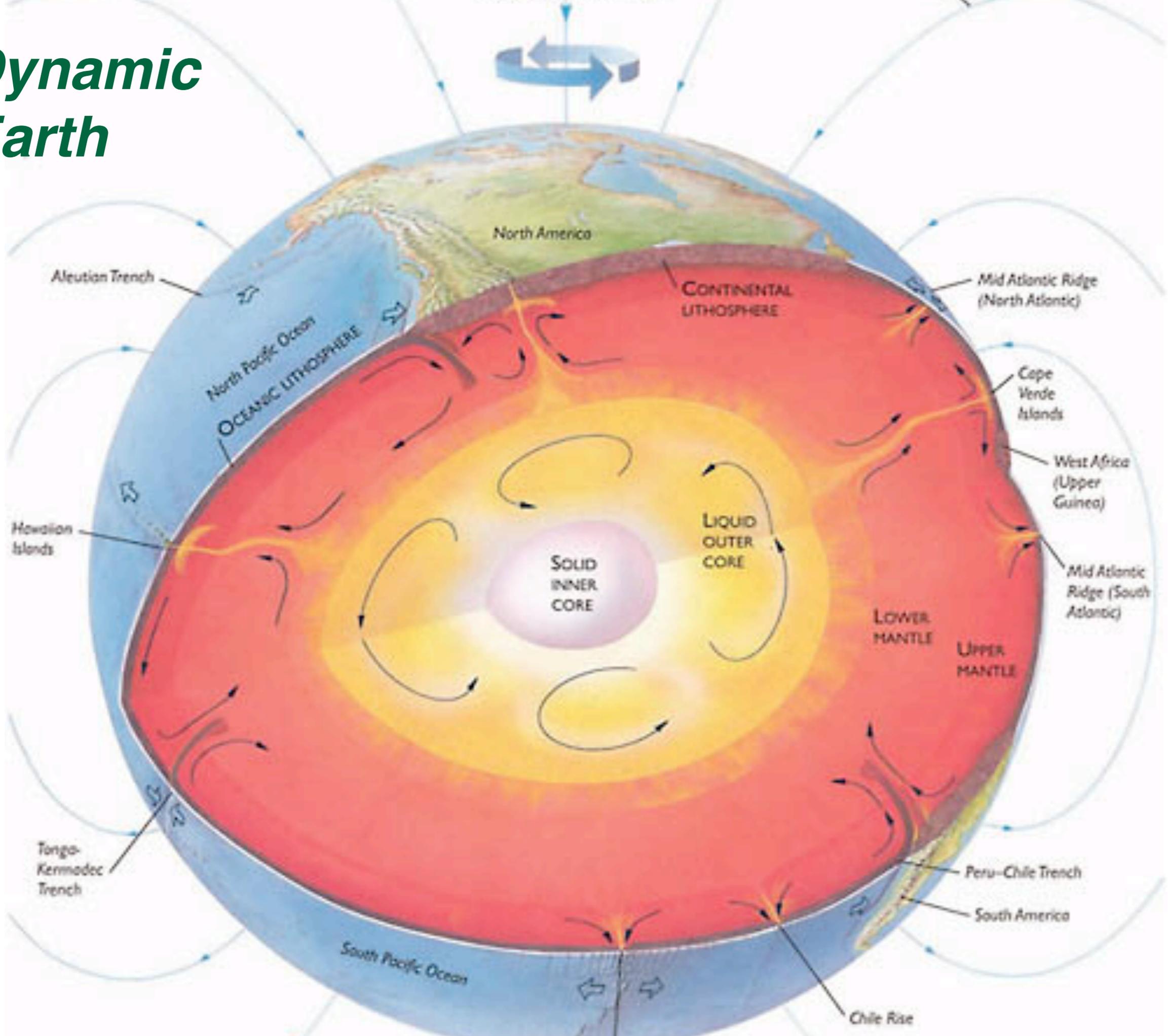
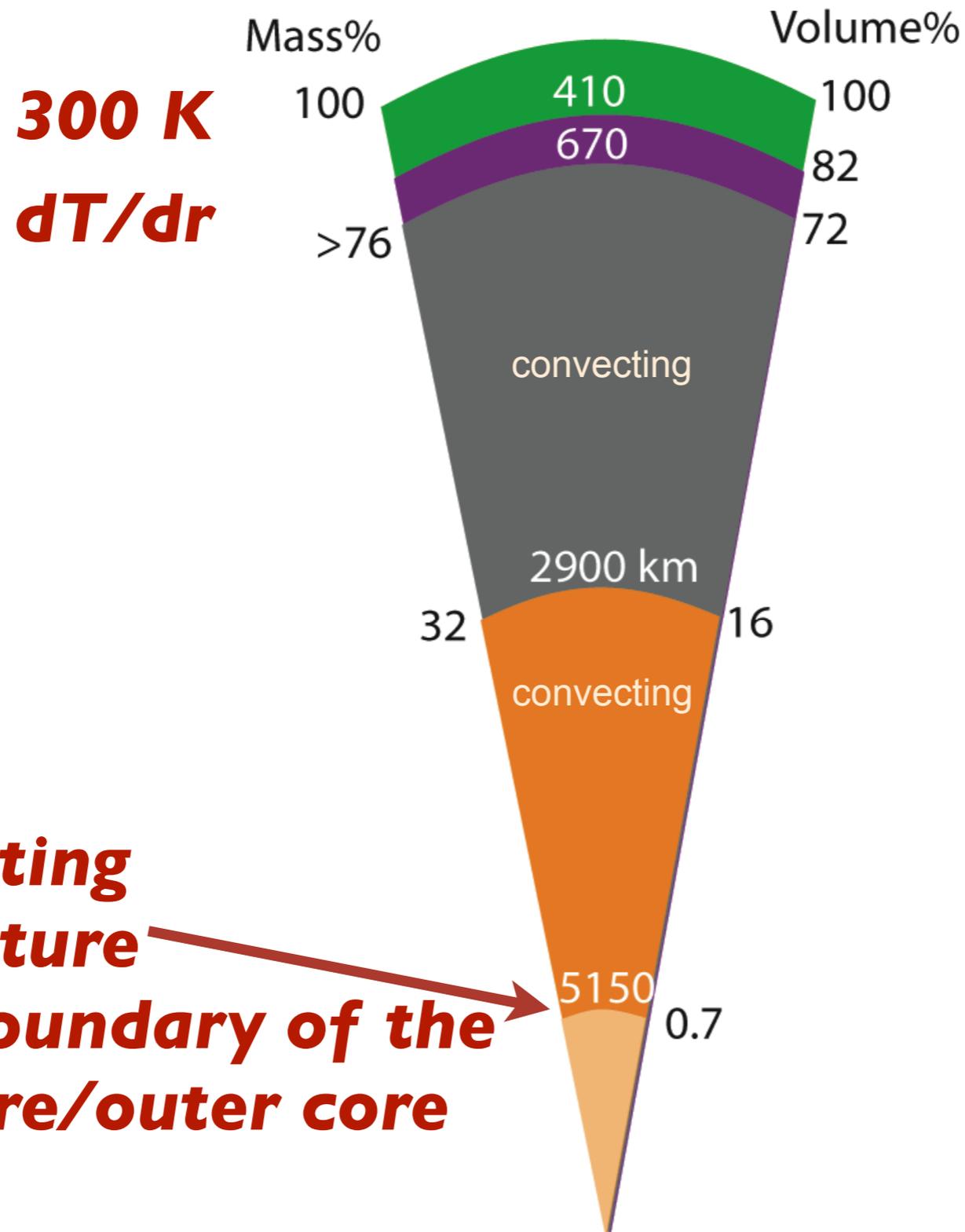


# Dynamic Earth



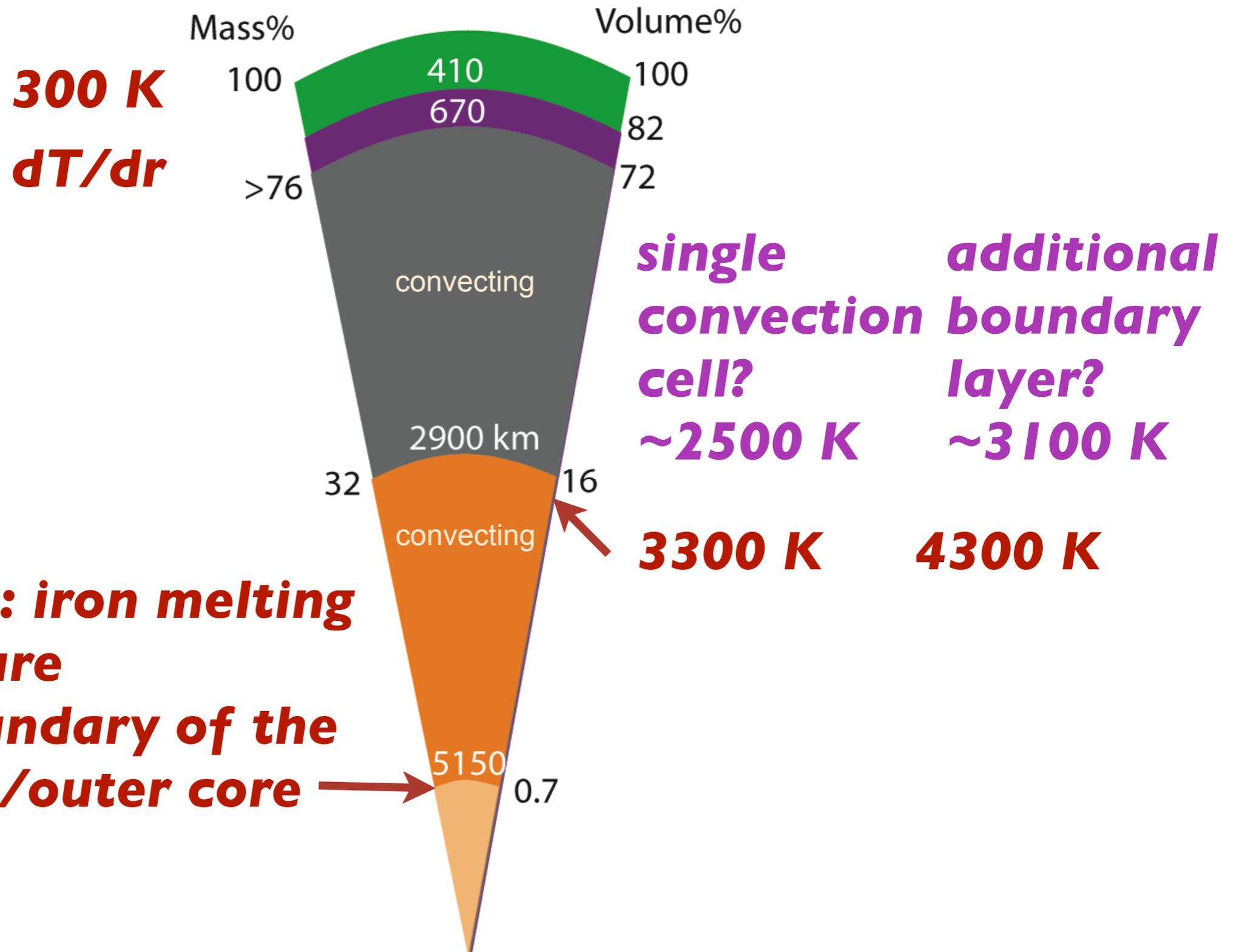
# Thermal State of the Earth

## Part I: Fixed Temperatures



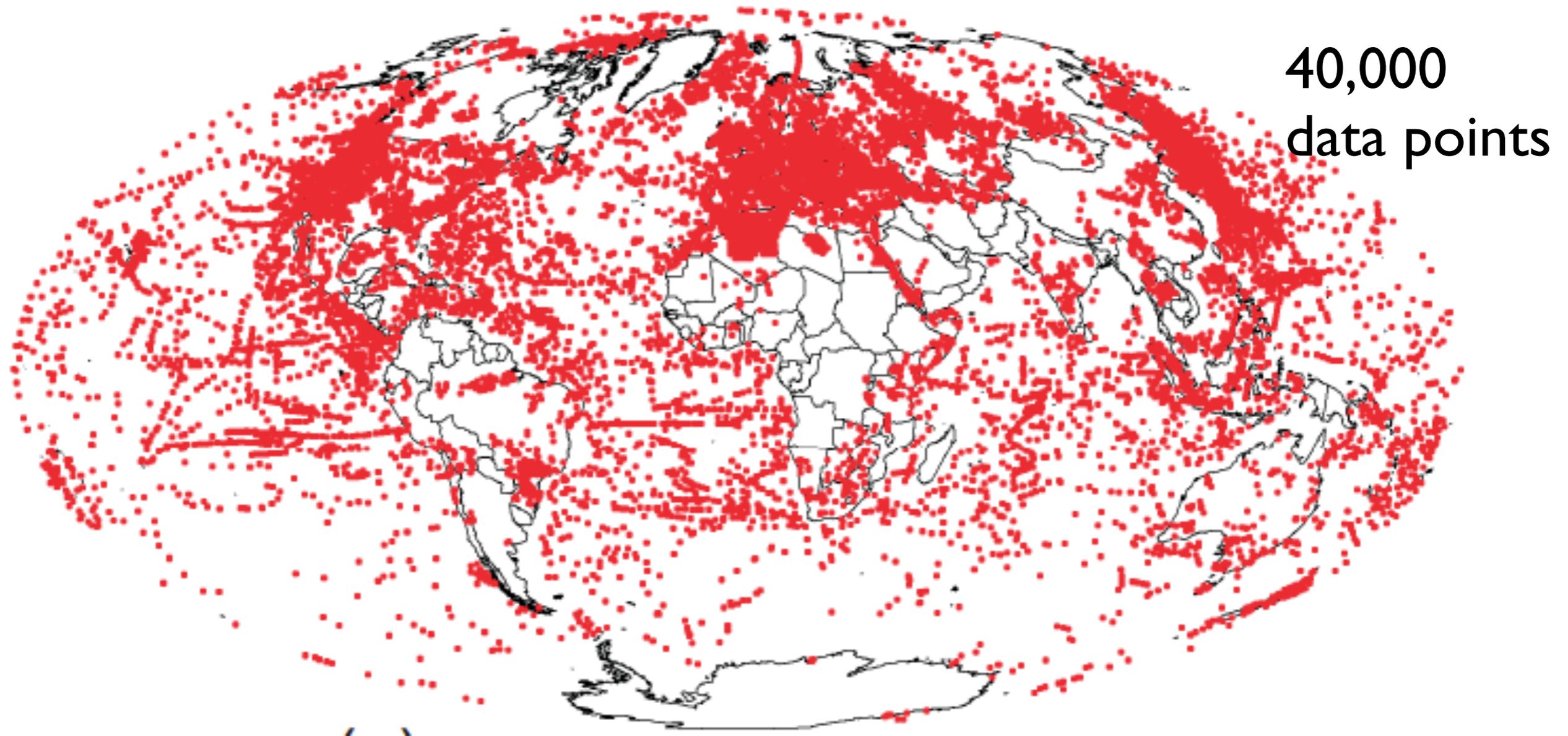
# Thermal State of the Earth

## Temperature inferences



# Thermal State of the Earth

## Part 2: Total heat flux constraints



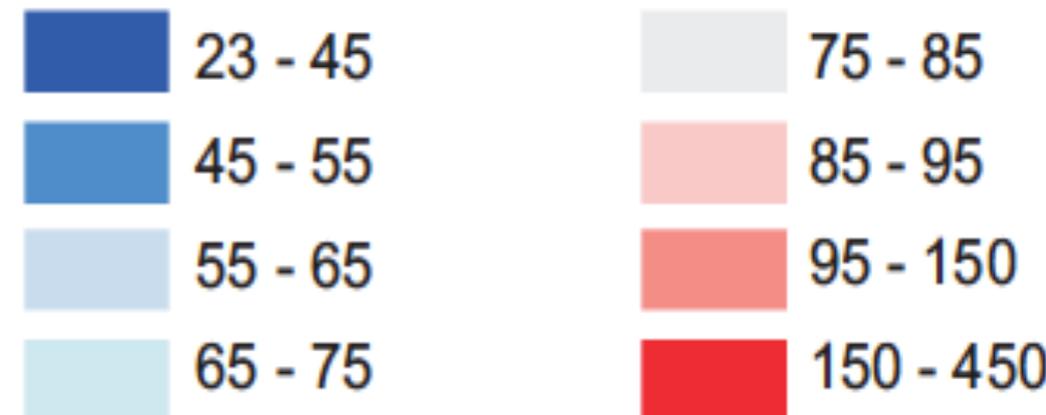
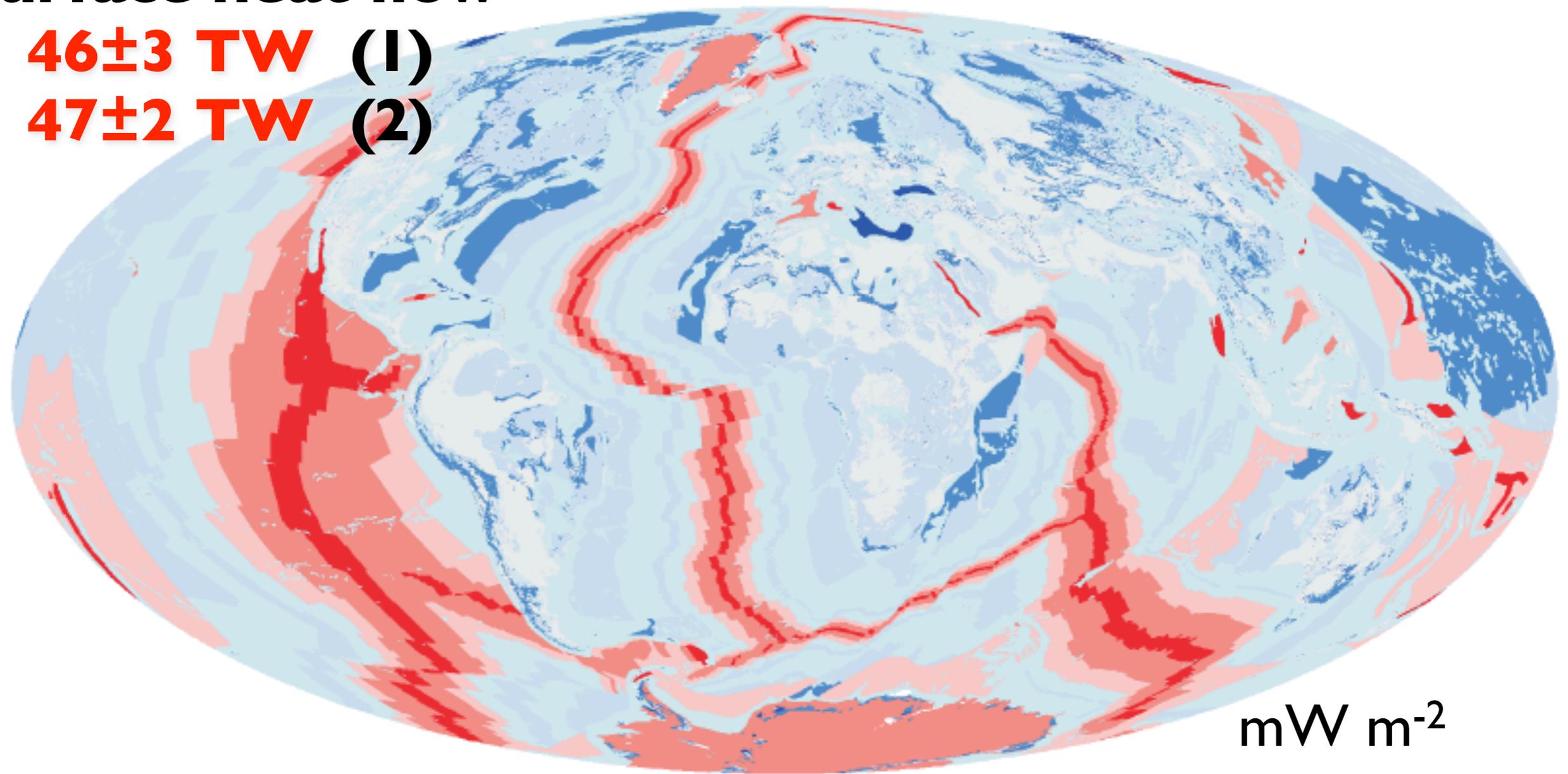
Conductive heat flow measured from borehole temperature gradient and conductivity

# Earth's Total Surface Heat Flow

## Surface heat flow

**46±3 TW (1)**

**47±2 TW (2)**



- (1) Jaupart et al (2008) *Treatise of Geophys.*  
(2) Davies and Davies (2010) *Solid Earth*

# Thermal State of the Earth

***~47(3) TW***

Radiogenic  
Contribution  
(U, Th, K)

