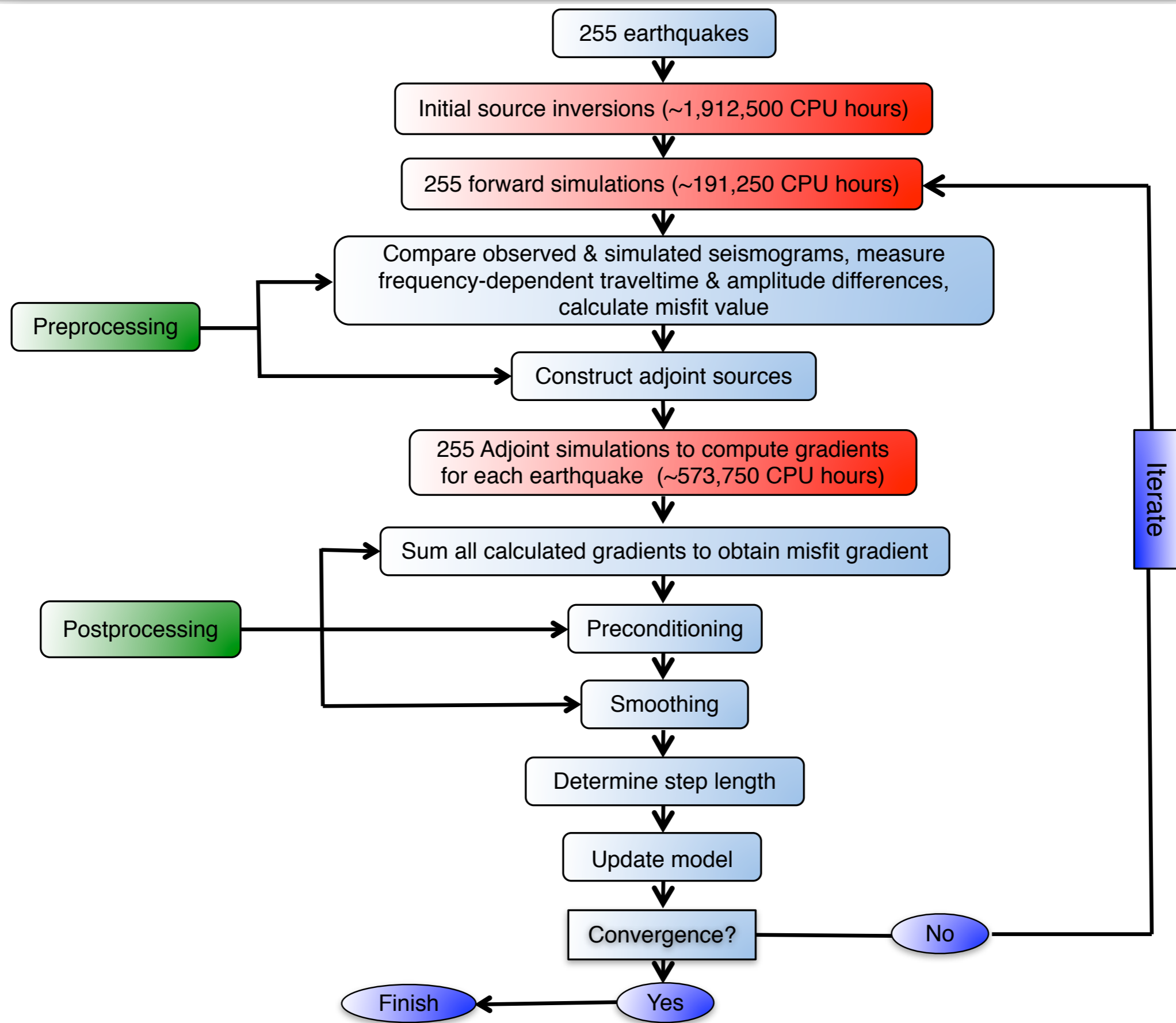
> 6000 global CMT events since 1999 ($5.5 \leq M_w \leq 7.0$)

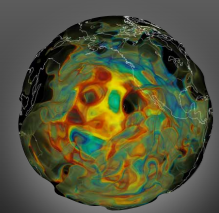


Use all available data!!