Earth  
Interior  
Cooling

# Constraints on the thermal state of the Earth

***~47(3) TW***

Crustal  
Radiogenic  
Contribution  
***~7 TW***

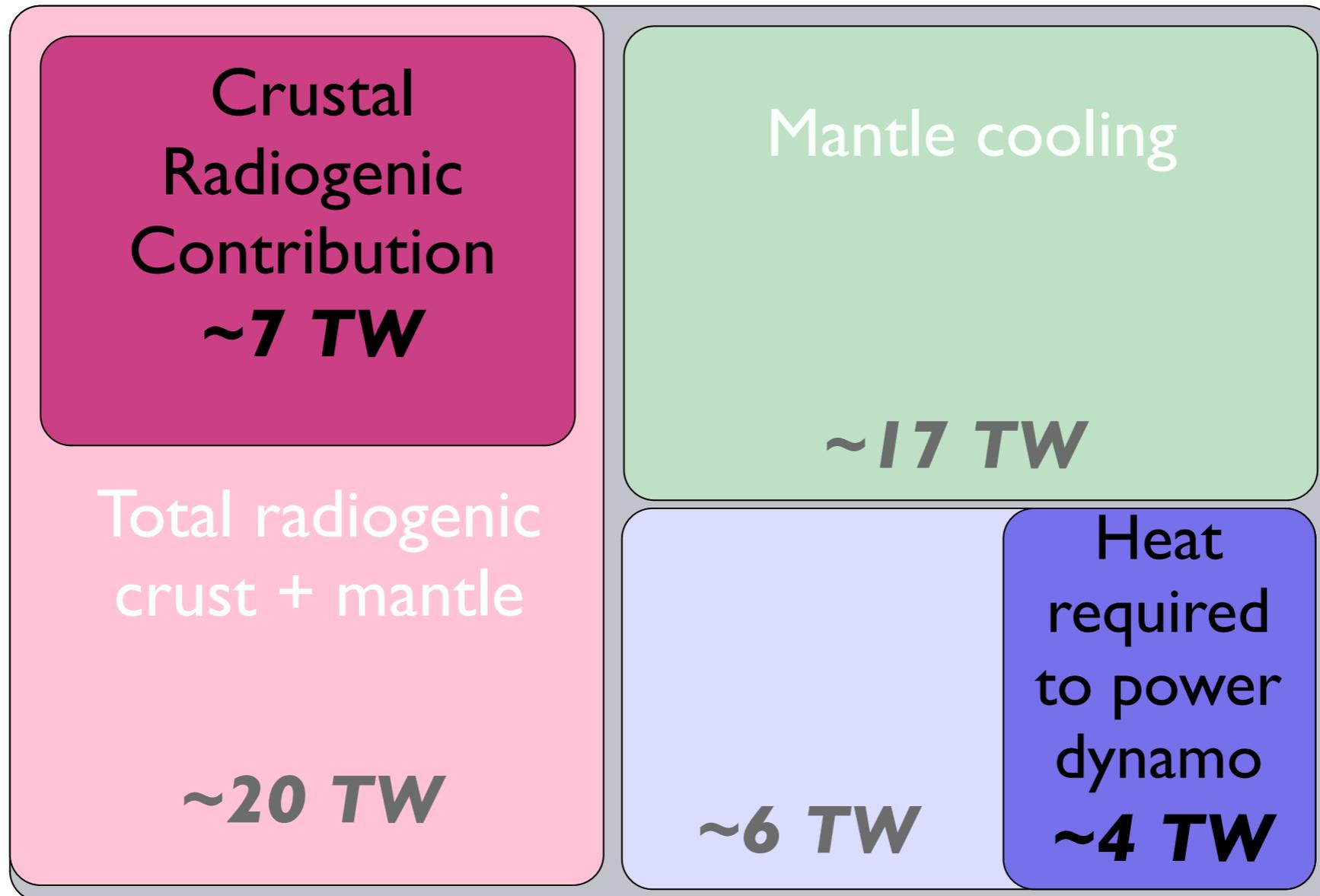
Heat  
required  
to power  
dynamo  
***~4 TW***

***values from McDonough & Sun,  
Chem. Geol., 120, 223-253, 1995***

***Buffett 2002***

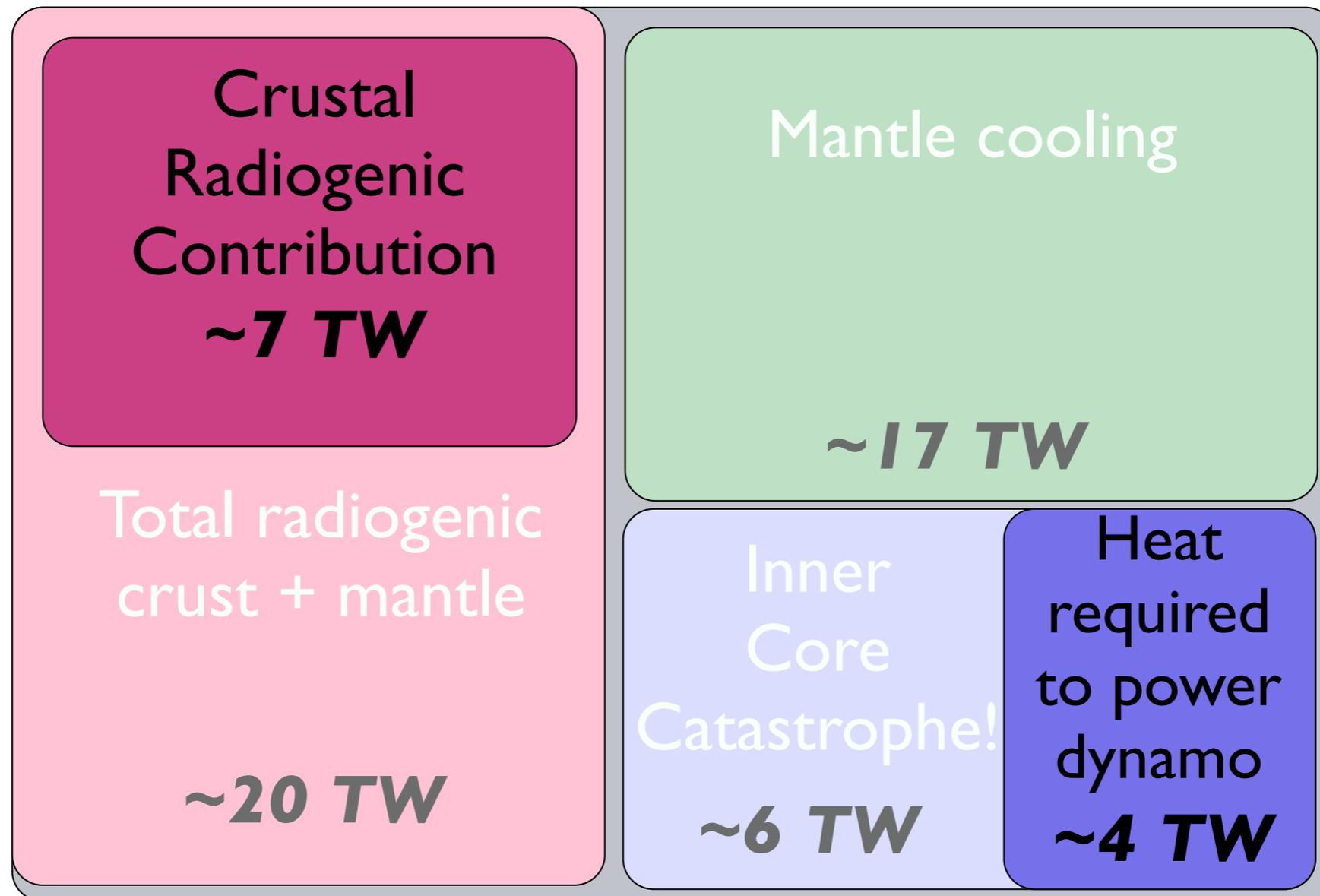
# Constraints on the thermal state of the Earth

**$\sim 47(3)$  TW**



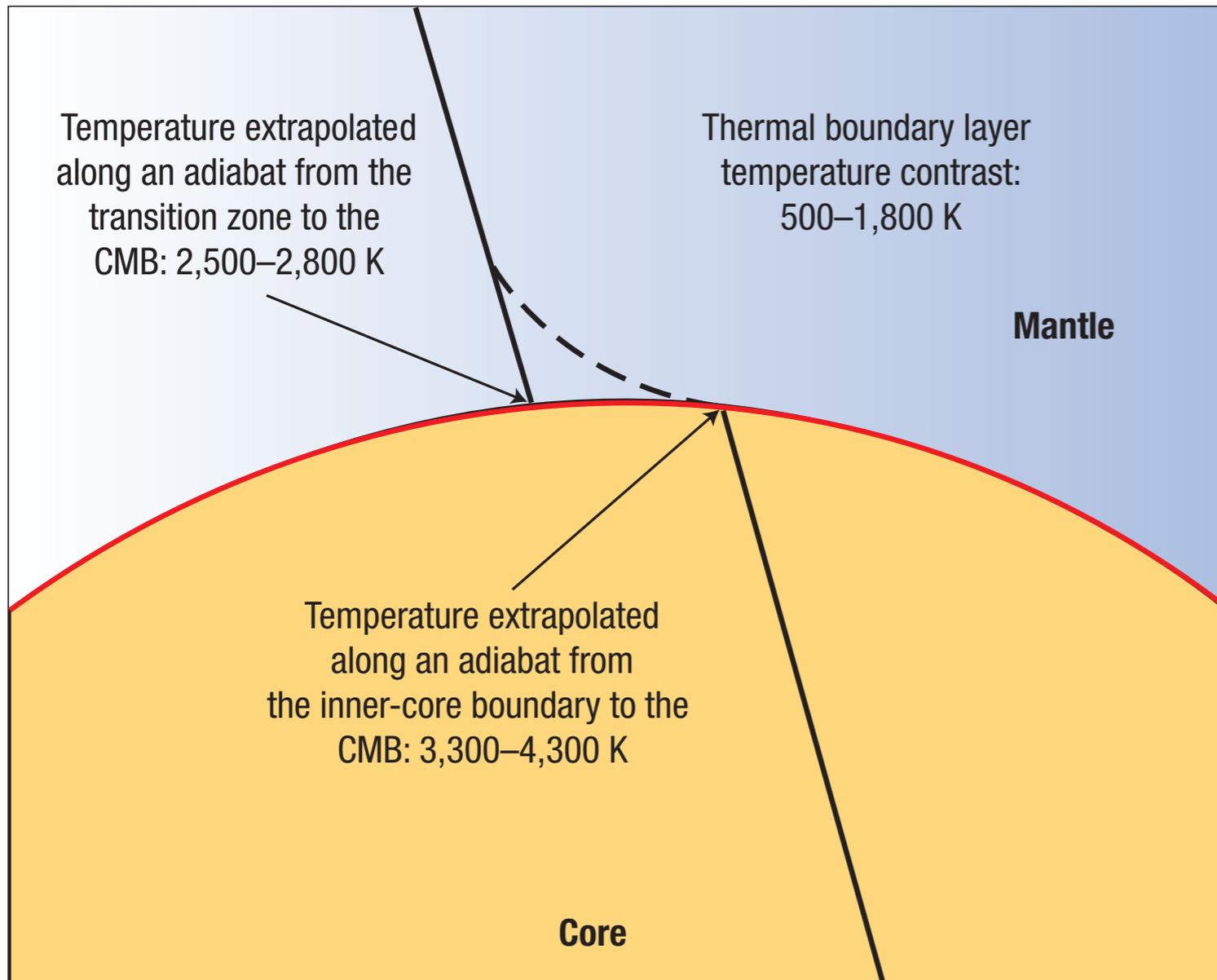
# Constraints on the thermal state of the Earth

***~47(3) TW***



***Buffett 2002--Inner core formed very recently/  
Early Temps were very hot!***

# Thermal Conductivity Definition



$$\text{Heat Flux} = K_{\text{cond}} \nabla T$$

**At core/mantle thermal boundary:**

**Length scale  $\sim 150(\pm 50)$  km**

**$\nabla T \sim 500\text{--}2000$  K**

**$K_{\text{cond}}$  very uncertain  
 $\sim 10 \pm 5$  W/m/K**

**Lay, Hernlund, Buffett,  
2008**

**Core/mantle heat flux  
uncertain  $\sim 4\text{--}20$  TW**

# Approaches to thermal conductivity of insulators

1. Determine phonon behavior

2. Measure heat flow via  $Q=dT/dx$

## THERMAL DIFFUSIVITY OF $MgSiO_3$ PEROVSKITE

Masahiro Osako

Section of Astronomy and Geophysics, National Science Museum, Tokyo

Eiji Ito

Institute for Study of the Earth's Interior, Okayama University

*Abstract.* Thermal diffusivity of  $MgSiO_3$  perovskite has been measured in the temperature range of 160 K to 340 K using a sample synthesized in a uniaxial split-sphere high-

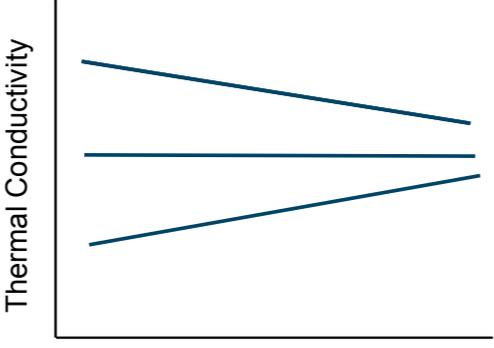
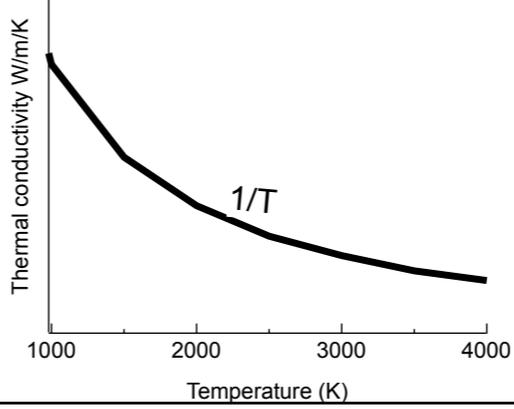
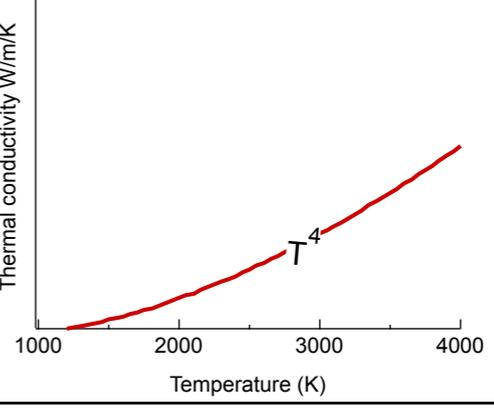
pressure apparatus. The results are compared with previous measurements and discuss the thermal conductivity under the lower mantle conditions.

and reaches to  $12 \text{ Wm}^{-1}\text{K}^{-1}$  in the v<sub>l</sub> mental

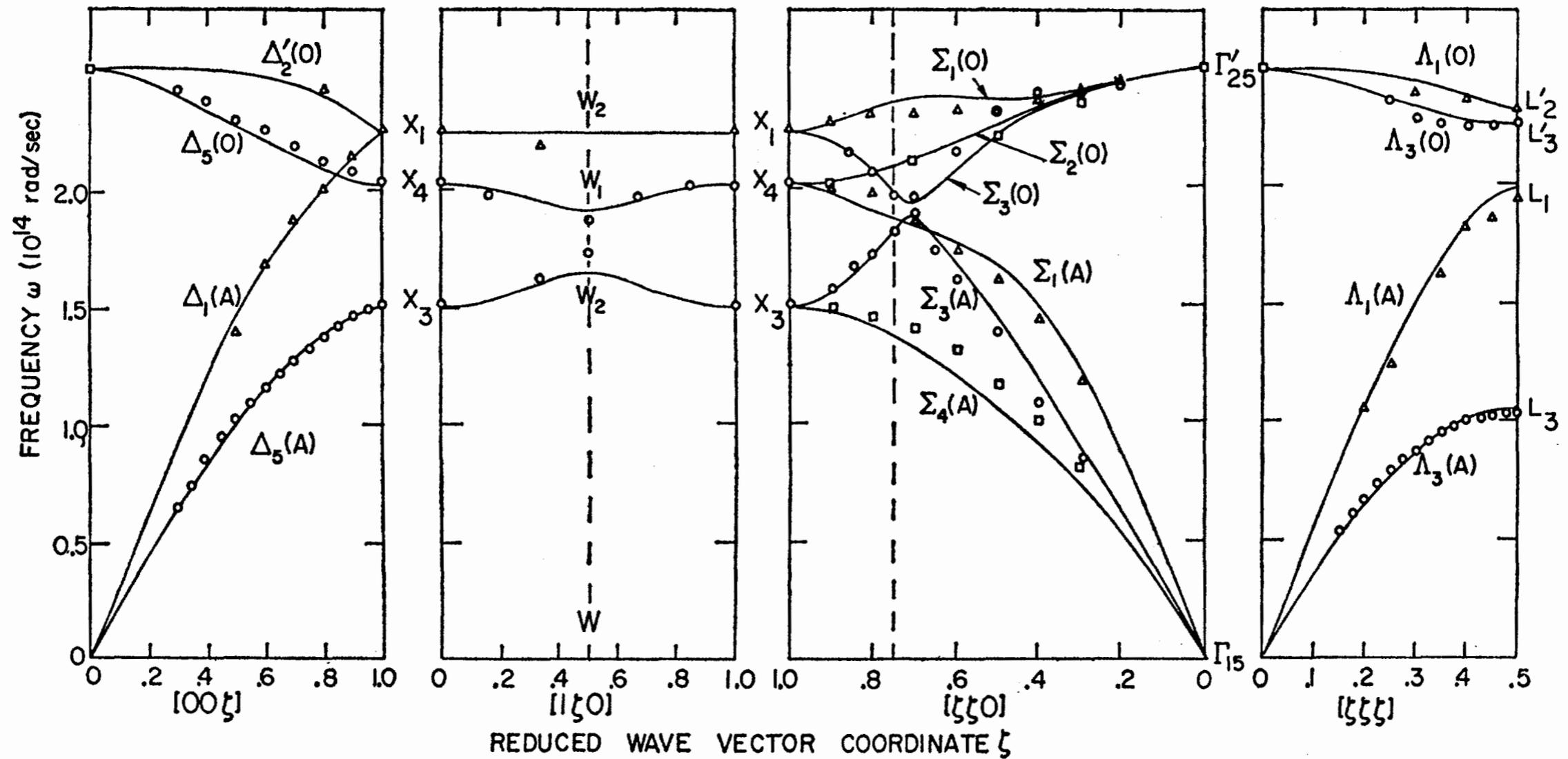
perovskite increases by a factor of 4 with depth throughout the lower mantle and reaches to  $12 \text{ Wm}^{-1}\text{K}^{-1}$  in the vicinity of the mantle-core boundary. The D'' layer might not be a thermal boundary layer insulating the high core-temperature, if this layer mainly consists of the perovskite.

was converted to  $MgSiO_3$  at 570 °C for one hour, using a uniaxial split-sphere high-pressure apparatus (USSA 5000) [Ito and Yamada, 1982]. The quenched perovskite sample, weighing 9.4 mg, had a slightly-distorted cylindrical shape and was of half transparency (Figure 1). The sample was confirmed to be a polycrystalline aggregate of  $MgSiO_3$  perovskite by a micro-focused X-ray diffractometer. For the

# Thermal Conductivity of Materials

<b>mechanisms:</b>	<b>electrons</b>	<b>atoms/ molecules</b>	<b>photons</b>
<b>Governing Equations</b>	$k\rho = LT$	$\kappa \sim C v^2 \tau_{eff}$	(solutions for radiative heat flow)
<b>Temperature dependence</b>	 <p>Thermal Conductivity vs Temperature</p>	 <p>Thermal conductivity W/m/K vs Temperature (K)</p>	 <p>Thermal conductivity W/m/K vs Temperature (K)</p>
<b>Examples (W/m/K)</b>	lead 35 iron 80 gold 310	asbestos 0.2-0.8 granite 2-4 Al <sub>2</sub> O <sub>3</sub> 30 diamond >1000	
<b>Trivia</b>	$k_{melt} < k_{solid}$	SiO <sub>2</sub> qtz 9.5/6.1 SiO <sub>2</sub> glass 1.46	goes in & out of fashion