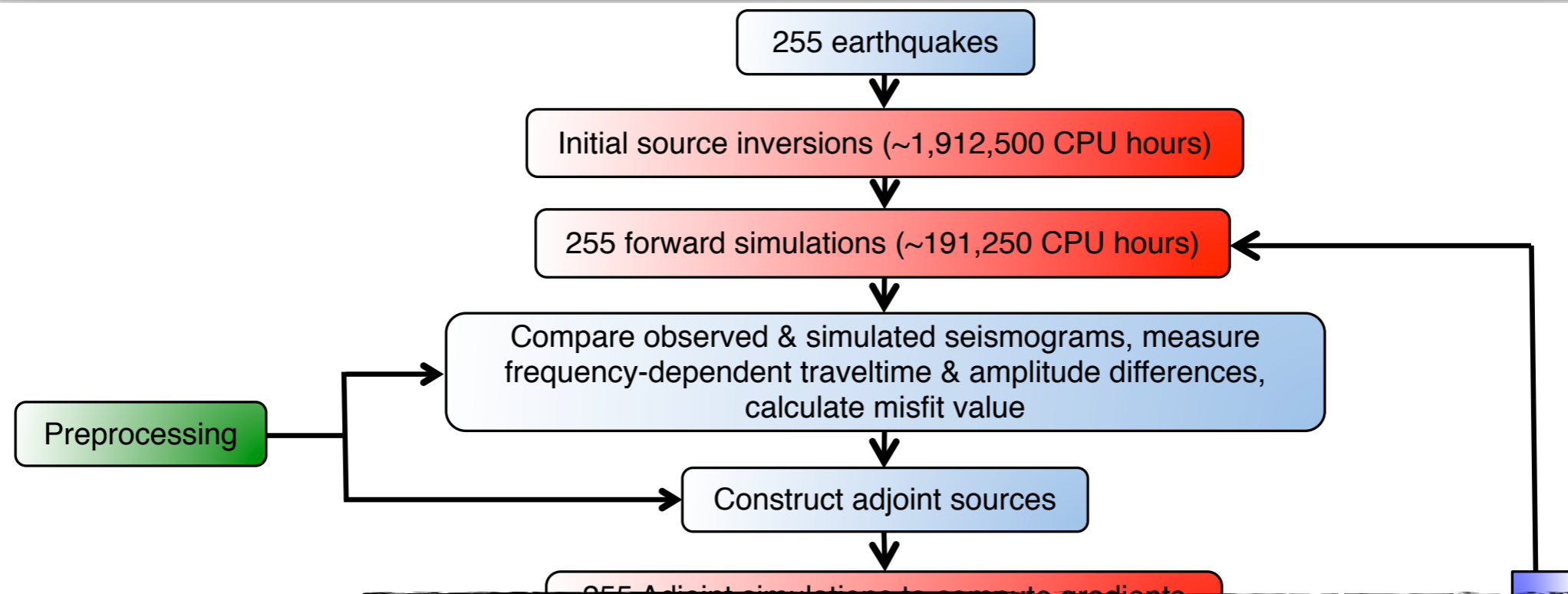


Adjoint tomography workflow





Adjoint tomography workflow



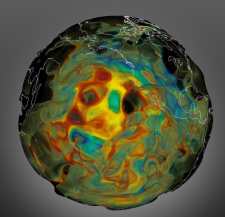
Bottlenecks:

I/O problems

- * mesher and solver outputs: mesh/kernel/model files, seismograms (SAC, ASCII, etc.), adjoint sources, ...
- * Pre-processing of SAC seismograms (can easily be *millions* of files), etc..

Data quality check for big data

Postprocessing



Data format



Adaptable Seismic Data Format

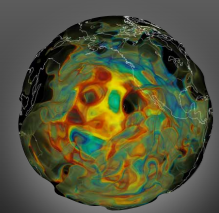


2013 R&D 100 Award Winner

ADIOS named one of top 100
tech products by R&D Magazine.

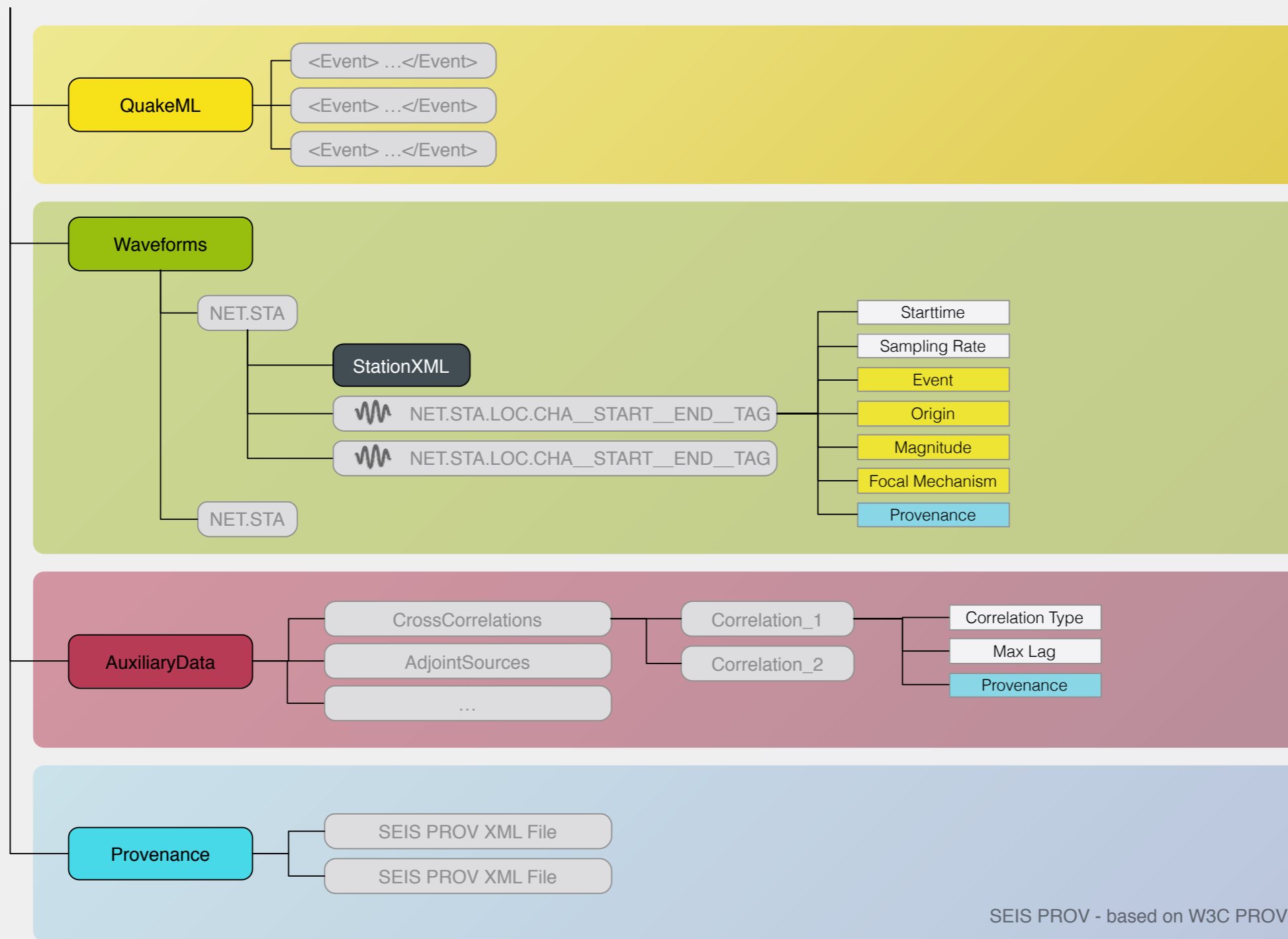


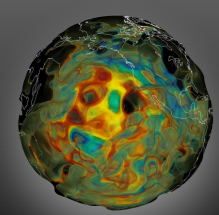
- **two ASDF files (observed & synthetics) per event**
 - **embarrassingly parallel pre-processing**
- **SPECFEM3D outputs: meshes / kernels / models, etc.**
 - **parallel post-processing**