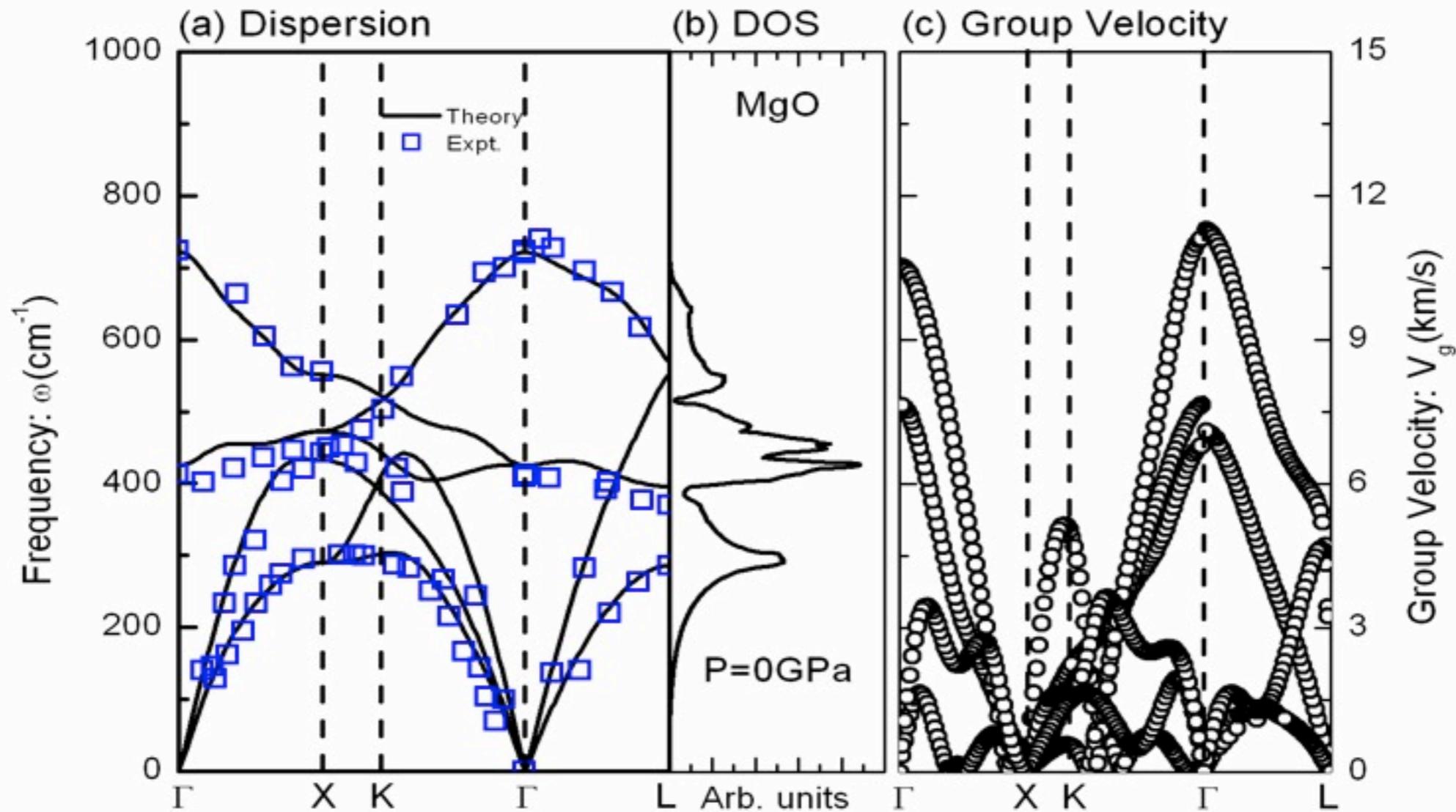
# Thermal Conductivity in Crystalline Insulators: Momentum transfer due to phonon interactions



**Phonon dispersion curve of diamond**  
**Warren et al., 1967**

# Approaches to thermal conductivity of insulators

1. Determine phonon behavior-experiment/theory
2. Measure heat flow directly via  $Q=dT/dx$

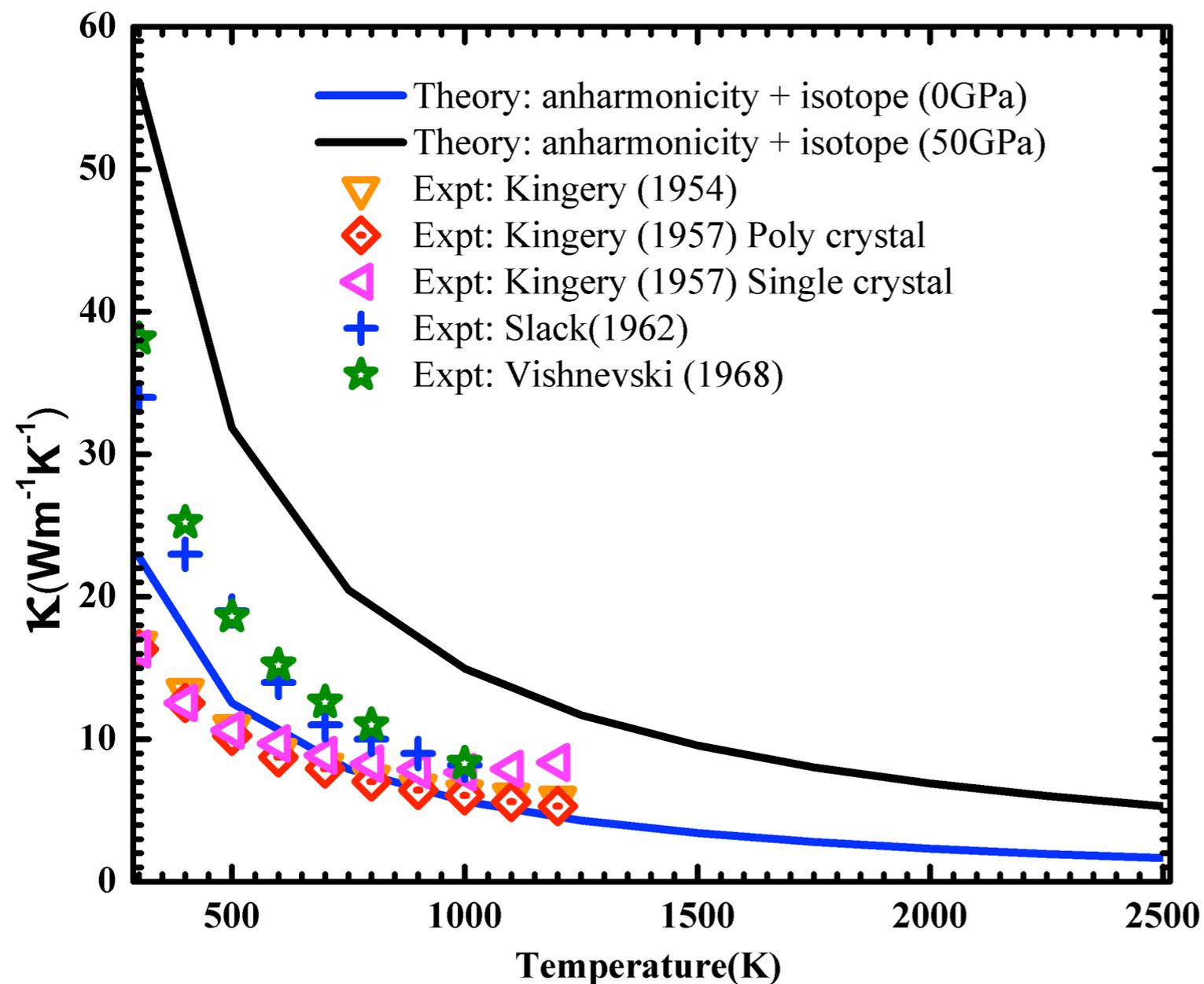


**MgO**

**Tang & Dong, 2010**

# Approaches to thermal conductivity of insulators

1. Determine phonon behavior-experiment/theory
2. Measure heat flow directly via  $Q=dT/dx$



***Thermal conductivity of MgO shows good agreement between theory and experiment***

***Tang & Dong 2010  
de Koker 2009***

# ***Thermal Conductivity of $MgSiO_3$ -pv and Heat Transport in the Earth's Mantle***

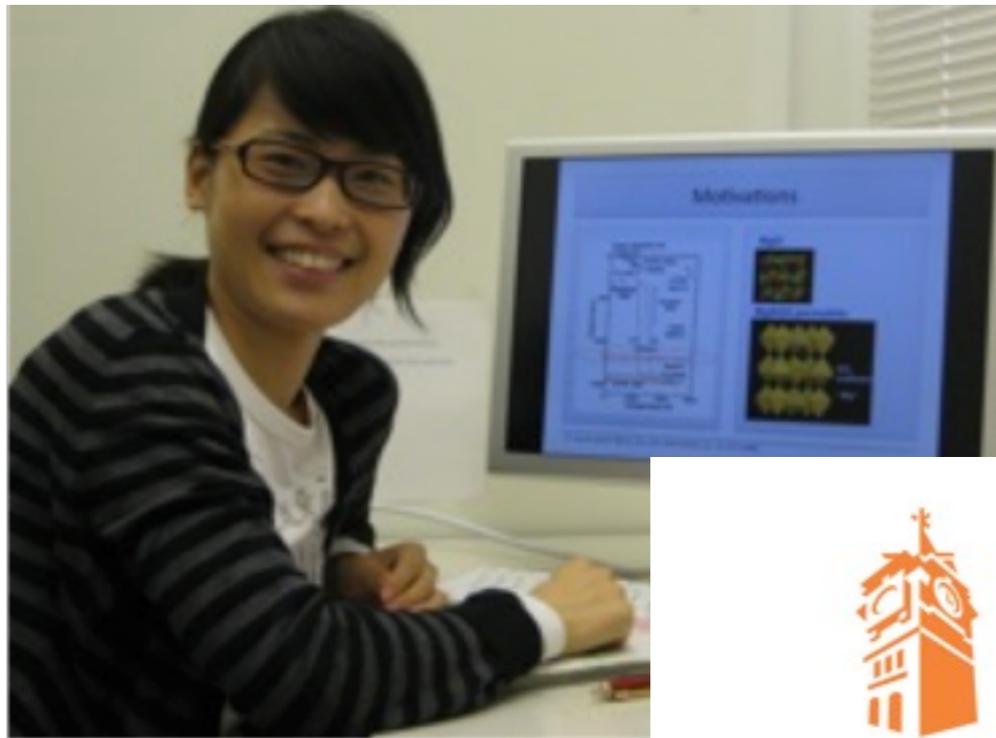
Xiaoli Tang (Now at: CalTech, Applied Physics)

Abby Kavner (Earth & Space Science, UCLA)

Jianjun Dong (Physics, Auburn University)

Laurent Pilon (Mech E., UCLA)

Emma Rainey (Earth & Space Science, UCLA)



# Approach: Solve the Peierls-Boltzmann kinetic transport equation for phonons

$$\kappa = \frac{1}{3} \sum_{i, \vec{q}} c_V(i, \vec{q}) v_g^2(i, \vec{q}) \tau(i, \vec{q})$$

**Lattice thermal conductivity**  
(W/m/K)

**Heat capacity**  
J/K/m<sup>3</sup>

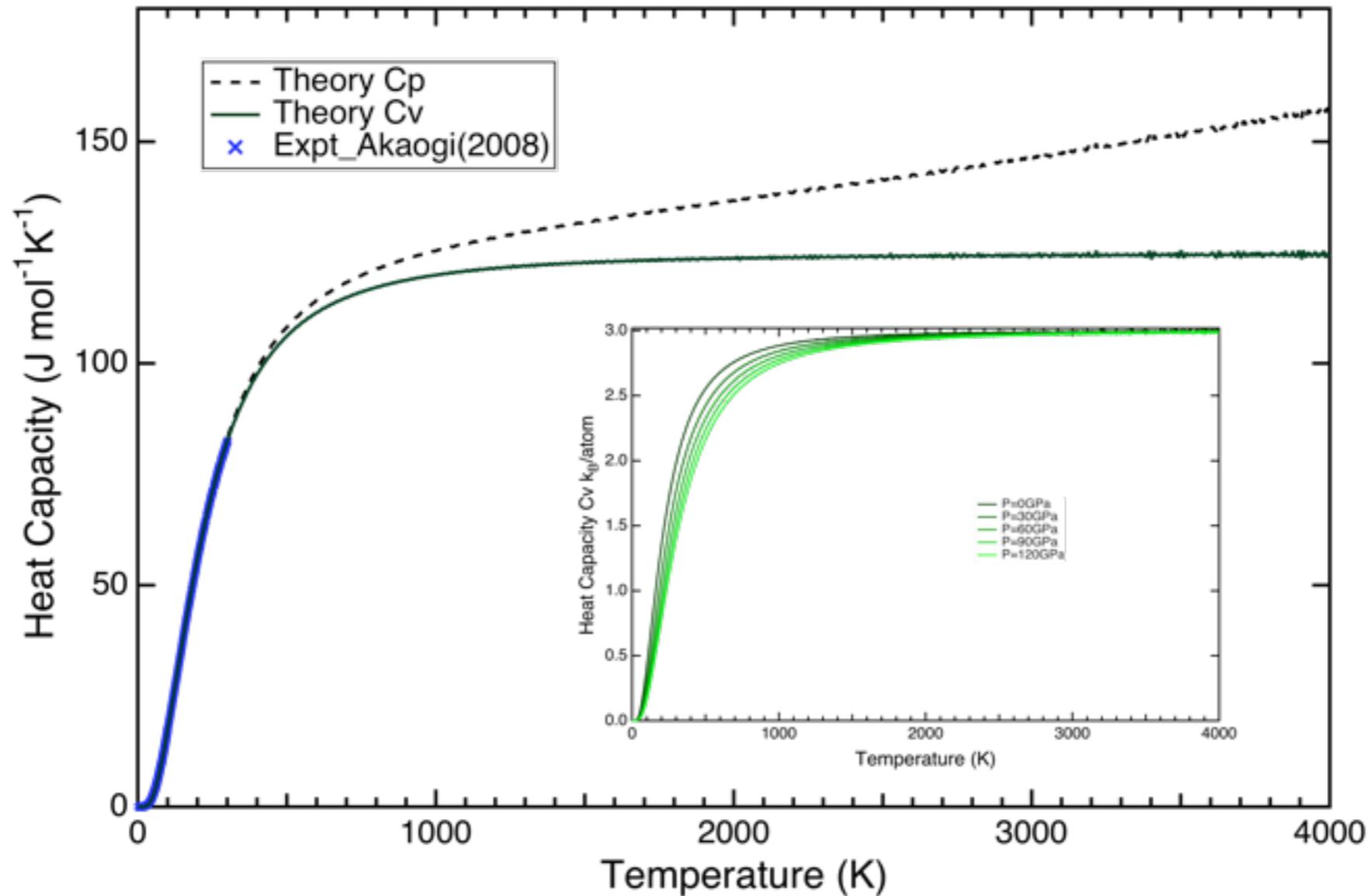
**Phonon group velocity**  
(m/s)<sup>2</sup>

**Phonon lifetime**  
s

***Integrate over all phonon modes in three dimensions***

# Heat Capacity $C_V$

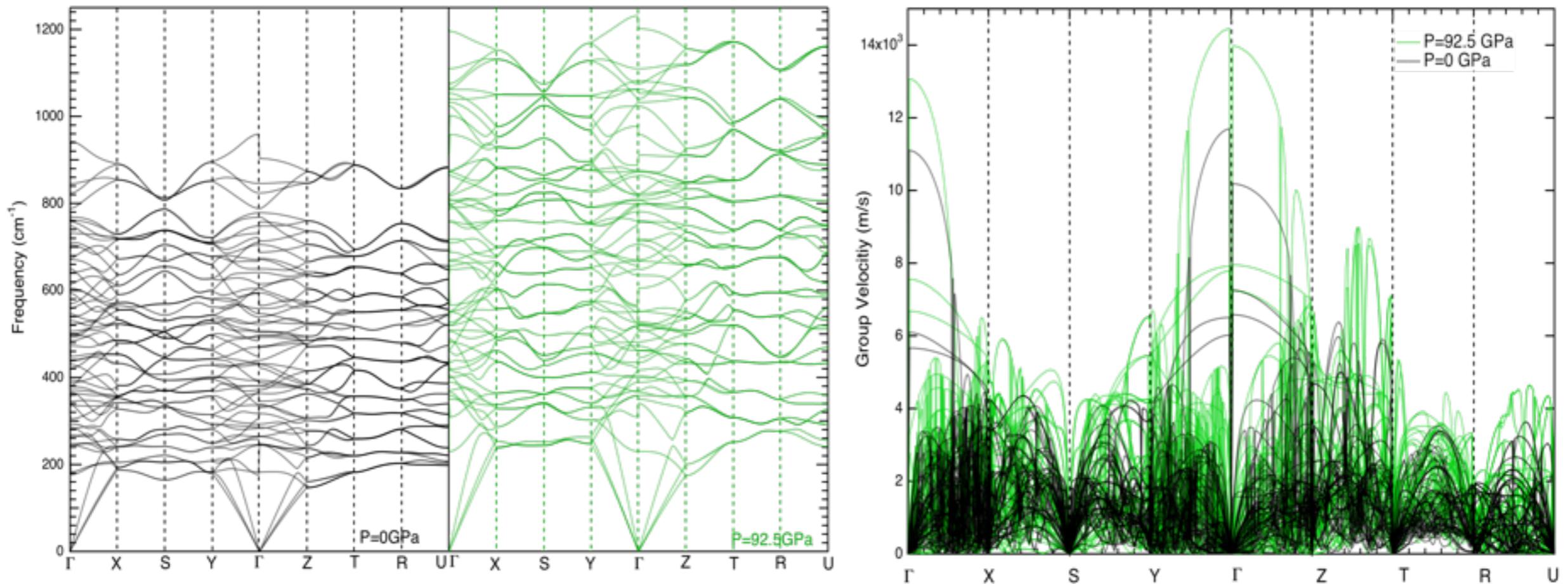
$$k = C_V v_g^2 \tau$$



- Heat capacity only varies in low-temperature range ( $T < \theta_D$ ).
- Heat capacity changes insignificantly with pressure.