ASDF (Adaptive Seismic Data Format)

CONTAINER (HDF5) => broad tool support, works on essentially all platforms of interest

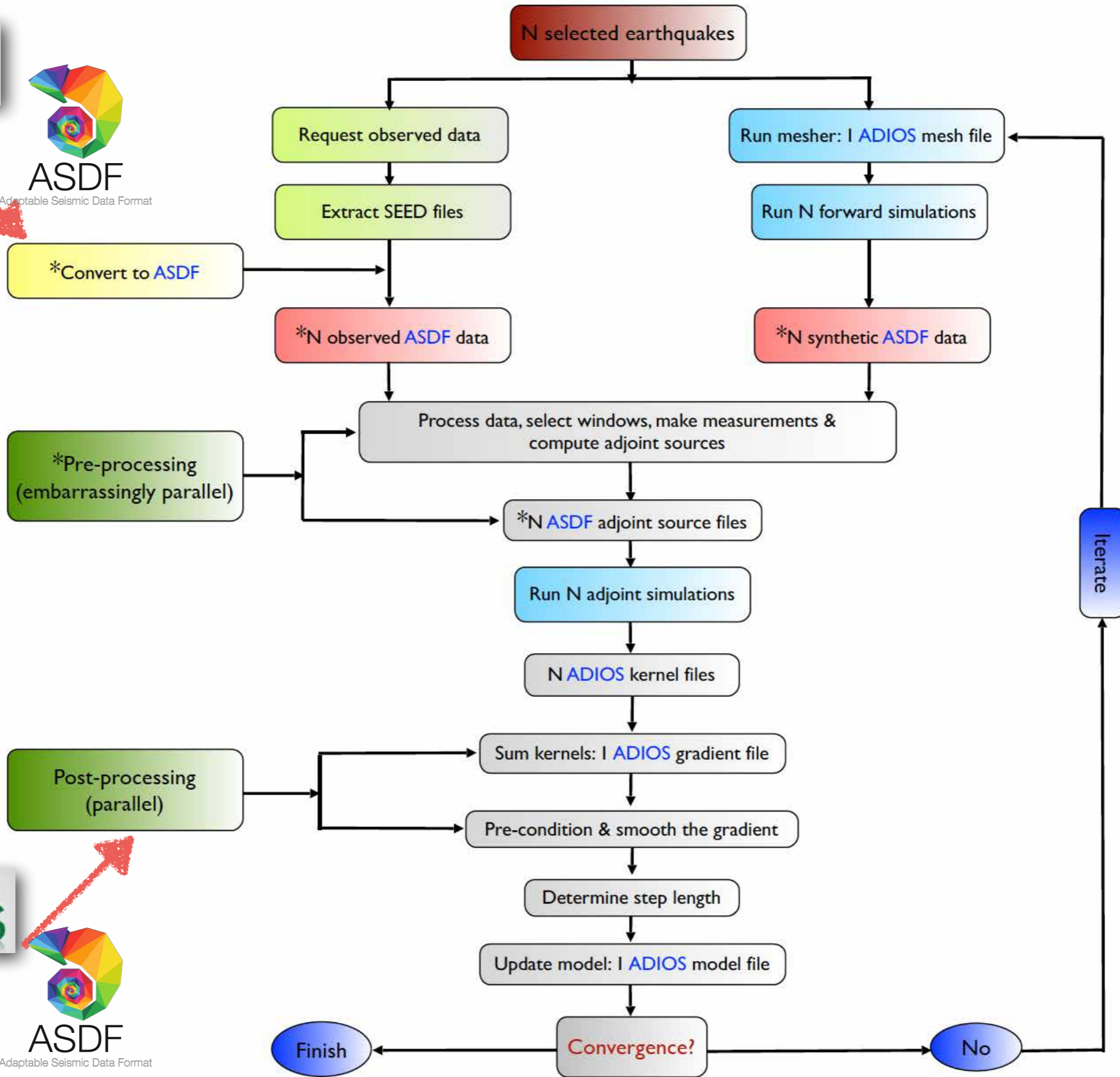




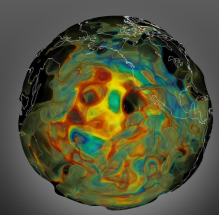
Adjoint tomography workflow



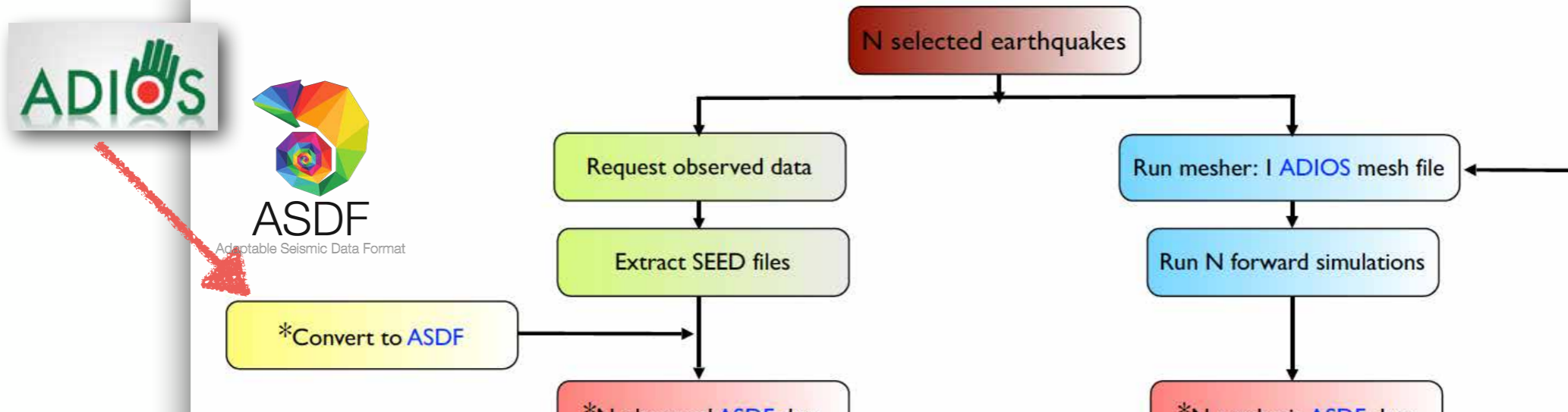
Adaptable Seismic Data Format



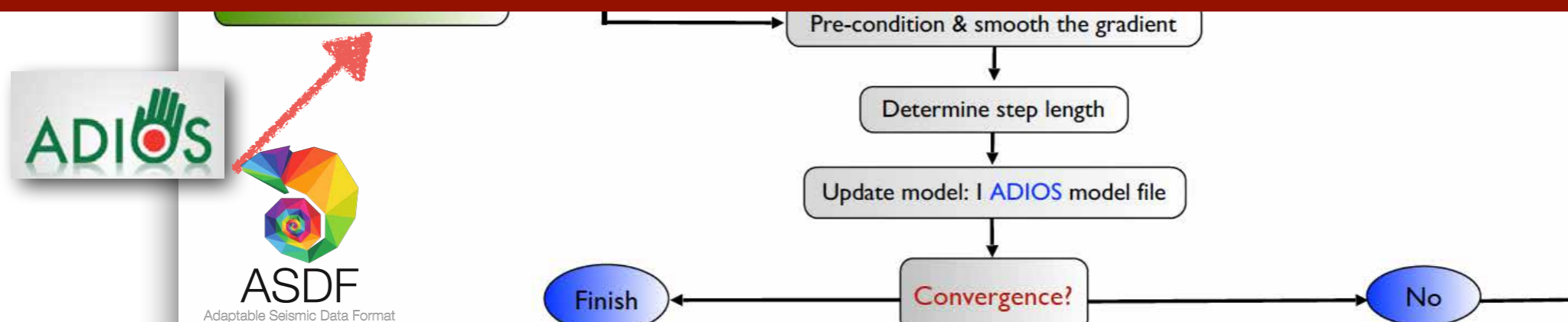
Adaptable Seismic Data Format

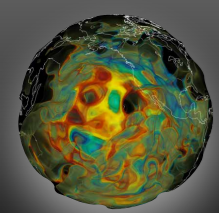


Adjoint tomography workflow



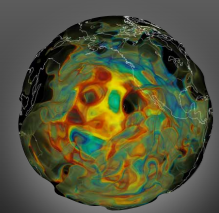
automatize the entire workflow:
 Radical EnTK workflow management environment
 Matthieu Lefebvre (Princeton) - Rutgers
 Balasubramanian et al. 2018





Conclusions / future directions

- **Next-generation models: Better parameterization**
 - Second-generation model: Azimuthal anisotropy**
 - Third-generation model: Anelasticity**
 - Internal topographies, crust, finite-source effects, etc...**
- **Increase the resolution of simulations, ultimately down to 1 Hz**
==> whole-Earth inversions
- **Approaching continental-scale resolution in global tomography**
in regions with good data coverage
- **Define better measurements / extract optimum information from**
each seismogram while keeping a balanced gradient
- **Resolution/uncertainty analysis ~ Model interpretation**

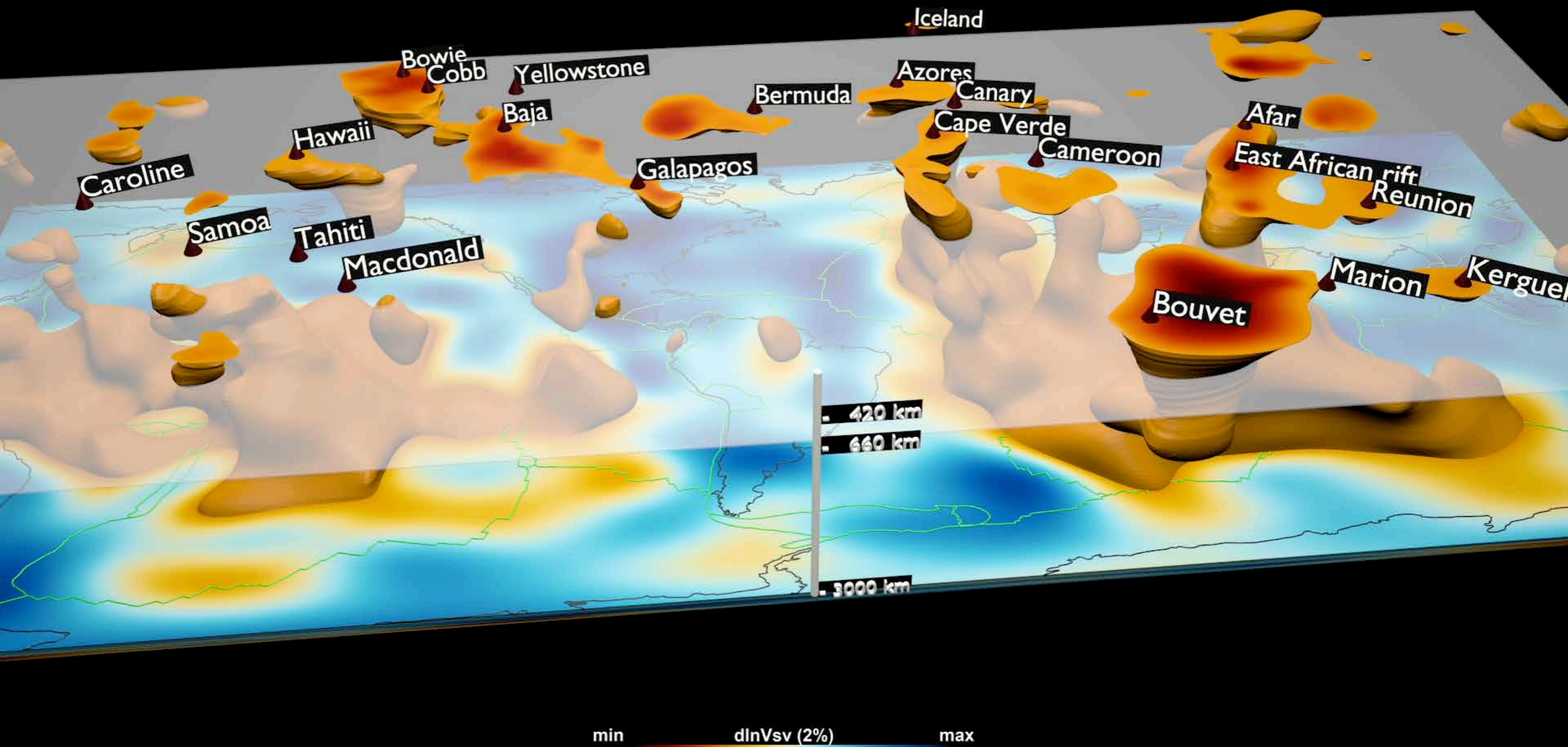


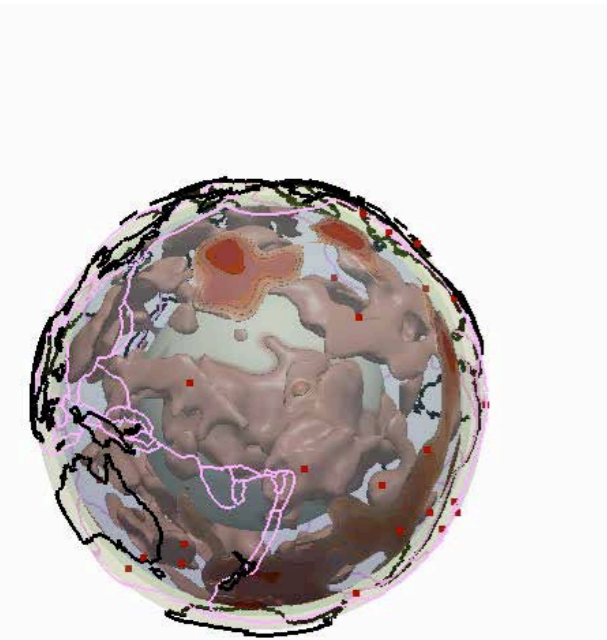
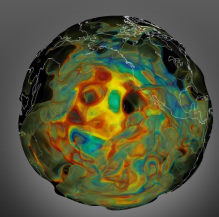
Visualizing Earth's mantle

GLAD-M15

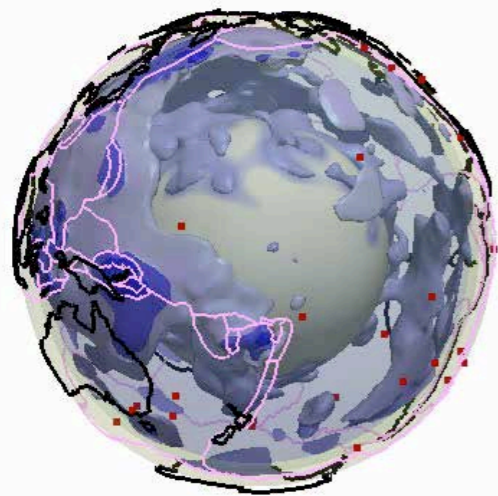
David Pugmire, ORNL

▲ Seismic Hotspots





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