# Phonon Group Velocity $V_g$

$$k = C_v \boxed{v_g^2} \tau$$



- Phonon frequencies increase as pressure goes up.
- Overall phonon group velocity also increases at higher pressure.
- Mode-by-mode group velocities have large variability

# Phonon Life Time $\tau_{eff}$

$$k = C_v v_g^2 \tau$$

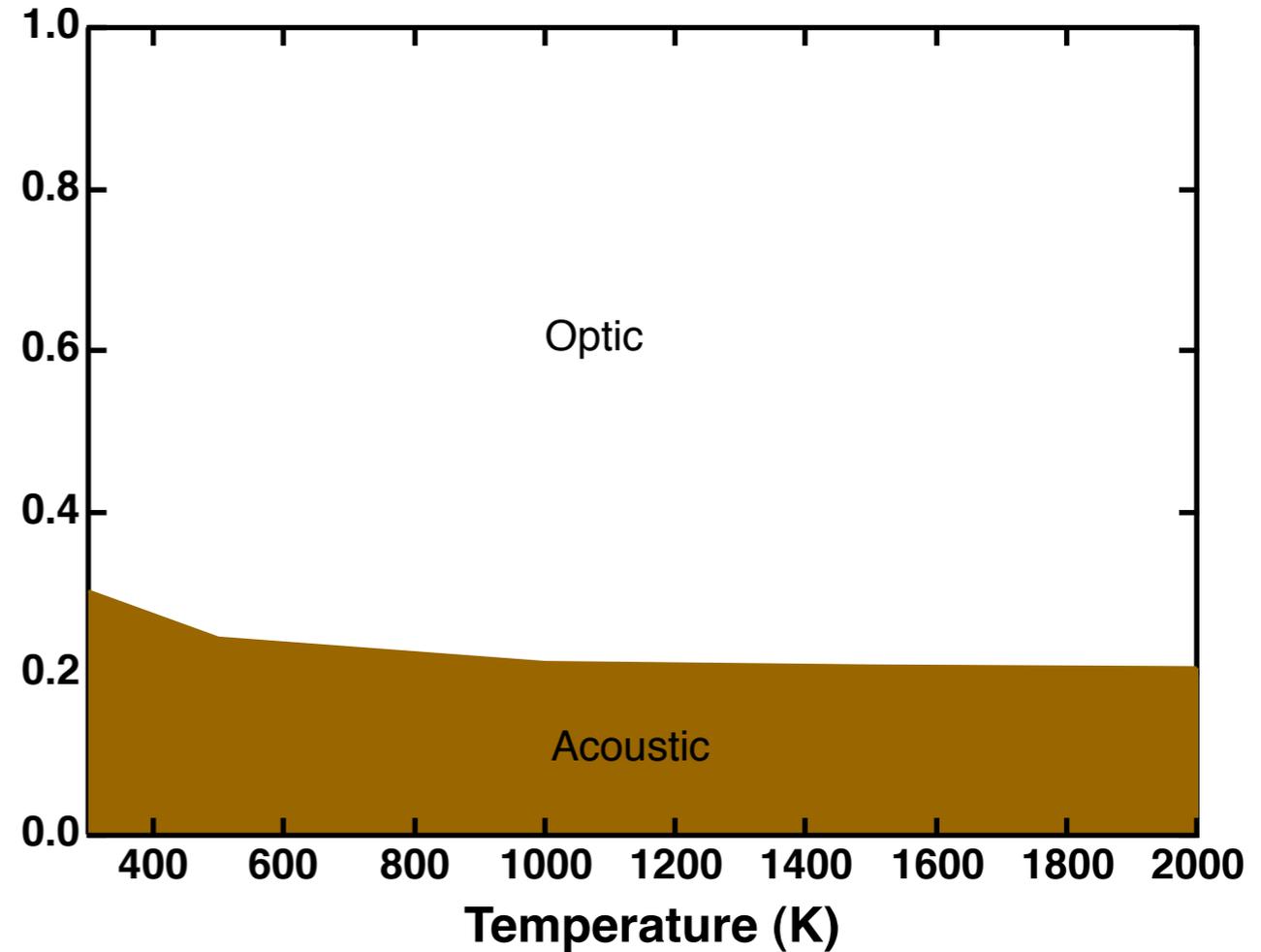
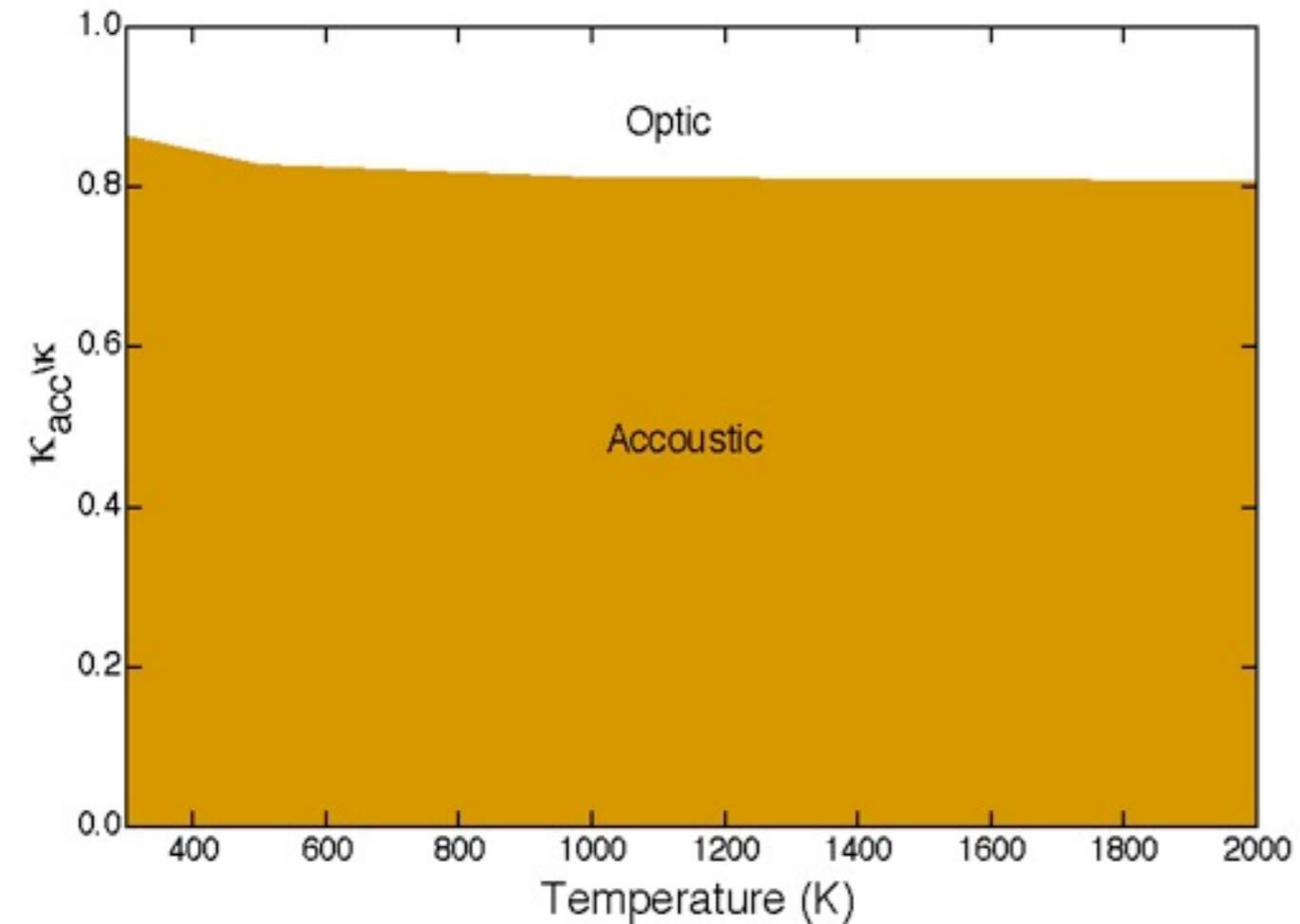
MgSiO<sub>3</sub>: (3 acoustic + 57 optic) v.s MgO: (3 acoustic + 3 optic)

<i>Typical</i> $\tau_{eff}$	$\tau_{eff}$ (acoustic) T=300K (ps)	$\tau_{eff}$ (acoustic) T=3000K (ps)	$\tau_{eff}$ (optic) T=300K (ps)	$\tau_{eff}$ (optic) T=3000K (ps)
<b>MgSiO<sub>3</sub></b>	1.7	0.14	0.16	0.008
<b>MgO</b>	8.5	3.03	0.91	0.34

Phonon lifetimes are much smaller in MgSiO<sub>3</sub> than in MgO

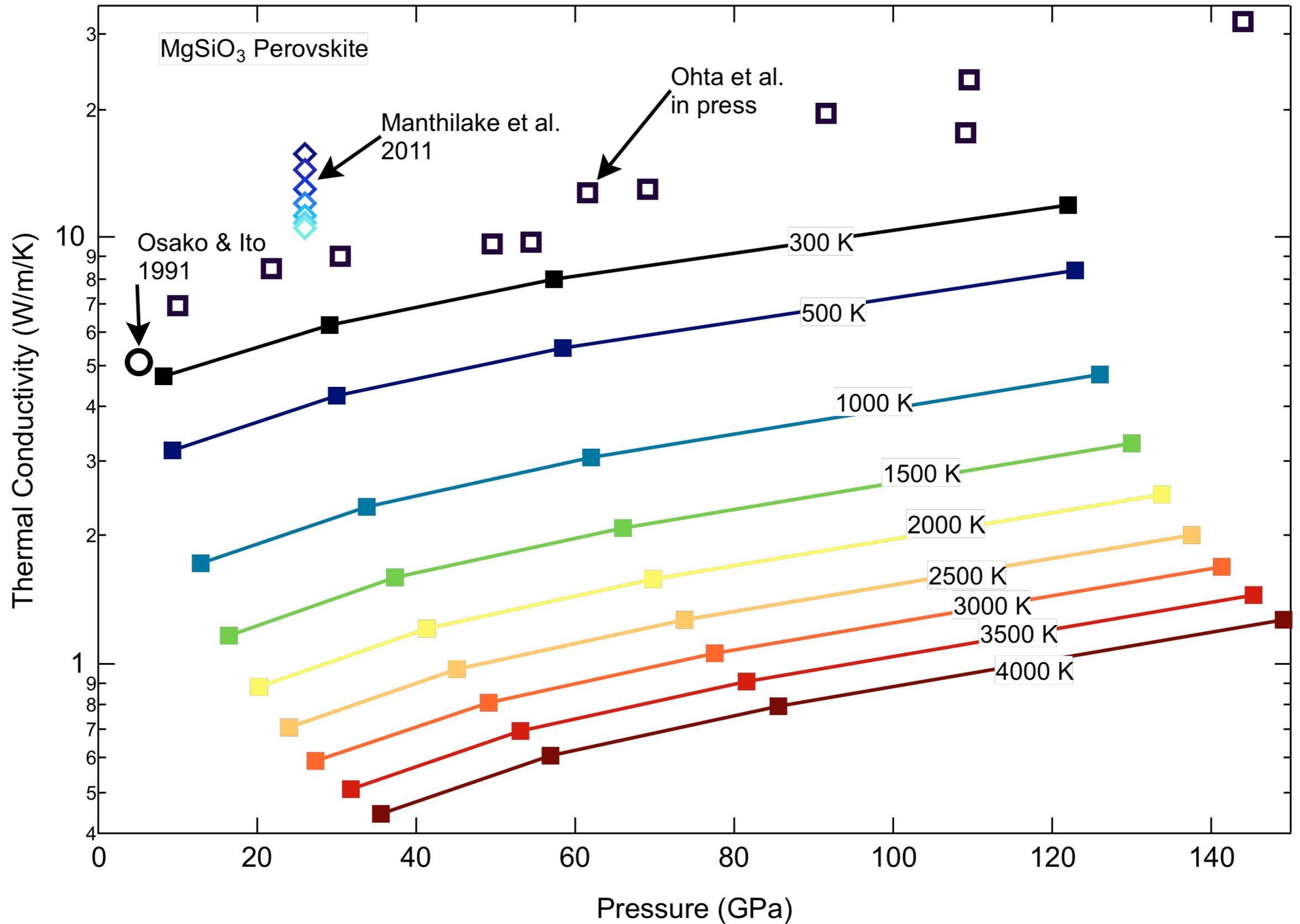
# Mode contributions: optic vs. acoustic

$$\kappa = \kappa_{acoustic} + \kappa_{optic}$$



quick quiz: which is MgO?  
which is MgSiO<sub>3</sub>? why?

# Results--Thermal conductivity MgSiO<sub>3</sub>



# Additional considerations for Earth's mantle

## ***1. Effect of iron***

2. Radiative contribution to heat flow
3. Composite material

**✓1. Changes mass**

**✓2. Adds impurities**

**3. Chemical effects**

for now:

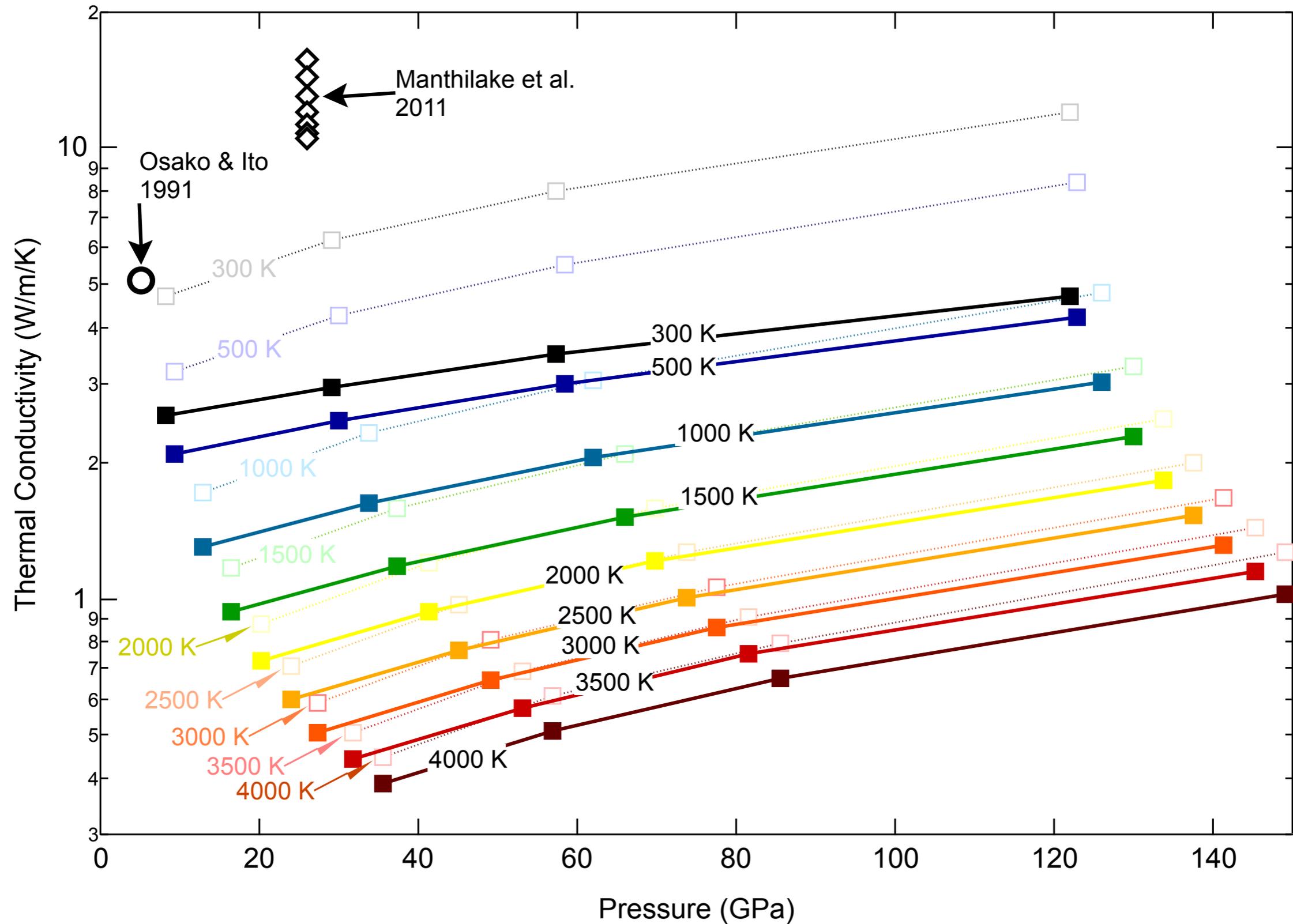
i. Fe<sup>+2</sup> substitution for Mg<sup>+2</sup>

ii. high spin in Pv

iii. low spin in MgO

*Reference : Moses C. Natm, Jianjun Dong,  
manuscript in preparation*

# Thermal conductivity (Mg<sub>0.85</sub>Fe<sub>0.15</sub>)SiO<sub>3</sub>



# Additional considerations for Earth's mantle

1. Effect of iron

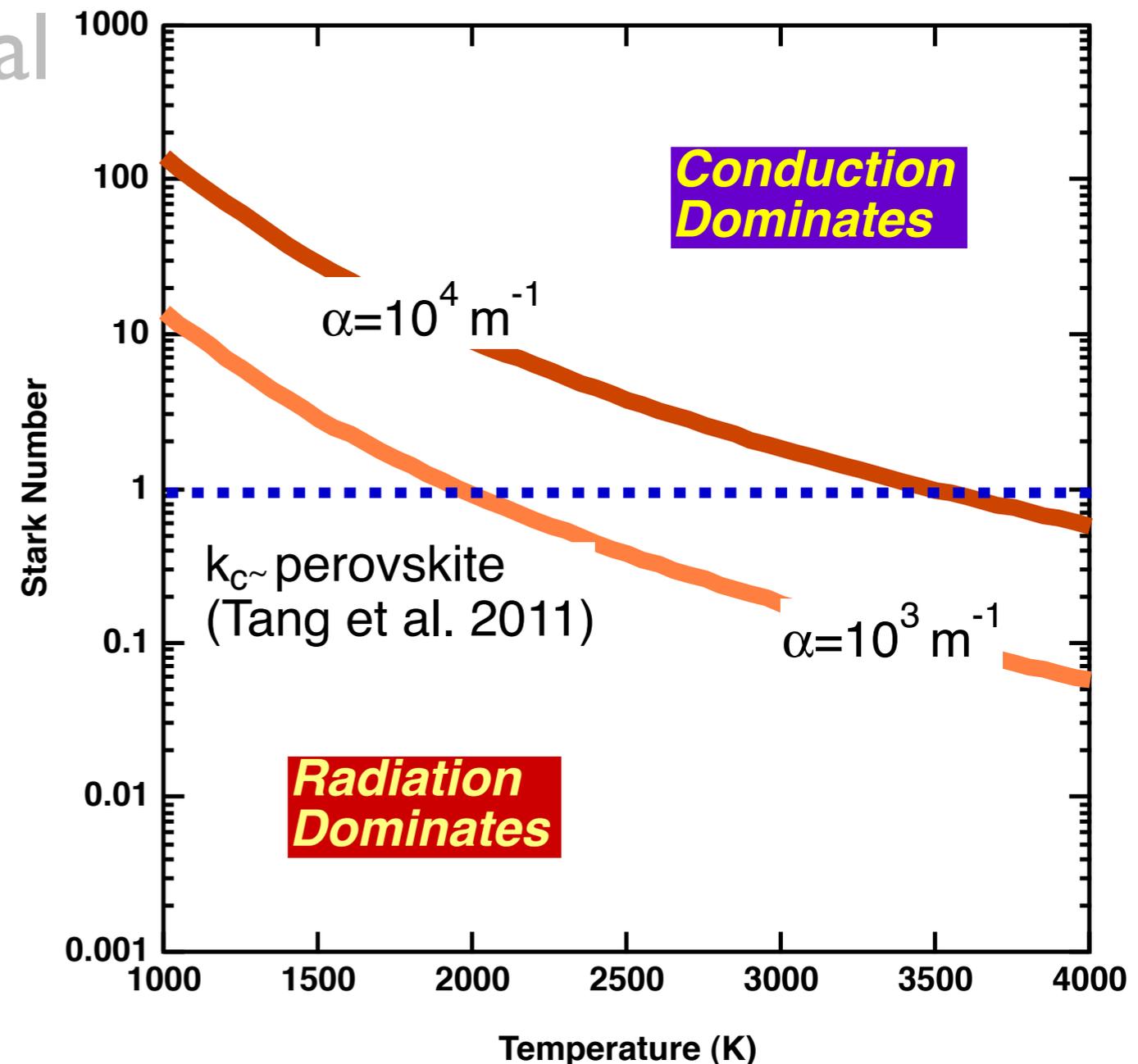
**2. Radiative contribution to heat flow**

3. Composite material

Stark Number:  
Non dimensional  
parameter to determine  
relative importance of  
radiative transport

$$N = k_{cond} \alpha_{abs} / 4\sigma T^3$$

$1/\alpha_{abs} \sim$  absorption  
length scale

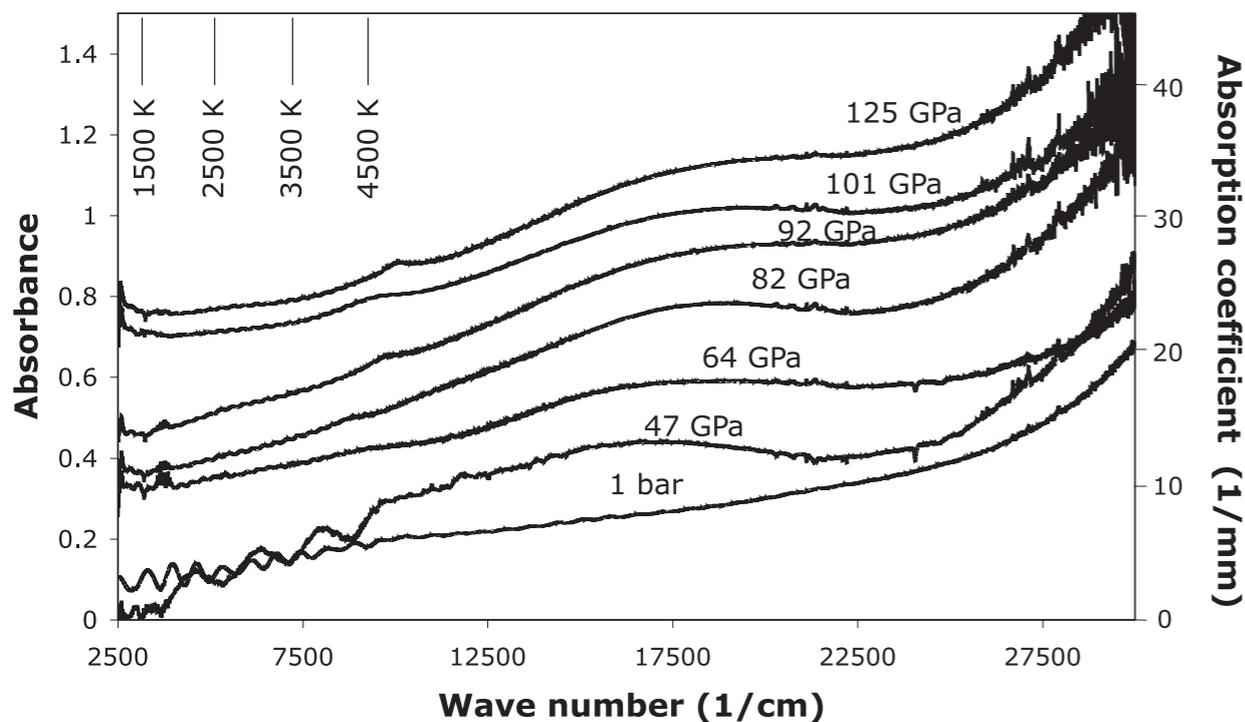


# Competing Measurements of Absorption Coefficients

**Science, 2008**

## Optical Absorption and Radiative Thermal Conductivity of Silicate Perovskite to 125 Gigapascals

Hans Keppler,<sup>1\*</sup> Leonid S. Dubrovinsky,<sup>1</sup> Olga Narygina,<sup>1</sup> Innokenty Kantor<sup>1,2</sup>



**Nature, 2008**

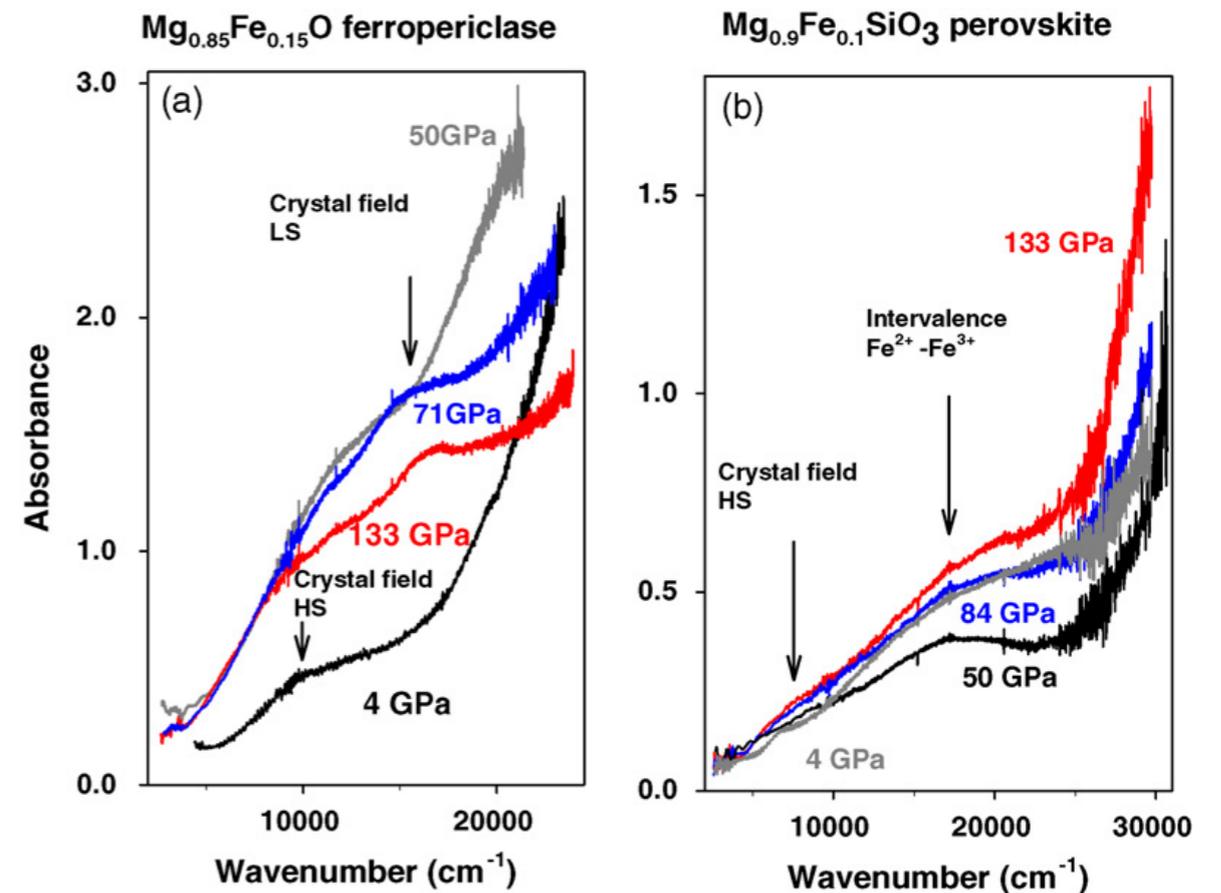
## Thermal conductivity of lower-mantle minerals

Alexander F. Goncharov<sup>a,\*</sup>, Pierre Beck<sup>a</sup>, Viktor V. Struzhkin<sup>a</sup>, Benjamin D. Haugen<sup>a,b</sup>, Steven D. Jacobsen<sup>c</sup>

<sup>a</sup> Geophysical Laboratory, Carnegie Institution of Washington, 5251 Broad Branch Road N.W., Washington, DC 20015-1305, USA

<sup>b</sup> Department of Geological Sciences, University of Colorado, Boulder, CO 80309, USA

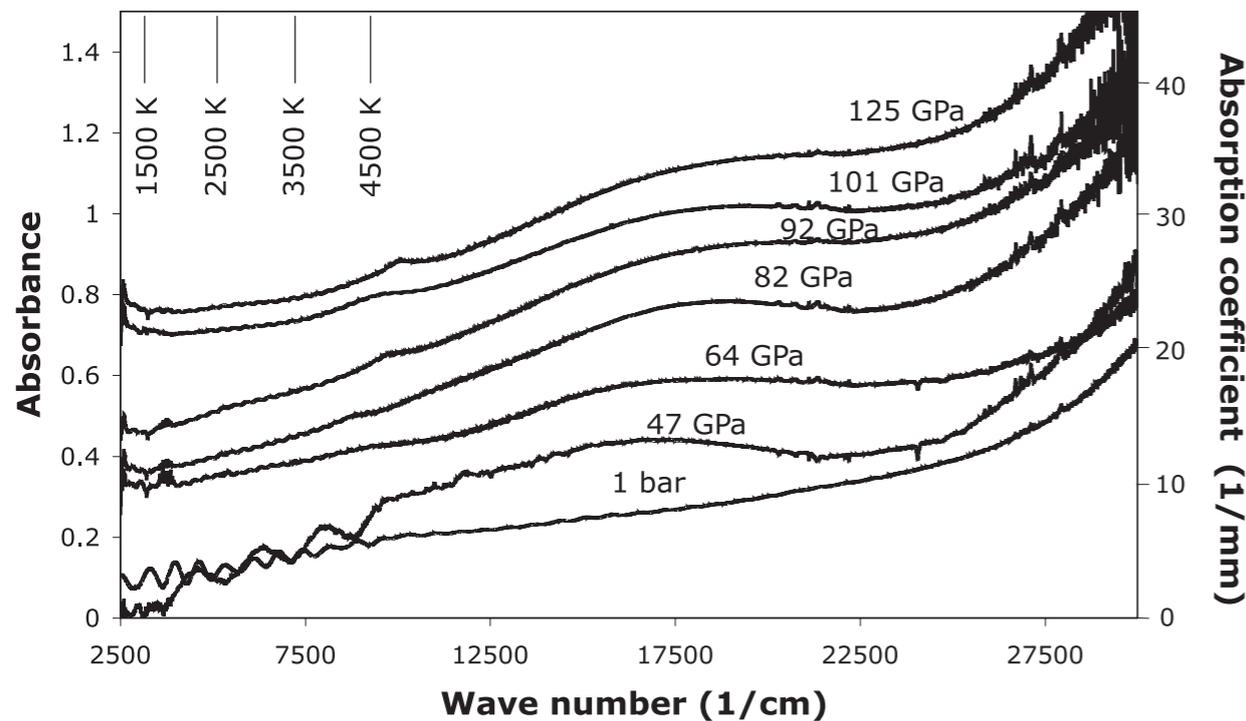
<sup>c</sup> Department of Earth and Planetary Sciences, Northwestern University, Evanston, IL 60208, USA



# Competing Measurements of Absorption Coefficients

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**“pristine mantle”  
radiative heat flow  
is important**

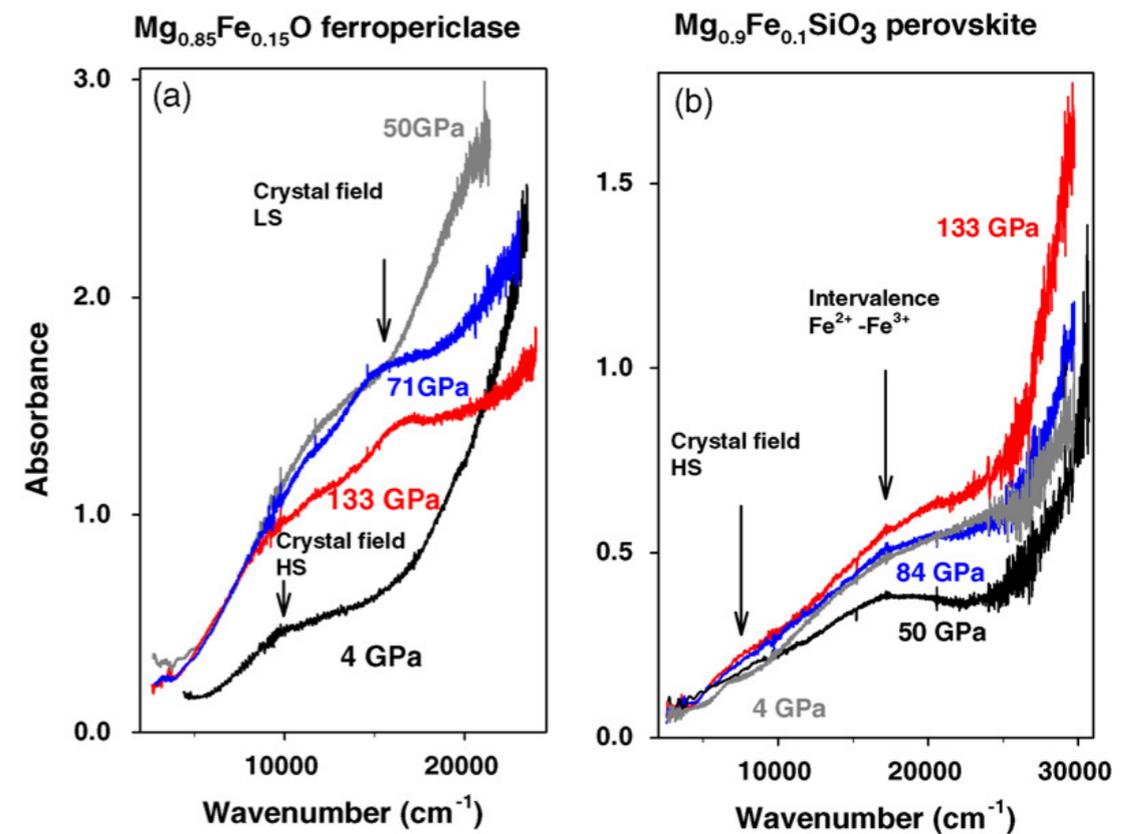
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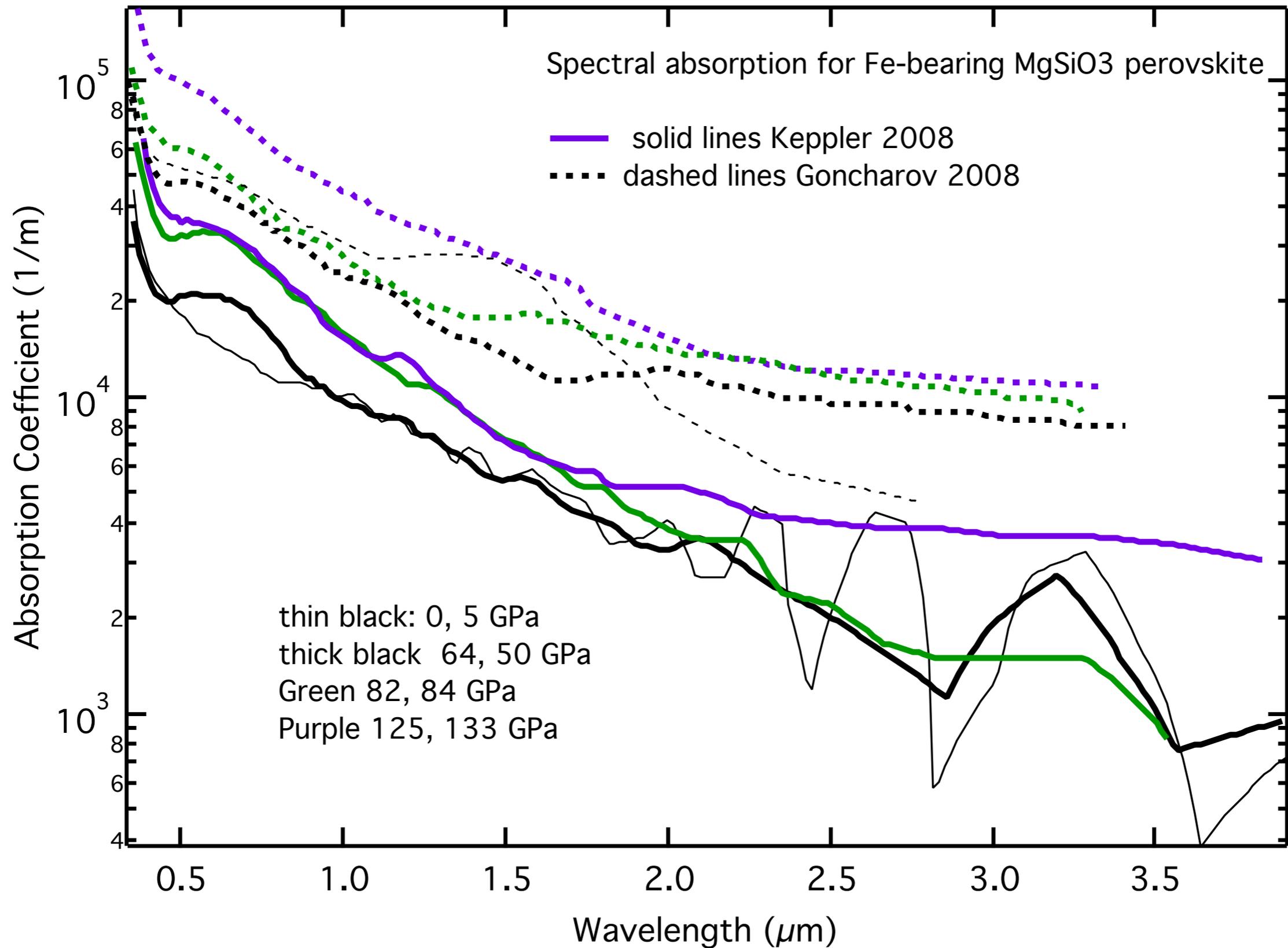
<sup>b</sup> Department of Geological Sciences, University of Colorado, Boulder, CO 80309, USA

<sup>c</sup> Department of Earth and Planetary Sciences, Northwestern University, Evanston, IL 60208, USA



**“dirty mantle”  
radiative heat flow  
is not important**

# step 1: plot everything on same plot



# step 2: Re-analyze data

radiative contribution to thermal conductivity calculated via Rosseland approximation (valid for optically thick)

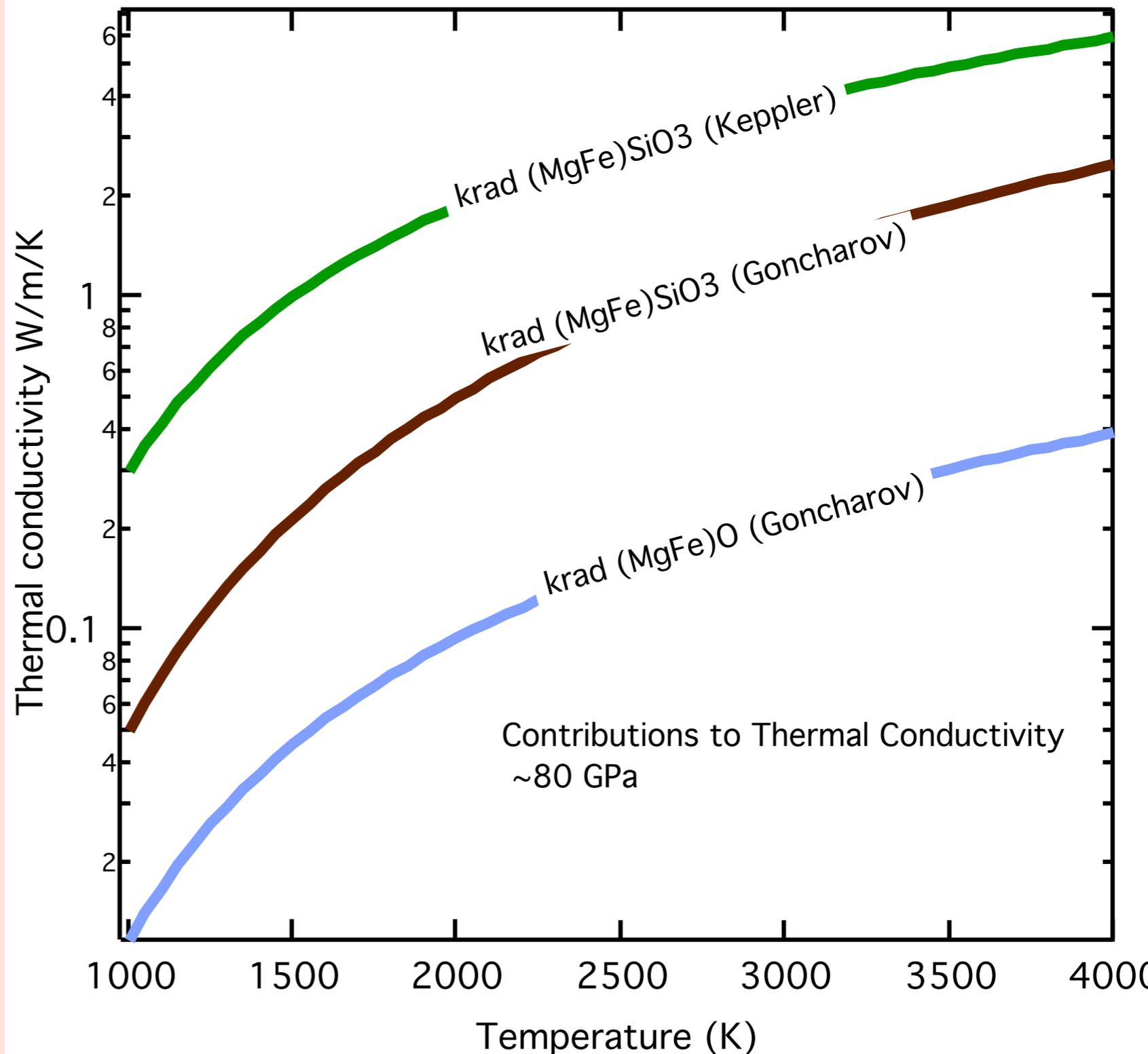
$$k_{rad}(T) = \frac{16n^2\sigma T^3}{3\beta_R}$$

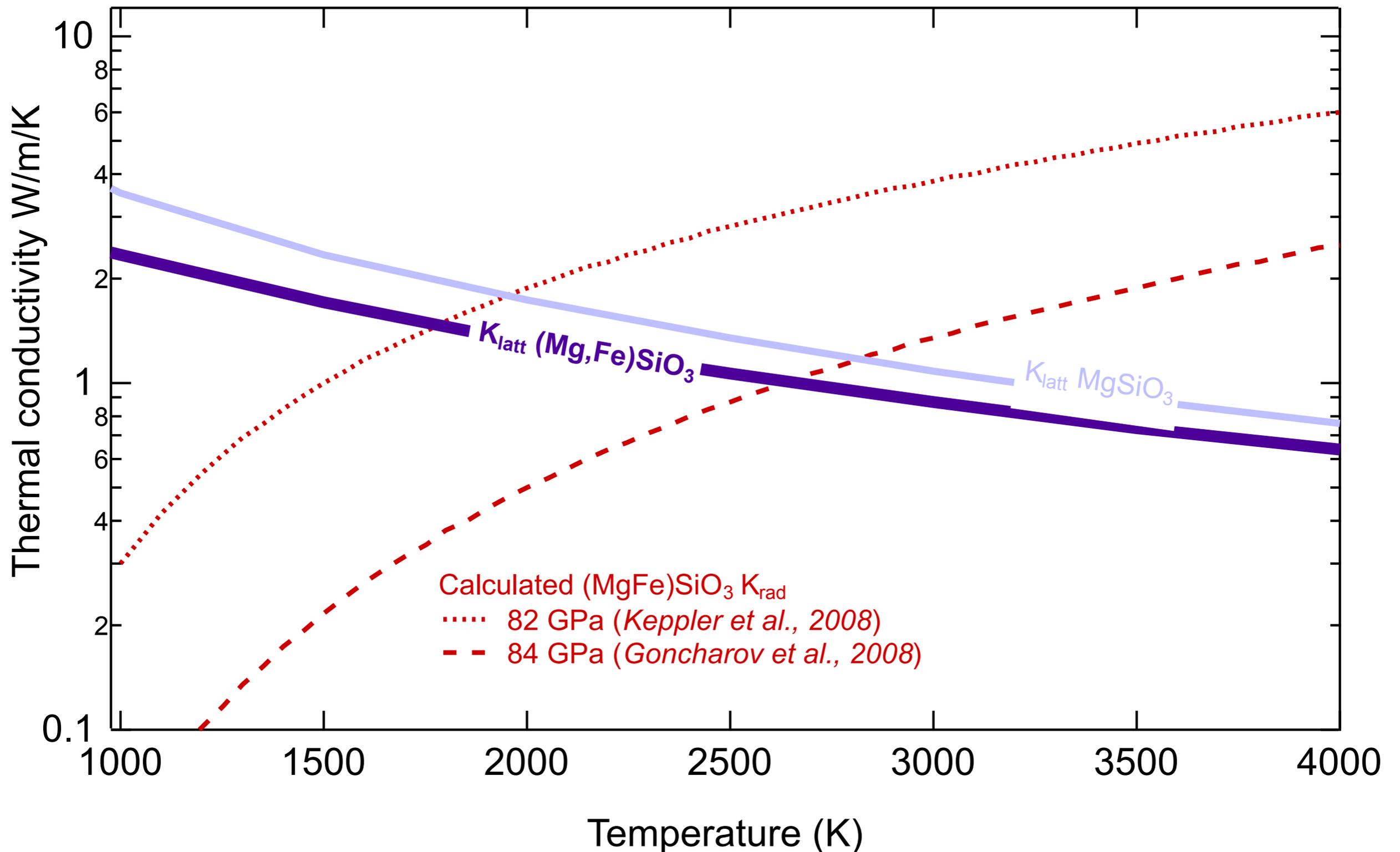
n: Index of refraction

T: Temperature

$\beta$ : Rosseland mean extinction coefficient

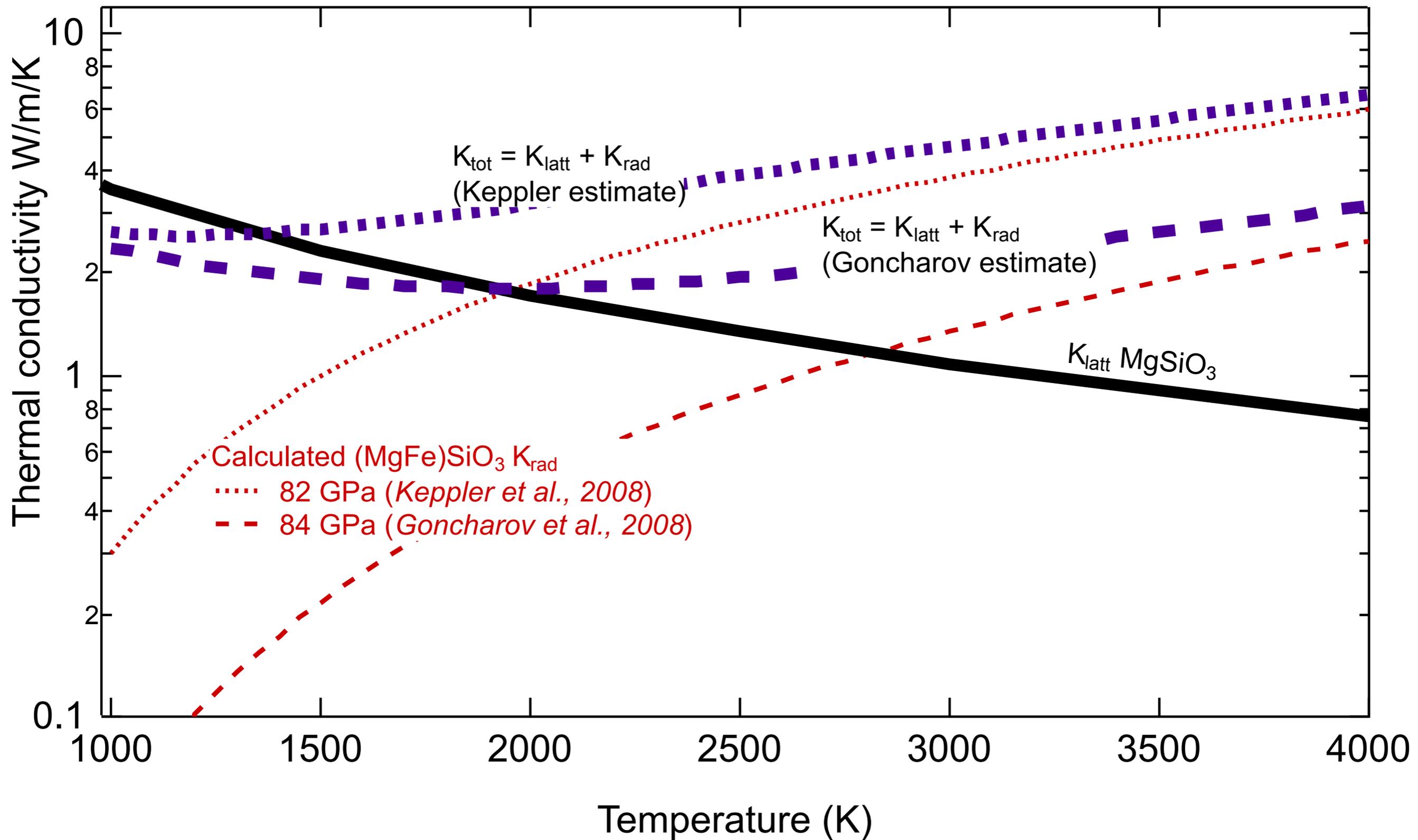
$$\frac{n^2}{\beta_R} = \frac{\pi}{4\sigma T^3} \int_0^\infty \frac{n_\lambda^2}{\beta_\lambda} \frac{dI_{b,\lambda}}{dT} d\lambda.$$



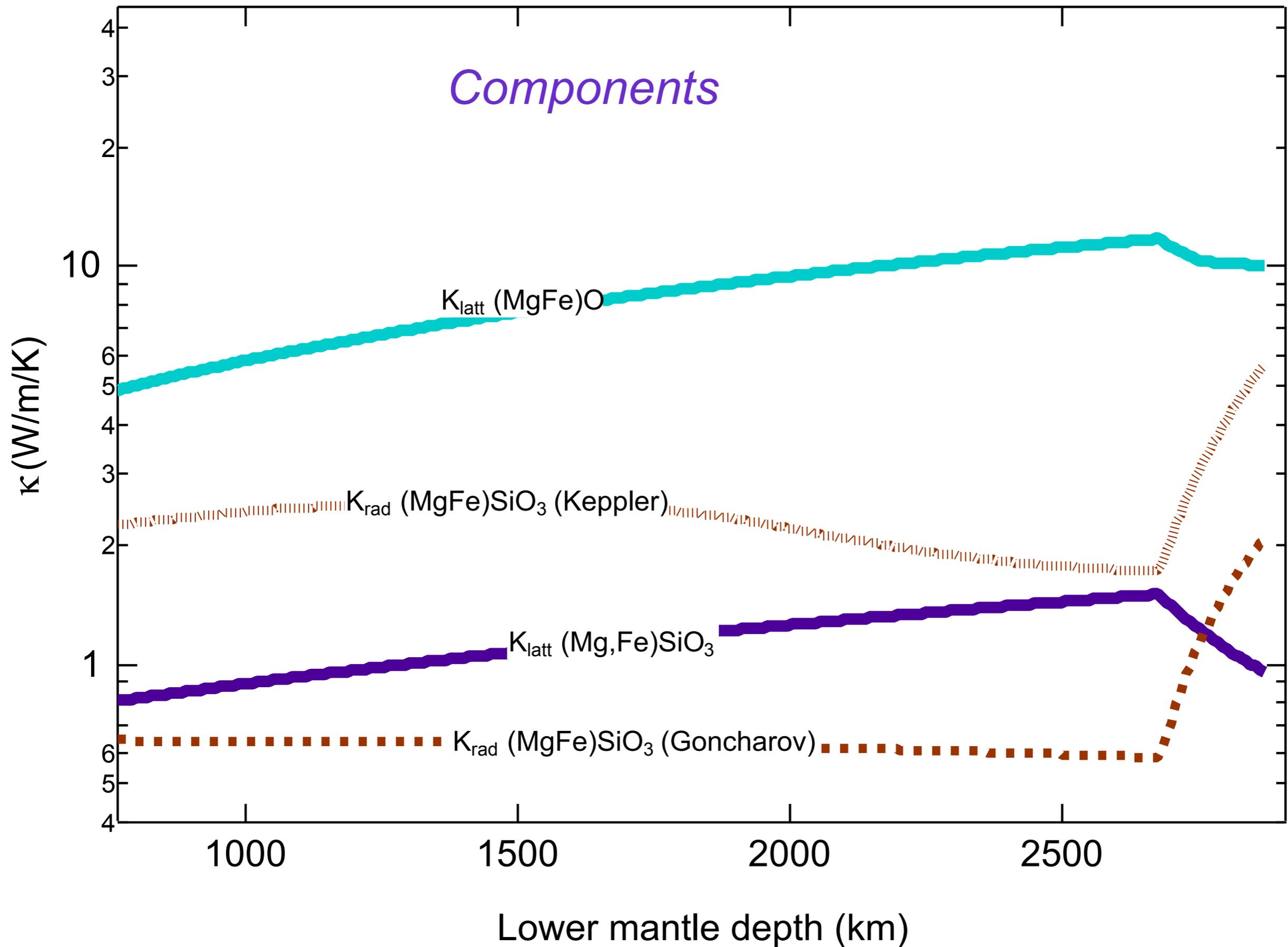


Lattice and radiative contributions to thermal conductivity have:  
 Similar magnitude  
 Opposite temperature dependence

# Total thermal conductivity of perovskite



# Depth-Dependent Contributions to Conductivity



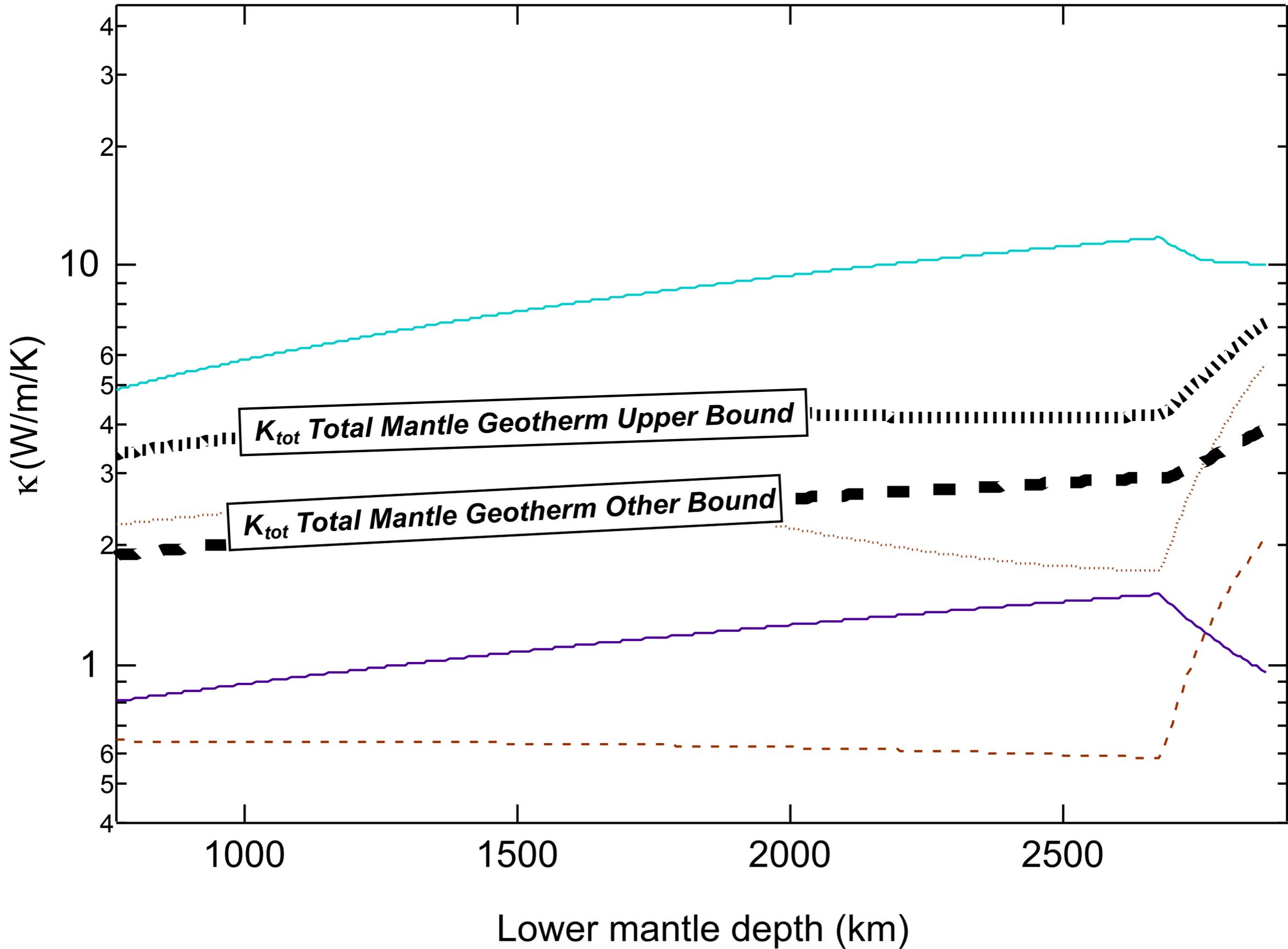
# Additional considerations for Earth's mantle

1. Effect of iron
2. Radiative contribution to heat flow
- 3. Composite material**

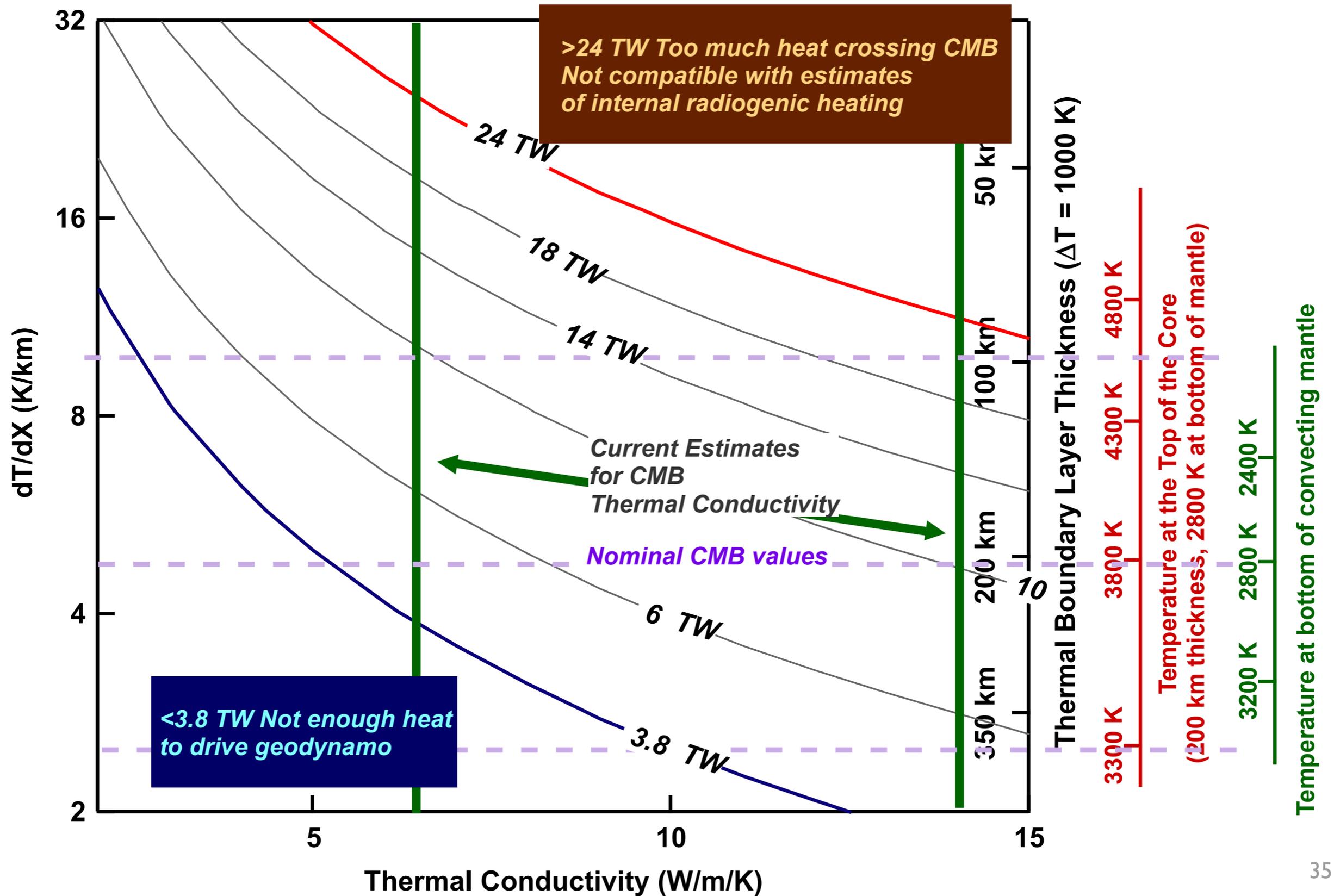
★ Maxwell-Garnett  
Composite

	$k_{latt}$	$k_{rad}$	
<b>MgSiO<sub>3</sub></b>	P, T, Fe	P, T, Fe	$k_{latt} + k_{rad}$ MgSiO <sub>3</sub>
<b>MgO</b>	P, T, Fe	P, T, Fe	★ $k_{latt} + k_{rad}$ MgO
<p><i>length scales:</i>  <math>k_{latt} &lt; 10 \text{ nm}</math>  <math>k_{rad} &lt; 100 \mu\text{m}</math>                      grain size <math>\sim 1 \text{ mm}</math></p>			

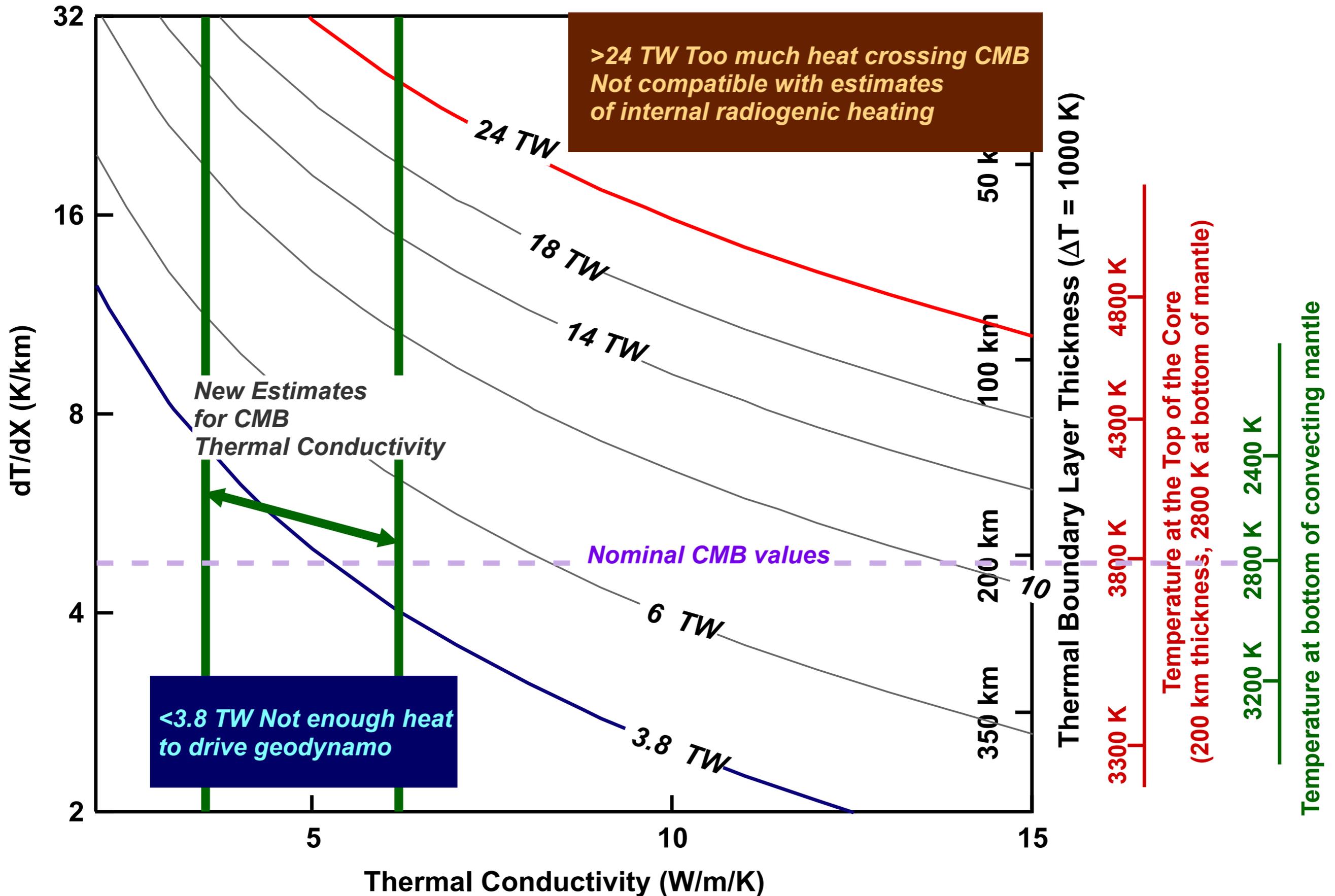
# Best Estimates Earth Mantle Thermal Conductivity



# A reminder--before this study



# Implications for Core/Mantle Boundary Heat flow



# Constraints on the thermal state of the Earth

***~47(3) TW***

Crustal  
Radiogenic  
Contribution  
***~7 TW***

A larger radiogenic  
heat budget?

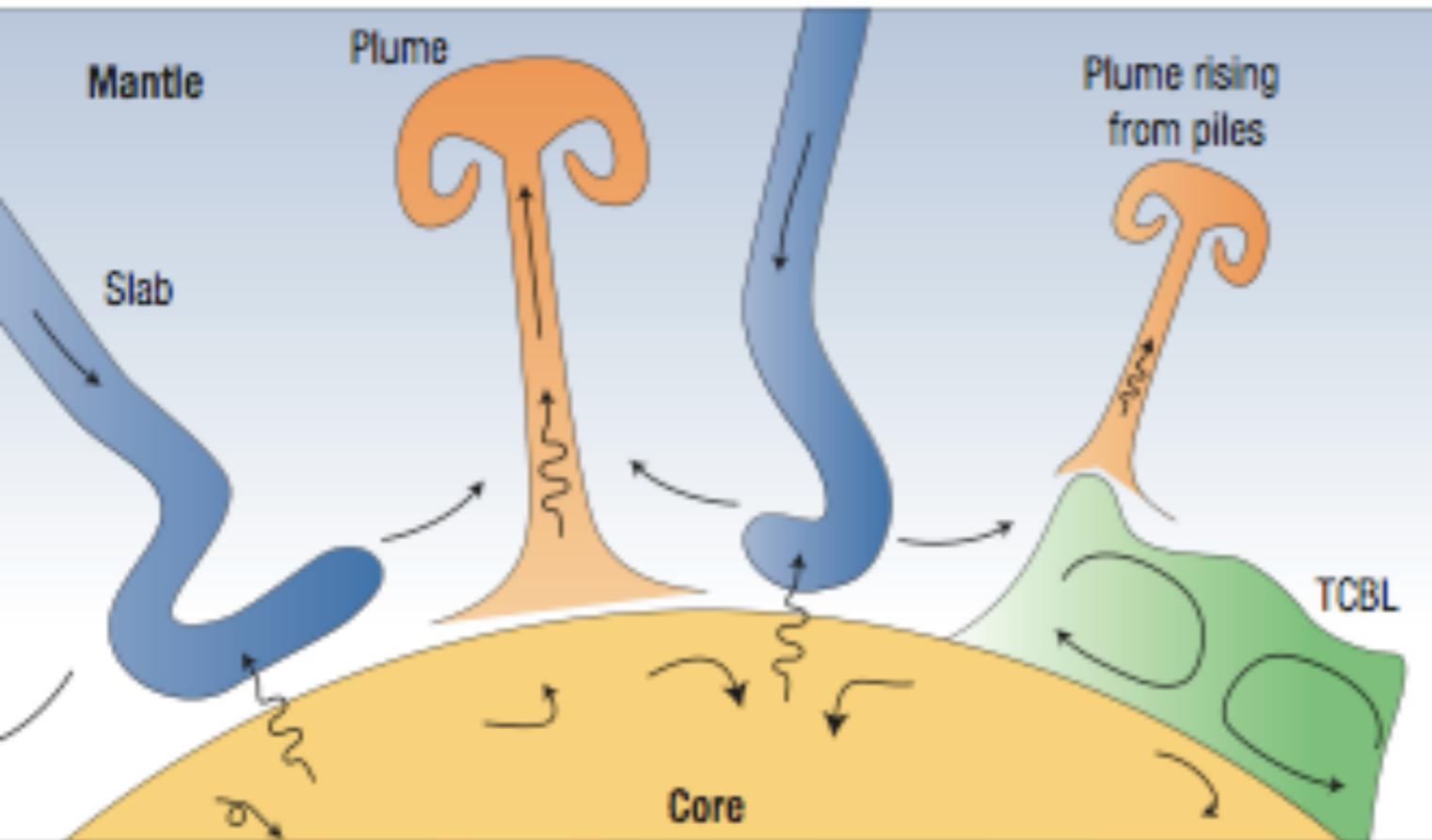
***~20 TW***

An additional  
thermal boundary  
layer in mantle?

***~17 TW***

New CMB heat flow  
***~3-6 TW***

# Resolution of the Goldilocks story: “Just Right” K



**Lay, Hernlund, Buffett, 2008**

## **Simple convecting mantle**

Thermal conductivity has  
\*weak\* pressure and  
temperature dependencies

## **Core/Mantle boundary is an insulating area**

1. Large temperature drop at core/mantle boundary
2. Partially molten lower mantle possible
3. Additional thermal boundary layer in the mantle or more radiogenic heating

## **Inner core timing**

Relaxes constraints on timing  
of inner core formation (e.g.  
Buffett, 2009)

# General Thermal Conductivity: General Boltzmann Transport Equation

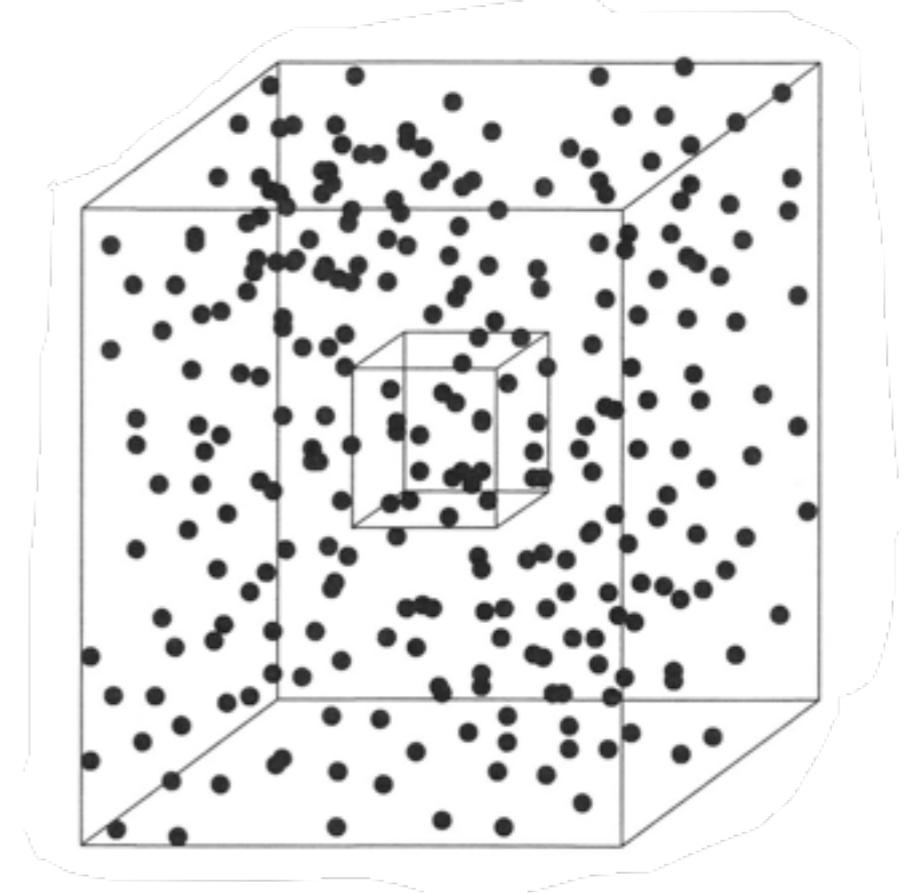
1. each particle defined by:  
position  $\mathbf{x}$  and momentum  $\mathbf{p}$

2. define a probability density function  
 $f(\mathbf{x}, \mathbf{p}, t)$  such that:  $dN = f(\mathbf{x}, \mathbf{p}, t) d^3\mathbf{x} d^3\mathbf{p}$

$$3. \frac{\partial f}{\partial t} = \left( \frac{\partial f}{\partial t} \right)_{diffusion} + \left( \frac{\partial f}{\partial t} \right)_{force} + \left( \frac{\partial f}{\partial t} \right)_{collisions}$$

4a. for particles with classical statistics:  
 $f(t_0) \sim \exp(-\mathbf{E}/kT)$

4a. for particles with quantum statistics:  
 $f(t_0) \sim (\exp(\mathbf{E}/kT) - 1)^{-1}$



N particles

**Distribution  
functions are  
temperature-  
dependent**

# Kinetic Transport Equation

$$\kappa = \frac{1}{3} \sum_{i, \vec{q}} c_V(i, \vec{q}) v_g^2(i, \vec{q}) \tau(i, \vec{q})$$

$$\kappa = C_V v_g^2 \tau$$

Quasi-Harmonic  
Approximation

$$C_V, v_g$$

Quantum  
Scattering Theory

$$\tau$$

2<sup>nd</sup>-order force  
constant matrix:

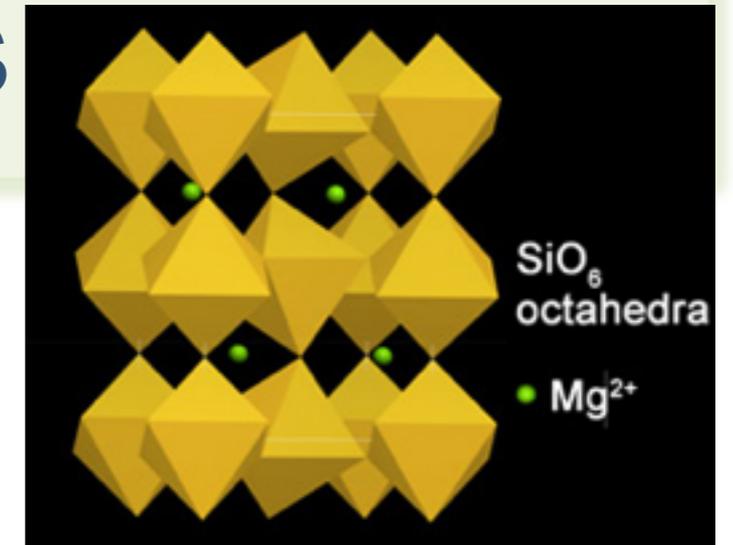
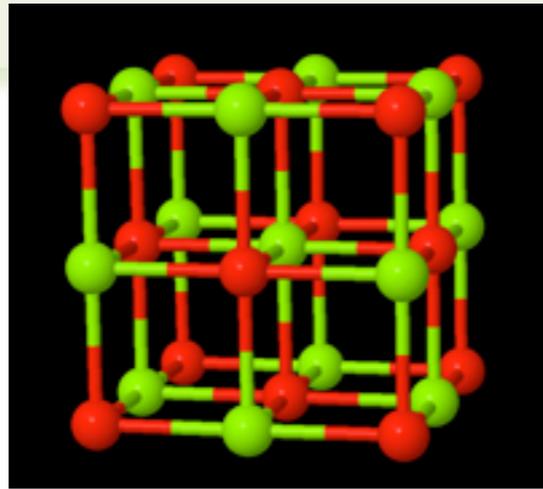
$$\phi_{ij} = \frac{\partial^2 E}{\partial x_i \partial x_j}$$

3<sup>rd</sup>-order lattice  
Anharmonicity:

$$A_{ijk} = \frac{\partial^3 E}{\partial x_i \partial x_j \partial x_k}$$

First-Principles calculation of atomic forces: Finite-Displacement Super-Cell Method

# Simulation Details

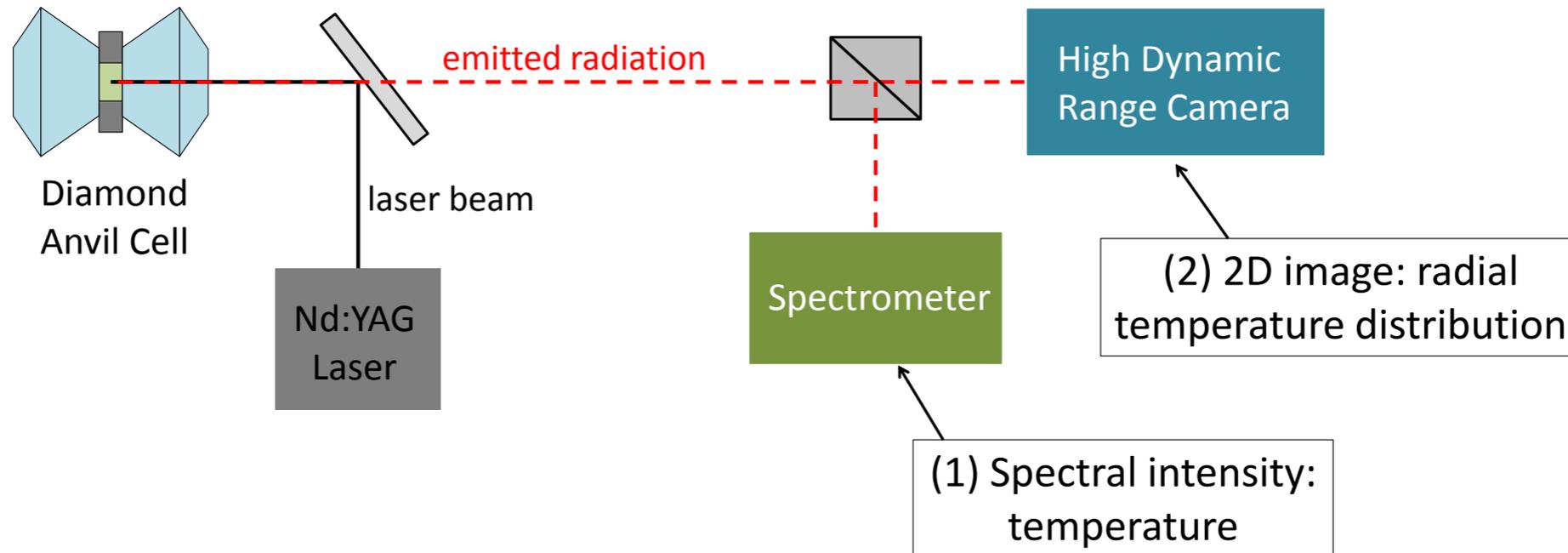


	MgO (Tang & Dong, 2009, 2010)	MgSiO <sub>3</sub>
Static Energy	LDA, PAW, PW; 2-atom unit cell; 12x12x12 Monkhorst k-grid	LDA, PAW, PW; 20-atom unit cell; 8x8x6 Monkhorst k-grid
Forces	128-atom super cell; Gamma point sampling	160-atom super cell; Gamma point sampling;
Conductivity ( $C_v$ , $V_g$ , $\tau$ )	16X16X16 q-point sampling; 5 X 9 grids in V-T domain	8X8X6 q-point sampling; 5 X 9 grids in V-T domain

# Current research & ongoing questions

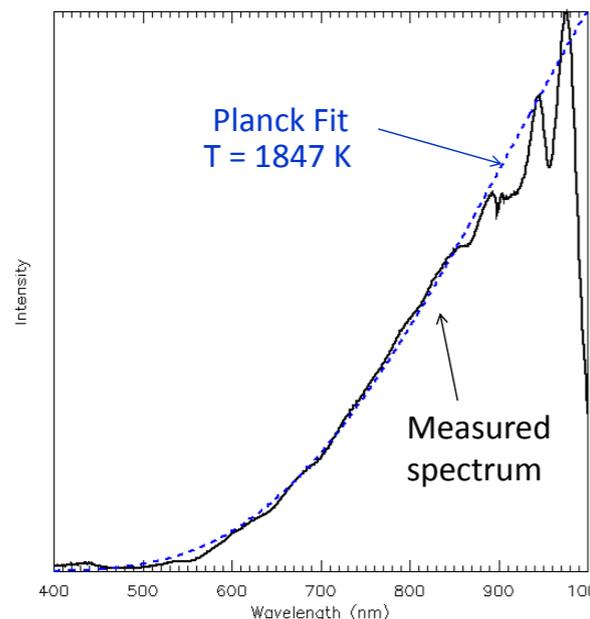
## LASER HEATING EXPERIMENTS

### UCLA Mineral Physics Lab Laser Heating System

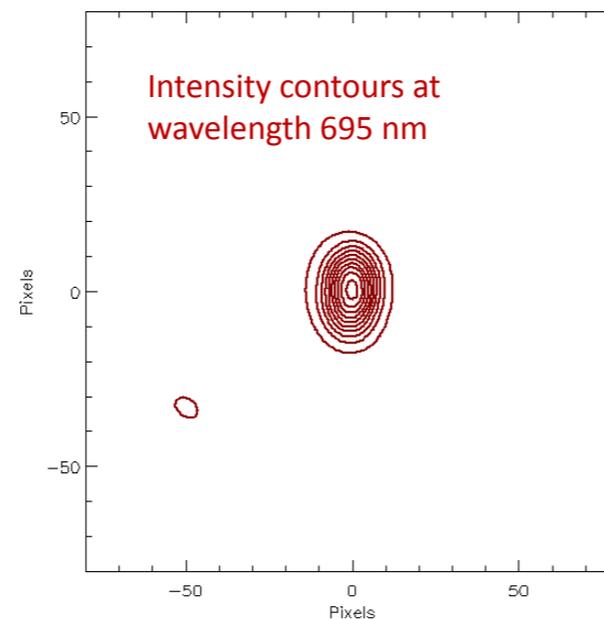


### Laser Heating Measurements

(1) Hotspot Spectrum



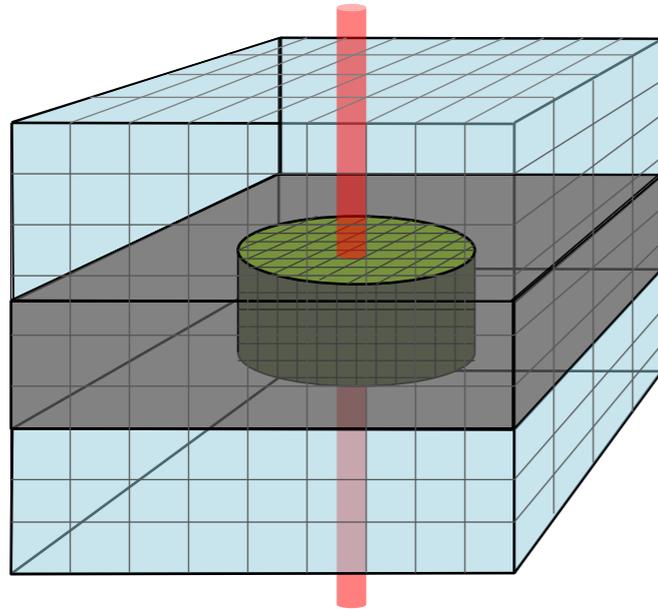
(2) 2D Hotspot Image



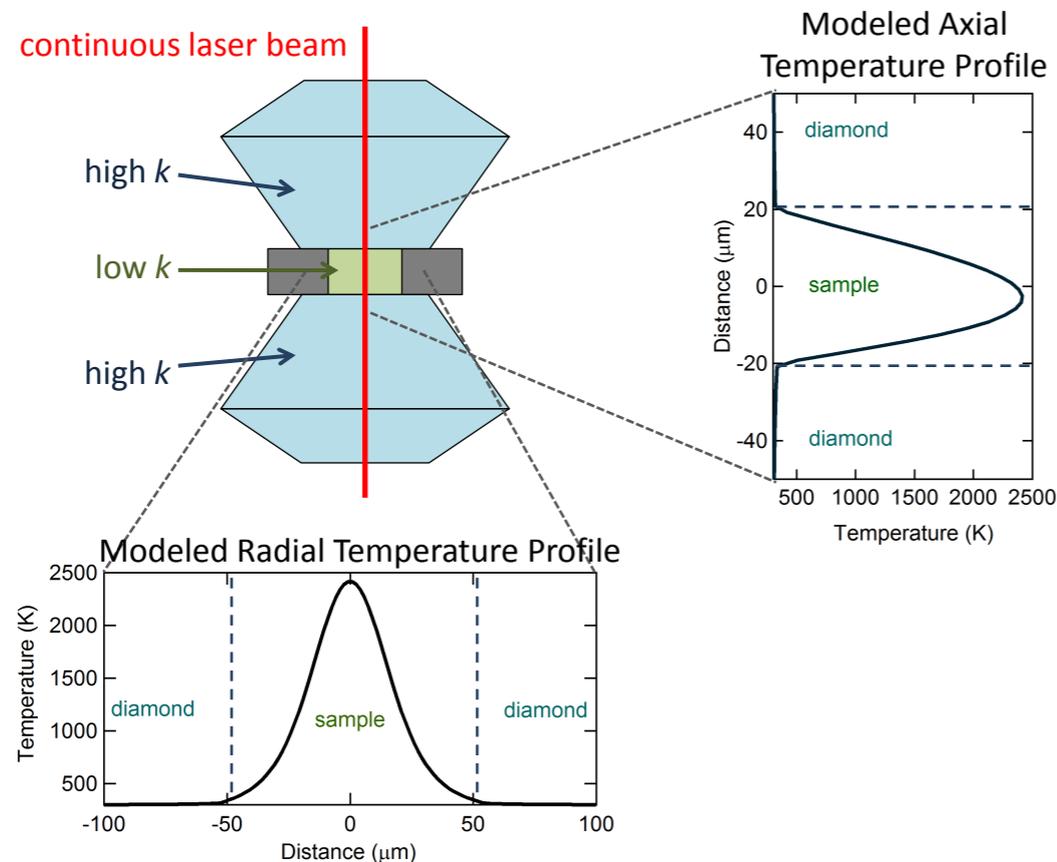
- Data taken during laser heating of AgI at pressure of 17 GPa in the LHDAC
- Planck fit for measured spectrum between 600 and 900 nm
- Entire hotspot spectrum was collected in order to minimize potential biases due to misalignment
- 2D hotspot intensity measured through bandpass filter with center wavelength 695 nm and width of 70 nm

# Current research & ongoing questions

## DIAMOND ANVIL CELL HEAT FLOW MODEL



- 3-D numerical model solves the steady-state heat equation using a full approximation storage multi-grid solver.
- Includes flexible laser and sample geometries and temperature-dependent material properties.
- Based on previous model for heat flow in multi-anvil press [1].

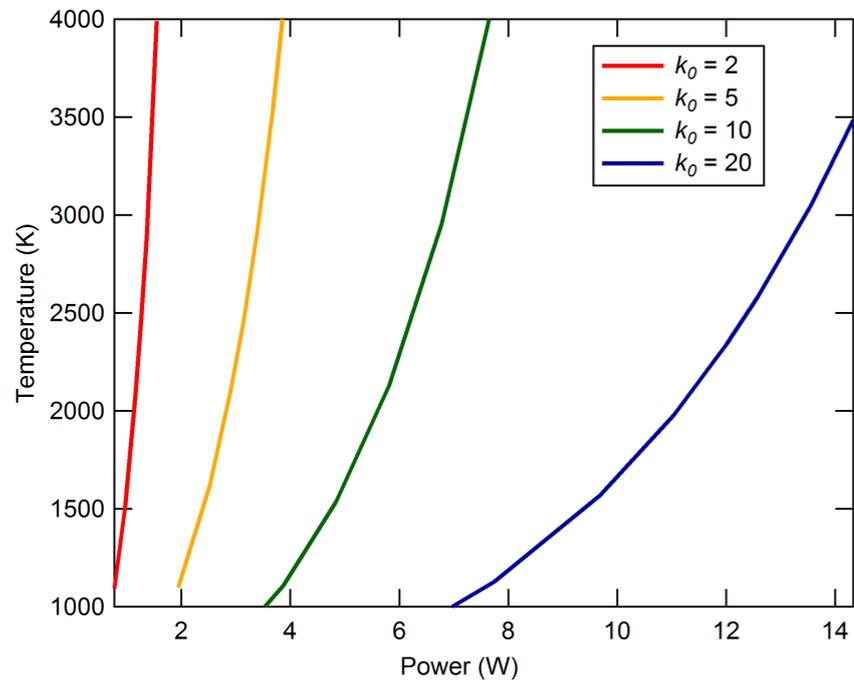


- Large 3-D temperature gradients form during continuous laser heating in the LHDAC.
- Temperature distribution depends on laser and sample geometry, and sample thermal conductivity and absorption properties.

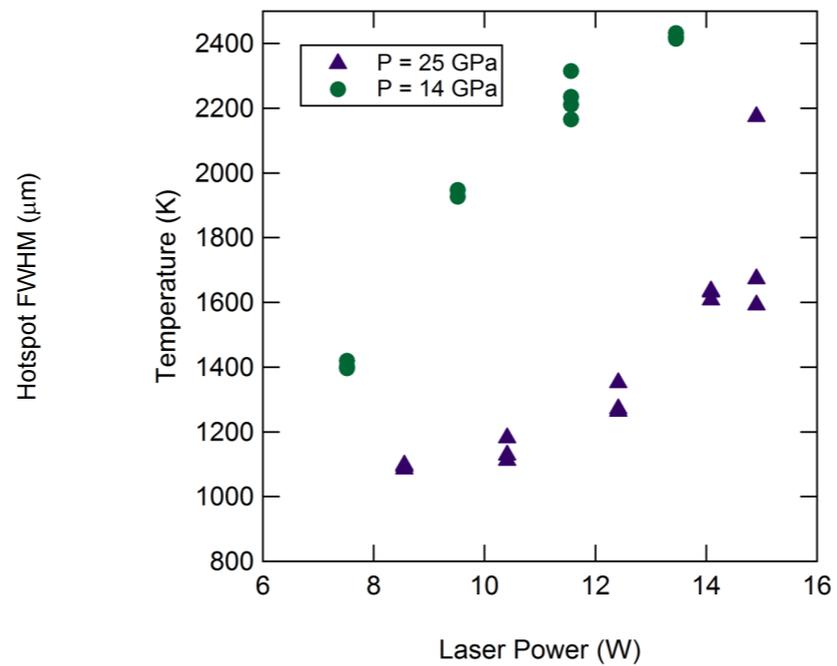
# Current research & ongoing questions

## PRELIMINARY EXPERIMENTAL DATA

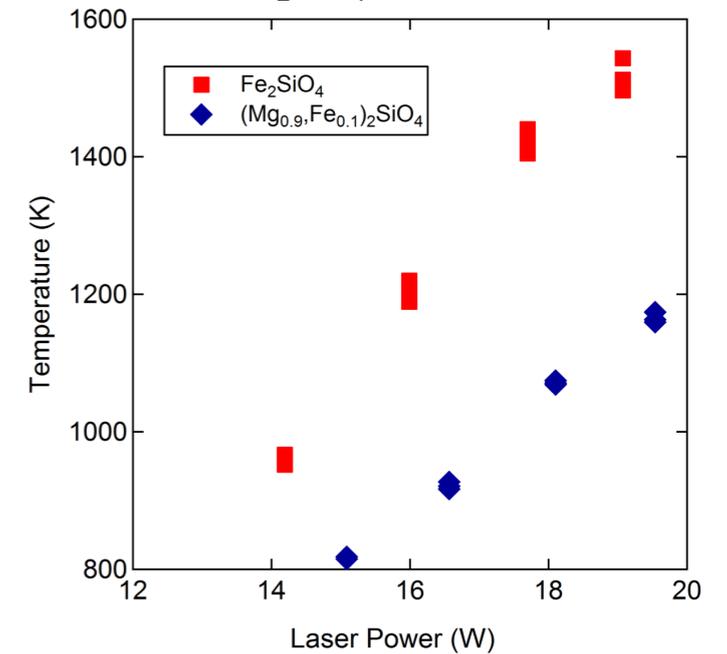
### Modeled Temperature vs. Power



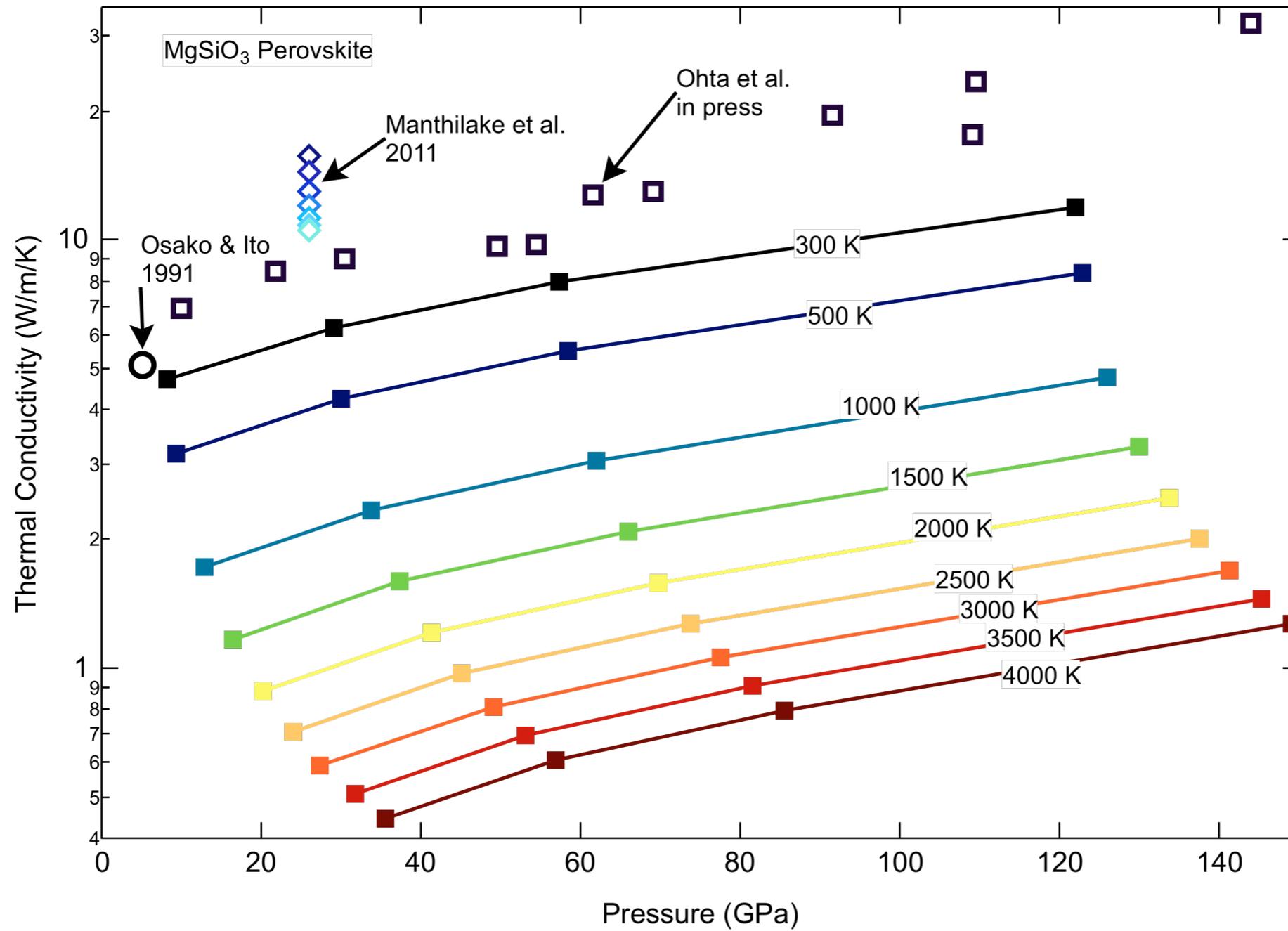
### Temperature vs. Laser Power for $(\text{Mg}_{0.9}\text{Fe}_{0.1})_2\text{SiO}_4$ at 14 and 24 GPa



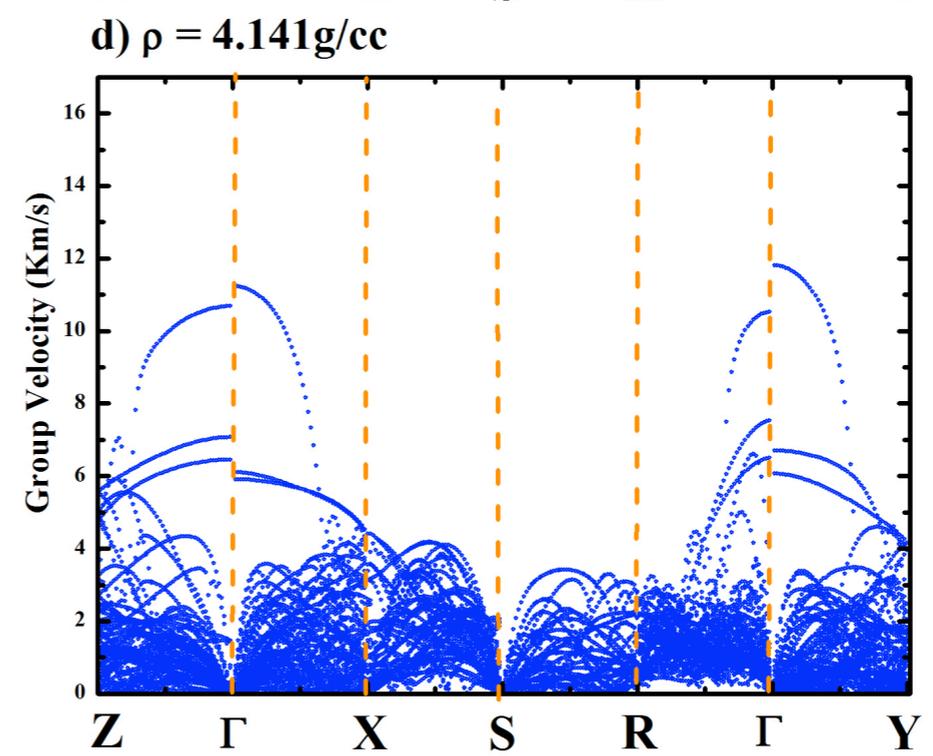
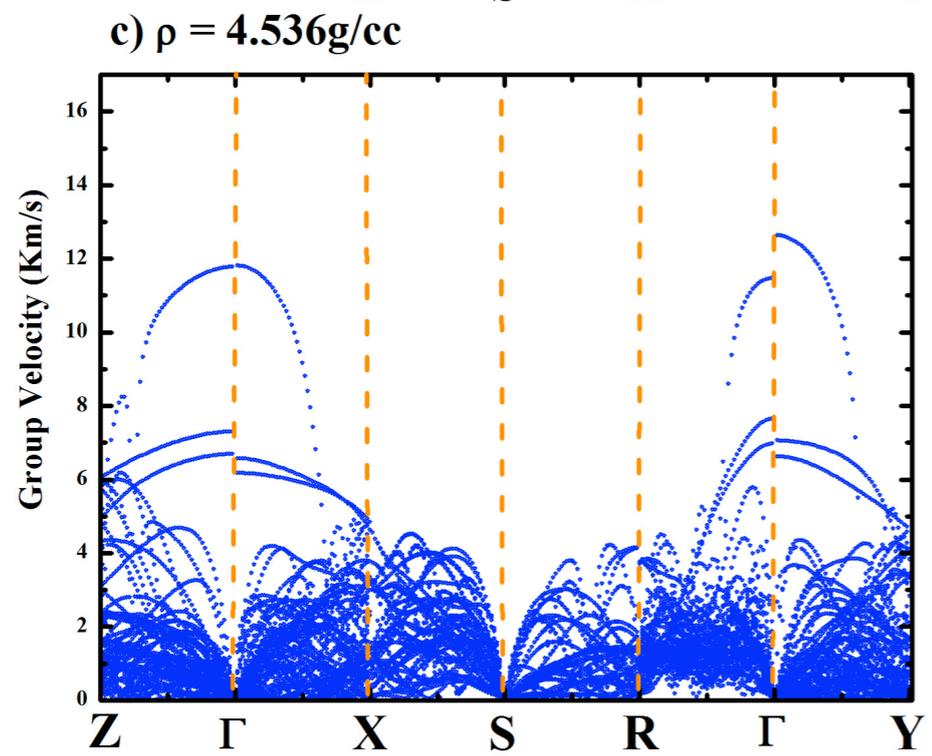
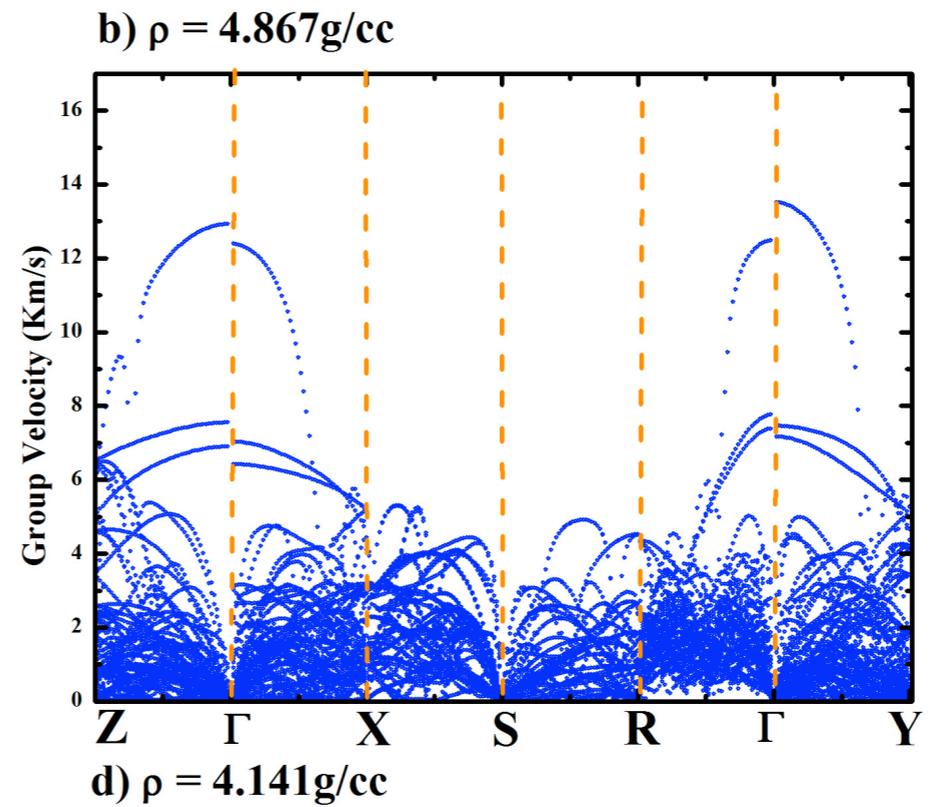
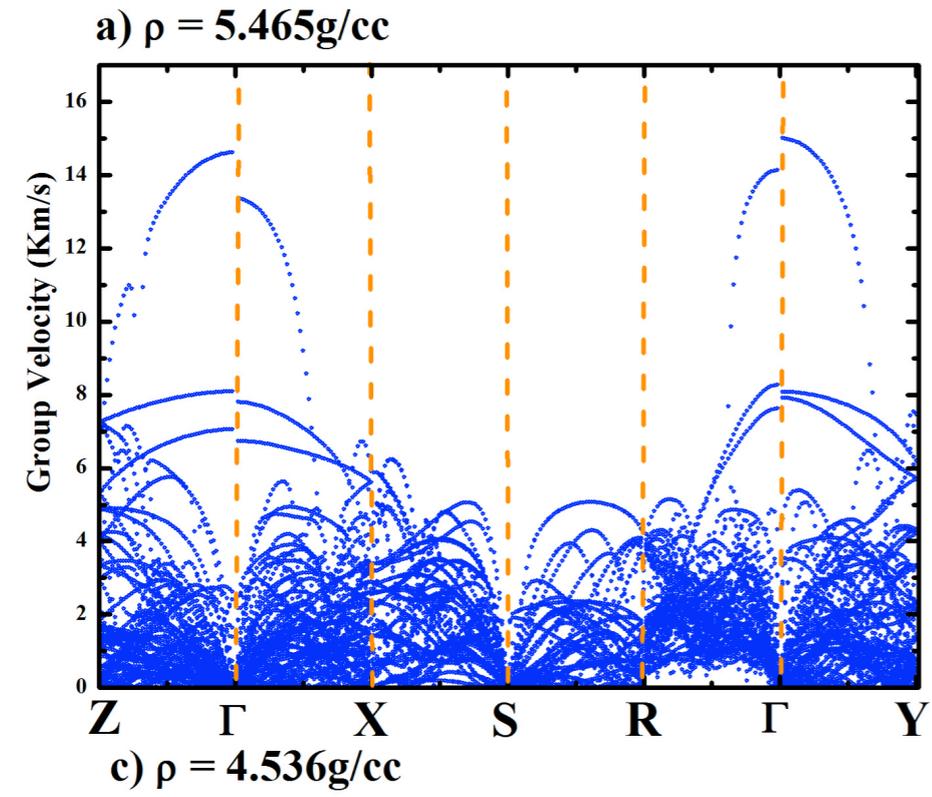
### Temperature vs. Laser Power for $(\text{Mg,Fe})_2\text{SiO}_4$ with variable Fe



# Thermal Conductivity of Perovskite



# pressure-dependent phonon calcs



# Temperature-dependence

